MULTIDIMENSIONAL LAPLACE TRANSFORMS OVER CAYLEY-DICKSON ALGEBRAS AND THEIR APPLICATIONS

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Multidimensional noncommutative Laplace transforms over octonions are studied. Theorems about direct and inverse transforms and other properties of the Laplace transforms over the Cayley-Dickson algebras are proved. Applications to partial differential equations including that of elliptic, parabolic and hyperbolic type are investigated. Moreover, partial differential equations of higher order with real and complex coefficients and with variable coefficients with or without boundary conditions are considered.

Key Words: multidimensional noncommutative Laplace transform, Cayley-Dickson algebras, partial differential equations, boundary conditions

1 Introduction.

The Laplace transform over the complex field is already classical and plays very important role in mathematics including complex analysis and differential equations [29, 12, 23]. The classical Laplace transform is used frequently for ordinary differential equations and also for partial differential equations sufficiently simple to be resolved, for example, of two variables. But it meets substantial difficulties or does not work for general partial differential equations even with constant coefficients especially for that of hyperbolic type.

To overcome these drawbacks of the classical Laplace transform in the present paper more general noncommutative multiparameter transforms over Cayley-Dickson algebras are investigated. In the preceding paper a noncommutative analog of the classical Laplace transform over the Cayley-Dickson algebras was defined and investigated [18]. This paper is devoted to its generalizations for several real parameters and also variables in the Cayley-Dickson algebras. For this the preceding results of the author on holomorphic, that is (super)differentiable functions, and meromorphic functions of the Cayley-Dickson numbers are used [17, 16]. The super-differentiability of functions of Cayley-Dickson variables is stronger than the Fréchet's differentiability. In those works also a noncommutative line integration was investigated.

We remind that quaternions and operations over them had been first defined and investigated by W.R. Hamilton in 1843 [8]. Several years later on Cayley and Dickson had introduced generalizations of quaternions known now as the Cayley-Dickson algebras [2, 9, 11, 25]. These algebras, especially quaternions and octonions, have found applications in physics. They were used by Maxwell, Yang and Mills while derivation of their equations, which they then have rewritten in the real form because of the insufficient development of mathematical analysis over such algebras in their time [4, 7, 13]. This is important, because noncommutative gauge fields are widely used in theoretical physics [27].

Each Cayley-Dickson algebra $\mathcal{A}_r$ over the real field $\mathbf{R}$ has 2^r generators $\{i_0, i_1, ..., i_{2^r-1}\}$ such that $i_0 = 1$, $i_j^2 = -1$ for each $j = 1, 2, ..., 2^r-1$, $i_j i_k = -i_k i_j$ for every $1 \le k \ne j \le 2^r-1$, where $r \ge 1$. The algebra $\mathcal{A}_{r+1}$ is formed from the preceding algebra $\mathcal{A}_r$ with the help of the so-called doubling procedure by generator i_{2^r}. In particular, $\mathcal{A}_1 = \mathbf{C}$ coincides with the field of complex numbers, $\mathcal{A}_2 = \mathbf{H}$ is the skew field of quaternions, $\mathcal{A}_3$ is the algebra of octonions, $\mathcal{A}_4$ is the algebra of sedenions. This means that a sequence of embeddings $... \hookrightarrow \mathcal{A}_r \hookrightarrow \mathcal{A}_{r+1} \hookrightarrow ...$ exists.

Generators of the Cayley-Dickson algebras have a natural physical meaning as generating operators of fermions. The skew field of quaternions is associative, and the algebra of octonions is alternative. The Cayley-Dickson algebra $\mathcal{A}_r$ is power associative, that is, $z^{n+m} = z^n z^m$

for each $n, m \in \mathbf{N}$ and $z \in \mathcal{A}_r$. It is non-associative and non-alternative for each $r \geq 4$. A conjugation $z^* = \tilde{z}$ of Cayley-Dickson numbers $z \in \mathcal{A}_r$ is associated with the norm $|z|^2 = zz^* = z^*z$. The octonion algebra has the multiplicative norm and is the division algebra. Cayley-Dickson algebras $\mathcal{A}_r$ with $r \geq 4$ are not division algebras and have not multiplicative norms. The conjugate of any Cayley-Dickson number z is given by the formula:

$(M1)$ $z^* := \xi^* - \eta \mathbf{l}$.

The multiplication in $\mathcal{A}_{r+1}$ is defined by the following equation:

$(M2)$ $(\xi + \eta \mathbf{l})(\gamma + \delta \mathbf{l}) = (\xi\gamma - \tilde{\delta}\eta) + (\delta\xi + \eta\tilde{\gamma})\mathbf{l}$

for each ξ, η, γ, $\delta \in \mathcal{A}_r$, $z := \xi + \eta \mathbf{l} \in \mathcal{A}_{r+1}$, $\zeta := \gamma + \delta \mathbf{l} \in \mathcal{A}_{r+1}$.

At the beginning of this article a multiparameter noncommutative transform is defined. Then new types of the direct and inverse noncommutative multiparameter transforms over the general Cayley-Dickson algebras are investigated, particularly, also over the quaternion skew field and the algebra of octonions. The transforms are considered in $\mathcal{A}_r$ spherical and $\mathcal{A}_r$ Cartesian coordinates. At the same time specific features of the noncommutative multiparameter transforms are elucidated, for example, related with the fact that in the Cayley-Dickson algebra $\mathcal{A}_r$ there are $2^r - 1$ imaginary generators $\{i_1, ..., i_{2^r-1}\}$ apart from one in the field of complex numbers such that the imaginary space in $\mathcal{A}_r$ has the dimension $2^r - 1$. Theorems about properties of images and originals in conjunction with the operations of linear combinations, differentiation, integration, shift and homothety are proved. An extension of the noncommutative multiparameter transforms for generalized functions is given. Formulas for noncommutative transforms of products and convolutions of functions are deduced.

Thus this solves the problem of non-commutative mathematical analysis to develop the multiparameter Laplace transform over the Cayley-Dickson algebras. Moreover, an application of the noncommutative integral transforms for solutions of partial differential equations is described. It can serve as an effective means (tool) to solve partial differential equations with real or complex coefficients with or without boundary conditions and their systems of different types. An algorithm is described which permits to write fundamental solutions and functions of Green's type. A moving boundary problem and partial differential equations with discontinuous coefficients are also studied with the use of the noncommutative transform.

Moreover, a decomposition theorem of linear partial differential operators over the Cayley-Dickson algebras is proved. A relation between fundamental solutions of an initial and component operators is demonstrated. In conjunction with a line integration over the Cayley-Dickson algebras and the decomposition theorem of partial differential operators it permits to solve partial differential equations linear with constant and variable coefficients and non-linear as well as boundary problems (see also [19]). Certainly, this approach effectively encompasses systems of partial differential equations, because each function f with values in the Cayley-Dickson algebra is the sum of functions $f_j i_j$, where each function f_j is real-valued.

All results of this paper are obtained for the first time.

2 Multidimensional noncommutative integral transforms.

1. Definitions. Transforms in $\mathcal{A}_r$ Cartesian coordinates.

Denote by $\mathcal{A}_r$ the Cayley-Dickson algebra, $0 \leq r$, which may be, in particular, $\mathbf{H} = \mathcal{A}_2$ the quaternion skew field or $\mathbf{O} = \mathcal{A}_3$ the octonion algebra. For unification of the notation we put $\mathcal{A}_0 = \mathbf{R}$, $\mathcal{A}_1 = \mathbf{C}$. A function $f : \mathbf{R}^n \to \mathcal{A}_r$ we call a function-original, where $2 \leq r$, $n \in \mathbf{N}$, if it fulfills the following conditions $(1 - 5)$.

(1). The function $f(t)$ is almost everywhere continuous on $\mathbf{R}^n$ relative to the Lebesgue measure λ_n on $\mathbf{R}^n$.

(2). On each finite interval in $\mathbf{R}$ each function $g_j(t_j) = f(t_1, ..., t_n)$ by t_j with marked all other variables may have only a finite number of points of discontinuity of the first kind, where $t = (t_1, ..., t_n) \in \mathbf{R}^n$, $t_j \in \mathbf{R}$, $j = 1, ..., n$. Recall that a point $u_0 \in \mathbf{R}$ is called a point of discontinuity of the first type, if there exist finite left and right limits $\lim_{u\to u_0, u<u_0} g(u) =: g(u_0 - 0) \in \mathcal{A}_r$ and $\lim_{u\to u_0, u>u_0} g(u) =: g(u_0 + 0) \in \mathcal{A}_r$.

(3). Every partial function $g_j(t_j)$=$f(t_1, ..., t_n)$ satisfies the Hölder condition: $|g_j(t_j + h_j) - g_j(t_j)| \le A_j|h_j|^{\alpha_j}$ for each $|h_j| < \delta$, where $0 < \alpha_j \le 1$, $A_j = const > 0$, δ_j>0 are constants for a given $t = (t_1, ..., t_n) \in \mathbf{R}^n$, $j = 1, ..., n$, everywhere on $\mathbf{R}^n$ may be besides points of discontinuity of the first type.

(4). The function $f(t)$ increases not faster, than the exponential function, that is there exist constants $C_v = const > 0$, $v = (v_1, ..., v_n)$, $a_{-1}, a_1 \in \mathbf{R}$, where $v_j \in \{-1, 1\}$ for every $j = 1, ..., n$, such that

$|f(t)| < C_v \exp((q_v, t))$ for each $t \in \mathbf{R}^n$ with $t_jv_j \ge 0$ for each $j = 1, ..., n$, $q_v = (v_1a_{v_1}, ..., v_na_{v_n})$; where

(5) $(x, y) := \sum_{j=1}^n x_jy_j$ denotes the standard scalar product in $\mathbf{R}^n$.

Certainly for a bounded original f it is possible to take $a_{-1} = a_1 = 0$.

Each Cayley-Dickson number $p \in \mathcal{A}_r$ we write in the form

(6) $p = \sum_{j=0}^{2^r-1} p_ji_j$, where $\{i_0, i_1, ..., i_{2^r-1}\}$ is the standard basis of generators of $\mathcal{A}_r$ so that $i_0 = 1$, $i_j^2 = -1$ and $i_0i_j = i_j = i_ji_0$ for each $j > 0$, $i_ji_k = -i_ki_j$ for each $j > 0$ and $k > 0$ with $k \ne j$, $p_j \in \mathbf{R}$ for each j.

If there exists an integral

(7) $F^n(p) := F^n(p; \zeta) := \int_{\mathbf{R}^n} f(t)e^{-<p,t)-\zeta}dt$,

then $F^n(p)$ is called the noncommutative multiparameter (Laplace) transform at a point $p \in \mathcal{A}_r$ of the function-original $f(t)$, where $\zeta - \zeta_0 = \zeta_1i_1 + ... + \zeta_{2^r-1}i_{2^r-1} \in \mathcal{A}_r$ is the parameter of an initial phase, $\zeta_j \in \mathbf{R}$ for each $j = 0, 1, ..., 2^r - 1$, $\zeta \in \mathcal{A}_r$, $n = 2^r - 1$, $dt = \lambda_n(dt)$,

(8) $< p, t) = p_0(t_1 + ... + t_{2^r-1}) + \sum_{j=1}^{2^r-1} p_jt_ji_j$, we also put

(8.1) $u(p, t; \zeta) =< p, t) + \zeta$.

For vectors $v, w \in \mathbf{R}^n$ we shall consider a partial ordering

(9) $v \prec w$ if and only if $v_j \le w_j$ for each $j = 1, ..., n$ and a k exists so that $v_k < w_k$, $1 \le k \le n$.

2. Transforms in $\mathcal{A}_r$ spherical coordinates.

Now we consider also the non-linear function $u = u(p, t; \zeta)$ taking into account non commutativity of the Cayley-Dickson algebra $\mathcal{A}_r$. Put

(1) $u(p, t) := u(p, t; \zeta) := p_0s_1 + M(p, t) + \zeta_0$, where

(2) $M(p, t) = M(p, t; \zeta) = (p_1s_1 + \zeta_1)\big[i_1 \cos(p_2s_2 + \zeta_2) + i_2 \sin(p_2s_2 + \zeta_2) \cos(p_3s_3 + \zeta_3) + ... + i_{2^r-2} \sin(p_2s_2 + \zeta_2)... \sin(p_{2^r-2}s_{2^r-2} + \zeta_{2^r-2}) \cos(p_{2^r-1}s_{2^r-1} + \zeta_{2^r-1}) + i_{2^r-1} \sin(p_2s_2 + \zeta_2)... \sin(p_{2^r-2}s_{2^r-2} + \zeta_{2^r-2}) \sin(p_{2^r-1}s_{2^r-1} + \zeta_{2^r-1})\big]$

for the general Cayley-Dickson algebra with $2 \le r < \infty$,

(2.1) $s_j := s_j(n; t) := t_j + ... + t_n$ for each $j = 1, ..., n$, $n = 2^r - 1$, so that $s_1 = t_1 + ... + t_n$, $s_n = t_n$. More generally, let

(3) $u(p, t) = u(p, t; \zeta) = p_0s_1 + w(p, t) + \zeta_0$, where $w(p, t)$ is a locally analytic function, $Re(w(p, t))$=0 for each $p \in \mathcal{A}_r$ and $t \in \mathbf{R}^{2^r-1}$, $Re(z) := (z+\tilde{z})/2$, $\tilde{z} = z^*$ denotes the conjugated number for $z \in \mathcal{A}_r$. Then the more general non-commutative multiparameter transform over

$\mathcal{A}_r$ is defined by the formula:

(4) $F^n_u(p;\zeta) := \int_{\mathbf{R}^n} f(t)\exp(-u(p,t;\zeta))dt$
for each Cayley-Dickson numbers $p \in \mathcal{A}_r$ whenever this integral exists as the principal value of either Riemann or Lebesgue integral, $n = 2^r - 1$. This non-commutative multiparameter transform is in $\mathcal{A}_r$ spherical coordinates, when $u(p,t;\zeta)$ is given by Formulas $(1,2)$.

At the same time the components p_j of the number p and ζ_j for ζ in $u(p,t;\zeta)$ we write in the p- and ζ-representations respectively such that

(5) $h_j = \left(-hi_j + i_j(2^r-2)^{-1}\left\{-h + \sum_{k=1}^{2^r-1} i_k(hi_k^*)\right\}\right)/2$ for each $j = 1,2,...,2^r-1$,

(6) $h_0 = \left(h + (2^r-2)^{-1}\left\{-h + \sum_{k=1}^{2^r-1} i_k(hi_k^*)\right\}\right)/2$,

where $2 \le r \in \mathbf{N}$, $h = h_0i_0 + ... + h_{2^r-1}i_{2^r-1} \in \mathcal{A}_r$, $h_j \in \mathbf{R}$ for each j, $i_k^* = \tilde{i}_k = -i_k$ for each $k > 0$, $i_0 = 1$, $h \in \mathcal{A}_r$. Denote $F^n_u(p;\zeta)$ in more details by $\mathcal{F}^n(f,u;p;\zeta)$.

Henceforth, the functions $u(p,t;\zeta)$ given by $1(8,8.1)$ or $(1,2,2.1)$ are used, if another form (3) is not specified. If for $u(p,t;\zeta)$ concrete formulas are not mentioned, it will be undermined, that the function $u(p,t;\zeta)$ is given in $\mathcal{A}_r$ spherical coordinates by Expressions $(1,2,2.1)$. If in Formulas 1(7) or (4) the integral is not by all, but only by $t_{j(1)},...,t_{j(k)}$ variables, where $1 \le k < n$, $1 \le j(1) < ... < j(k) \le n$, then we denote a noncommutative transform by $F^{k;t_{j(1)},...,t_{j(k)}}_u(p;\zeta)$ or $\mathcal{F}^{k;t_{j(1)},...,t_{j(k)}}(f,u;p;\zeta)$. If $j(1) = 1$,...,$j(k) = k$, then we denote it shortly by $F^k_u(p;\zeta)$ or $\mathcal{F}^k(f,u;p;\zeta)$. Henceforth, we take $\zeta_m = 0$ and $t_m = 0$ and $p_m = 0$ for each $1 \le m \notin \{j(1),...,j(k)\}$ if something other is not specified.

3. Remark. The spherical $\mathcal{A}_r$ coordinates appear naturally from the following consideration of iterated exponents:

(1) $\exp(i_1(p_1s_1+\zeta_1)\exp\big(-i_3(p_2s_2+\zeta_2)\exp(-i_1(p_3s_3+\zeta_3))\big))$
$= \exp\big(i_1(p_1s_1+\zeta_1)\exp\big(-(p_2s_2+\zeta_2)(i_3\cos(p_3s_3+\zeta_3) - i_2\sin(p_3s_3+\zeta_3))\big)\big)$
$= \exp\big(i_1(p_1s_1+\zeta_1)(\cos(p_2s_2+\zeta_2) - \sin(p_2s_2+\zeta_2)(i_3\cos(p_3s_3+\zeta_3) - i_2\sin(p_3s_3+\zeta_3)))\big)$
$= \exp\big((p_1s_1+\zeta_1)(i_1\cos(p_2s_2+\zeta_2)+i_2\sin(p_2s_2+\zeta_2)\cos(p_3s_3+\zeta_3)+i_3\sin(p_2s_2+\zeta_2)\sin(p_3s_3+\zeta_3))\big)$.

Consider i_{2^r} the generator of the doubling procedure of the Cayley-Dickson algebra $\mathcal{A}_{r+1}$ from the Cayley-Dickson algebra $\mathcal{A}_r$, such that $i_ji_{2^r} = i_{2^r+j}$ for each $j = 0,...,2^r-1$. We denote now the function $M(p,t;\zeta)$ from Definition 2 over $\mathcal{A}_r$ in more details by ${}_rM$.

Then by induction we write:

$$(2)\quad \exp\big({}_{r+1}M(p,t;\zeta)\big) = \exp\big\{{}_rM\big((i_1p_1 + ... + i_{2^r-1}p_{2^r-1}),(t_1,...,t_{2^r-2},(t_{2^r-1}+s_{2^r}));$$

$$(i_1\zeta_1 + ... + i_{2^r-1}\zeta_{2^r-1})\exp(-i_{2^r+1}(p_{2^r}s_{2^r}+\zeta_{2^r})\exp(-{}_rM((i_1p_{2^r+1} + ... + i_{2^r-1}p_{2^{r+1}-1}),$$

$$(t_{2^r+1},...,t_{2^{r+1}-1});(i_1\zeta_{2^r+1} + ... + i_{2^r-1}\zeta_{2^{r+1}-1}))))\big\},$$

where $t = (t_1,...,t_n)$, $n = n(r+1) = 2^{r+1}-1$, $s_j = s_j(n(r+1);t)$ for each $j = 1,...,n(r+1)$, since $s_m(n(r+1);t) = t_m + ... + t_{n(r+1)} = s_m(n(r);t) + s_{2^r}(n(r+1);t)$ for each $m = 1,...,2^r-1$.

An image function can be written in the form

(3) $F^n_u(p;\zeta) := \sum_{j=0}^{2^r-1} i_jF^n_{u,j}(p;\zeta)$,
where a function f is decomposed in the form

(3.1) $f(t) = \sum_{j=0}^{2^r-1} i_jf_j(t)$, $f_j : \mathbf{R}^n \to \mathbf{R}$ for each $j = 0,1,...,2^r-1$, $F^n_{u,j}(p;\zeta)$ denotes the image of the function-original f_j.

If an automorphism of the Cayley-Dickson algebra $\mathcal{A}_r$ is taken and instead of the standard generators $\{i_0,...,i_{2^r-1}\}$ new generators $\{N_0,...,N_{2^r-1}\}$ are used, this provides also $M(p,t;\zeta) =$

$M_N(p,t;\zeta)$ relative to new basic generators, where $2 \le r \in \mathbf{N}$. In this more general case we denote by ${}_NF^n_u(p;\zeta)$ an image for an original $f(t)$, or in more details we denote it by ${}_N\mathcal{F}^n(f,u;p;\zeta)$.

Formulas 1(7) and 2(4) define the right multiparameter transform. Symmetrically is defined a left multiparameter transform. They are related by conjugation and up to a sign of basic generators. For real valued originals they certainly coincide. Henceforward, only the right multiparameter transform is investigated.

Particularly, if $p = (p_0, p_1, 0, ..., 0)$ and $t = (t_1, 0, ..., 0)$, then the multiparameter non-commutative Laplace transforms 1(7) and 2(4) reduce to the complex case, with parameters a_1, a_{-1}. Thus, the given above definitions over quaternions, octonions and general Cayley-Dickson algebras are justified.

4. Theorem. *If an original $f(t)$ satisfies Conditions 1$(1-4)$ and $a_1<a_{-1}$, then its image $\mathcal{F}^n(f,u;p;\zeta)$ is $\mathcal{A}_r$-holomorphic (that is locally analytic) by p in the domain $\{z \in \mathcal{A}_r : a_1 < Re(z) < a_{-1}\}$, as well as by $\zeta \in \mathcal{A}_r$, where $1\le r\in\mathbf{N}$, $2^{r-1}\le n\le 2^r-1$, the function $u(p,t;\zeta)$ is given by 1(8,8.1) or 2(1,2, 2.1).*

Proof. At first consider the characteristic functions $\chi_{U_v}(t)$, where $\chi_U(t) = 1$ for each $t \in U$, while $\chi_U(t) = 0$ for every $t \in \mathbf{R}^n \setminus U$, $U_v := \{t \in \mathbf{R}^n : v_jt_j \ge 0 \ \forall j = 1, ..., n\}$ is the domain in the Euclidean space $\mathbf{R}^n$ for any v from §1. Therefore,

(1) $F^n_u(p;\zeta) := \sum_{[v=(v_1,...,v_n):v_1,...,v_n\in\{-1,1\}]} \int_{U_v} f(t)\exp(-u(p,t;\zeta))dt,$

since $\lambda_n(U_v \cap U_w) = 0$ for each $v \ne w$. Each integral $\int_{U_v} f(t)\exp(-u(p,t;\zeta))dt$ is absolutely convergent for each $p \in \mathcal{A}_r$ with the real part $a_1 < Re(p) < a_{-1}$, since it is majorized by the converging integral

(2) $\left|\int_{U_v} f(t)\exp(-u(p,t;\zeta))dt\right| \le \int_0^\infty ... \int_0^\infty C_v \exp\{-v_1(w-a_{v_1})y_1 - ... - v_n(w-a_{v_n})y_n - \zeta_0\}$ $dy_1...dy_n = C_ve^{-\zeta_0}\prod_{j=1}^n v_j(w-a_{v_j})^{-1}$,

where $w = Re(p)$, since $|e^z| = \exp(Re(z))$ for each $z \in \mathcal{A}_r$ in view of Corollary 3.3 [16]. While an integral, produced from the integral (1) differentiating by p converges also uniformly:

(3) $\left|\int_{U_v} f(t)\left[\partial\exp(-u(p,t;\zeta))/\partial p\right].hdt\right|$
$\le \int_0^\infty ... \int_0^\infty C_v\big|(h_0(v_1y_1 + ... + v_ny_n), h_1(v_1y_1 + ... + v_ny_n), ..., h_{n-1}(v_{n-1}y_{n-1} + v_ny_n), h_nv_ny_n)\big|$
$\exp\{-v_1(w-a_{v_1})y_1 - ... - v_n(w-a_{v_n})y_n - \zeta_0\}dy_1...dy_n$
$\le |h|C_ve^{-\zeta_0}\prod_{j=1}^n(w-a_{v_j})^{-2}$

for each $h \in \mathcal{A}_r$, since each $z \in \mathcal{A}_r$ can be written in the form $z = |z|\exp(M)$, where $|z|^2 = z\tilde{z} \in [0,\infty) \subset \mathbf{R}$, $M \in \mathcal{A}_r$, $Re(M) := (M + \tilde{M})/2 = 0$ in accordance with Proposition 3.2 [16]. In view of Equations 2(5, 6):

(4) $\partial\left(\int_{\mathbf{R}^n} f(t)\exp(-u(p,t;\zeta))dt\right)/\partial\tilde{p} = 0$ and

(5) $\partial\left(\int_{\mathbf{R}^n} f(t)\exp(-u(p,t;\zeta))dt\right)/\partial\tilde{\zeta} = 0$, while

(6) $\left|\int_{U_v} f(t)\left[\partial\exp(-u(p,t;\zeta))/\partial\zeta\right].hdt\right| \le |h|\int_0^\infty ... \int_0^\infty C_v\exp\big\{-v_1(w-a_{v_1})y_1 - ... - v_n(w-a_{v_n})y_n - \zeta_0\big\}dy_1...dy_n = |h|C_ve^{-\zeta_0}\prod_{j=1}^n v_j(w-a_{v_j})^{-1}$

for each $h \in \mathcal{A}_r$. In view of convergence of integrals given above (1−6) the multiparameter non-commutative transform $F^n_u(p;\zeta)$ is (super)differentiable by p and ζ, moreover, $\partial F^n_u(p;\zeta)/\partial\tilde{p} = 0$ and $\partial F^n_u(p;\zeta)/\partial\tilde{\zeta} = 0$ in the considered (p,ζ)-representation. In accordance with [17, 16] a function $g(p)$ is locally analytic by p in an open domain U in the Cayley-Dickson algebra $\mathcal{A}_r$, $2 \le r$, if and only if it is (super)differentiable by p, in another words $\mathcal{A}_r$-holomorphic. Thus, $F^n_u(p;\zeta)$ is $\mathcal{A}_r$-holomorphic by $p \in \mathcal{A}_r$ with $a_1 < Re(p) < a_{-1}$ and $\zeta \in \mathcal{A}_r$ due to Theorem 2.6 [18].

4.1. Corollary. *Let suppositions of Theorem 4 be satisfied. Then the image $\mathcal{F}^n(f,u;p;\zeta)$*

with $u = u(p,t;\zeta)$ given by 2(1,2) has the following periodicity properties:

(1) $\mathcal{F}^n(f,u;p;\zeta+\beta i_j) = \mathcal{F}^n(f,u;p;\zeta)$ *for each $j = 1,...,n$ and $\beta \in 2\pi\mathbf{Z}$;*

(2) $\mathcal{F}^n(f,u;p^1;\zeta^1) = (-1)^\kappa \mathcal{F}^n(f,u;p^2;\zeta^2)$ *for each $j = 1,...,n-1$ so that $\zeta_0^1 = \zeta_0^2$ and $\zeta_j^1 = -\zeta_j^2$, $\zeta_{j+1}^1 = \pi + \zeta_{j+1}^2$, $\zeta_s^1 = \zeta_s^2$ for each $s \ne j$ and $s \ne j+1$, while either $p_j^1 = -p_j^2$ and $p_l^1 = p_l^2$ for each $l \ne j$ with $\kappa = 2$ or $p^1 = p^2$ and $f(t)$ is an even function with $\kappa = 2$ by the $s_j = (t_j + ... + t_n)$ variable or an odd function by $s_j = (t_j + ... + t_n)$ with $\kappa = 1$;*

(3) $\mathcal{F}^n(f,u;p;\zeta+\pi i_1) = -\mathcal{F}^n(f,u;p;\zeta)$.

Proof. In accordance with Theorem 4 the image $\mathcal{F}^n(f,u;p;\zeta)$ exists for each $p \in W_f := \{z \in \mathcal{A}_r : a_1 < Re(z) < a_{-1}\}$ and $\zeta \in \mathcal{A}_r$, where $1 \le r$. Then from the 2π periodicity of sine and cosine functions the first statement follows. From $\sin(-\phi) = -\sin(\phi)$, $\cos(\phi) = \cos(-\phi)$, $\sin(\pi+\phi) = -\sin(\phi)$, $\cos(\phi+\pi) = -\cos(\phi)$ we get that $\cos(p_js_j + \zeta_j^1) = \cos(-p_js_j + \zeta_j^2)$, $\sin(p_js_j+\zeta_j^1)\cos(p_{j+1}s_{j+1}+\zeta_{j+1}^1) = (-\sin(-p_js_j+\zeta_j^2))(-\cos(p_{j+1}s_{j+1}+\zeta_{j+1}^2))$ and $\sin(p_js_j + \zeta_j^1)\sin(p_{j+1}s_{j+1}+\zeta_{j+1}^1) = (-\sin(-p_js_j+\zeta_j^2))(-\sin(p_{j+1}s_{j+1}+\zeta_{j+1}^2))$ for each $t \in \mathbf{R}^n$. On the other hand, either $p_j^1 = -p_j^2$ and $p_l^1 = p_l^2$ for each $l \ne j \ge 1$ with $\kappa = 2$ or $p^1 = p^2$ and $f(t_1,...,s_{j-1}+s_j, -s_j - s_{j+1}, t_{j+1},...,t_n) = (-1)^\kappa f(t_1,...,s_{j-1}-s_j, s_j - s_{j+1}, t_{j+1},...,t_n)$ is an even with $\kappa = 2$ or odd with $\kappa = 1$ function by the $s_j = (t_j + ... + t_n)$ variable for each $t = (t_1,...,t_n) \in \mathbf{R}^n$, where $t_j = s_j - s_{j+1}$ for $j = 1,...,n$, $s_{n+1} = s_{n+1}(n;t) = 0$. From this and Formulas 2(1,2,4) the second and the third statements of this corollary follow.

5. Remark. For a subset U in $\mathcal{A}_r$ we put $\pi_{\mathsf{s},\mathsf{p},\mathsf{t}}(U) := \{\mathsf{u} : z \in U, z = \sum_{\mathsf{v}\in\mathbf{b}} w_\mathsf{v}\mathsf{v},\ \mathsf{u} = w_\mathsf{s}\mathsf{s} + w_\mathsf{p}\mathsf{p}\}$ for each $\mathsf{s} \ne \mathsf{p} \in \mathbf{b}$, where $\mathsf{t} := \sum_{\mathsf{v}\in\mathbf{b}\setminus\{\mathsf{s},\mathsf{p}\}} w_\mathsf{v}\mathsf{v} \in \mathcal{A}_{r,\mathsf{s},\mathsf{p}} := \{z \in \mathcal{A}_r : z = \sum_{\mathsf{v}\in\mathbf{b}} w_\mathsf{v}\mathsf{v},\ w_\mathsf{s} = w_\mathsf{p} = 0,\ w_\mathsf{v} \in \mathbf{R}\ \forall \mathsf{v} \in \mathbf{b}\}$, where $\mathbf{b} := \{i_0, i_1, ..., i_{2^r-1}\}$ is the family of standard generators of the Cayley-Dickson algebra $\mathcal{A}_r$. That is, geometrically $\pi_{\mathsf{s},\mathsf{p},\mathsf{t}}(U)$ means the projection on the complex plane $\mathbf{C}_{\mathsf{s},\mathsf{p}}$ of the intersection U with the plane $\tilde{\pi}_{\mathsf{s},\mathsf{p},\mathsf{t}} \ni \mathsf{t}$, $\mathbf{C}_{\mathsf{s},\mathsf{p}} := \{a\mathsf{s} + b\mathsf{p} : a, b \in \mathbf{R}\}$, since $\mathsf{s}\mathsf{p}^* \in \hat{b} := \mathbf{b} \setminus \{1\}$. Recall that in §§2.5-7 [16] for each continuous function $f : U \to \mathcal{A}_r$ it was defined the operator $\hat{f}$ by each variable $z \in \mathcal{A}_r$. For the non-commutative integral transformations consider, for example, the left algorithm of calculations of integrals.

A Hausdorff topological space X is said to be n-connected for $n \ge 0$ if each continuous map $f : S^k \to X$ from the k-dimensional real unit sphere into X has a continuous extension over $\mathbf{R^{k+1}}$ for each $k \le n$ (see also [28]). A 1-connected space is also said to be simply connected.

It is supposed further, that a domain U in $\mathcal{A}_r$ has the property that U is (2^r-1)-connected; $\pi_{\mathsf{s},\mathsf{p},\mathsf{t}}(U)$ is simply connected in $\mathbf{C}$ for each $k = 0,1,...,2^{r-1}$, $\mathsf{s} = i_{2k}$, $\mathsf{p} = i_{2k+1}$, $\mathsf{t} \in \mathcal{A}_{r,\mathsf{s},\mathsf{p}}$ and $\mathsf{u} \in \mathbf{C}_{\mathsf{s},\mathsf{p}}$, for which there exists $z = \mathsf{u} + \mathsf{t} \in U$.

6. Theorem. *If a function $f(t)$ is an original (see Definition 1), such that ${}_NF^n_u(p;\zeta)$ is its image multiparameter non-commutative transform, where the functions f and F^n_u are written in the forms given by 3(3, 3.1), $f(\mathbf{R}^n) \subset \mathcal{A}_r$ over the Cayley-Dickson algebra $\mathcal{A}_r$, where $1 \le r \in \mathbf{N}$, $2^{r-1} \le n \le 2^r-1$.*

Then at each point t, where $f(t)$ satisfies the Hölder condition the equality is accomplished:

$$(1)\qquad f(t) = \left\{\left[(2\pi N_n)^{-1}\int_{-N_n\infty}^{N_n\infty}\right]\left(...\left(\left[(2\pi N_1)^{-1}\int_{-N_1\infty}^{N_1\infty}\right] {}_NF^n_u(a+p;\zeta)\right.\right.\right.$$

$$\left.\left.\left.\exp\{u(a+p,t;\zeta)\}\right)...\right)dp\right\} =: (\mathcal{F}^n)^{-1}({}_NF^n_u(a+p;\zeta),u,t;\zeta),$$

where either $u(p,t;\zeta) = <p,t) + \zeta$ or $u(p,t;\zeta) = p_0s_1 + M_N(p,t;\zeta) + \zeta_0$ (see §§1 and 2), the integrals are taken along the straight lines $p(\tau_j) = N_j\tau_j \in \mathcal{A}_r$, $\tau_j \in \mathbf{R}$ for each $j = 1,...,n$; $a_1 < Re(p) = a < a_{-1}$ and this integral is understood in the sense of the principal value, $t = (t_1,...,t_n) \in \mathbf{R}^n$, $dp = (...((d[p_1N_1])d[p_2N_2])...)d[p_nN_n]$.

Proof. In Integral (1) an integrand $\eta(p)dp$ certainly corresponds to the iterated integral as $(...(\eta(p)d[p_1N_1])...)d[p_nN_n]$, where $p = p_1N_1 + ... + p_nN_n$, $p_1, ..., p_n \in \mathbf{R}$. Using Decomposition 3(3.1) of a function f it is sufficient to consider the inverse transformation of the real valued function f_j, which we denote for simplicity by f. We put

$$ {}_NF^n_{u,j}(p;\zeta) := \int_{\mathbf{R}^n} f_j(t)\exp(-u(p,t;\zeta))dt. $$

If η is a holomorphic function of the Cayley-Dickson variable, then locally in a simply connected domain U in each ball $B(\mathcal{A}_r, z_0, R)$ with the center at z_0 of radius $R > 0$ contained in the interior $Int(U)$ of the domain U there is accomplished the equality

$\left(\partial\left[\int_{z_0}^{z}\eta(a+\zeta)d\zeta\right]/\partial z\right).1 = \eta(a+z),$

where the integral depends only on an initial z_0 and a final z points of a rectifiable path in $B(\mathcal{A}_r, z_0, R)$, $a \in \mathbf{R}$ (see also Theorem 2.14 [18]). Therefore, along the straight line $N_j\mathbf{R}$ the restriction of the antiderivative has the form $\int_{\theta_0}^{\theta}\eta(a+N_j\tau_j)d\tau_j$, since

(2) $\int_{z_0=N_j\theta_0}^{z=N_j\theta}\eta(a+\zeta)d\zeta = \int_{\theta_0}^{\theta}\hat{\eta}(a+N_j\tau_j).N_jd\tau_j,$

where $\partial\eta(a+z)/\partial\theta = (\partial\eta(a+z)/\partial z).N_j$ for the (super)differentiable by $z \in U$ function $\eta(z)$, when $z{=}\theta N_j$, $\theta \in \mathbf{R}$. For the chosen branch of the line integral specified by the left algorithm this antiderivative is unique up to a constant from $\mathcal{A}_r$ with the given z-representation ν of the function η [16, 17, 18]. On the other hand, for analytic functions with real expansion coefficients in their power series non-commutative integrals specified by left or right algorithms along straight lines coincide with usual Riemann integrals by the corresponding variables. The functions $\sin(z)$, $\cos(z)$ and e^z participating in the multiparameter non-commutative transform are analytic with real expansion coefficients in their series by powers of $z \in \mathcal{A}_r$.

Using Formula 4(1) we reduce the consideration to $\chi_{U_v}(t)f(t)$ instead of $f(t)$. By symmetry properties of such domains and integrals and utilizing change of variables it is sufficient to consider U_v with $v = (1, ..., 1)$. In this case $\int_{\mathbf{R}^n}$ for the direct multiparameter non-commutative transform 1(7) and 2(4) reduces to $\int_0^\infty ... \int_0^\infty$. Therefore, we consider in this proof below the domain $U_{1,...,1}$ only. Using Formulas $3(3, 3.1)$ and $2(1, 2, 2.1)$ we mention that any real algebra with generators $N_0 = 1$, N_k and N_j with $1 \le k \ne j$ is isomorphic with the quaternion skew field $\mathbf{H}$, since $Re(N_jN_k) = 0$ and $|N_j| = 1$, $|N_k| = 1$ and $|N_jN_k| = 1$. Then $\exp(\alpha + M\beta)\exp(\gamma + M\omega) = \exp((\alpha+\gamma) + M(\beta+\omega))$ for each real numbers $\alpha, \beta, \gamma, \delta$ and a purely imaginary Cayley-Dickson number M.

The octonion algebra $\mathbf{O}$ is alternative, while the real field $\mathbf{R}$ is the center of the Cayley-Dickson algebra $\mathcal{A}_r$. We consider the integral

(3) $g_b(t) := \left[(2\pi N_n)^{-1}\int_{-N_nb}^{N_nb}\right]\left(...\left(\left[(2\pi N_1)^{-1}\int_{-N_1b}^{N_1b}\right] {}_NF^n_{u,j}(a{+}p;\zeta)\exp\{u(a{+}p,t;\zeta)\}\right)...\right)dp$

for each positive value of the parameter $0 < b < \infty$. With the help of generators of the Cayley-Dickson algebra $\mathcal{A}_r$ and the Fubini Theorem for real valued components of the function the integral can be written in the form:

$$ (4) \qquad g_b(t) = \left[(2\pi N_n)^{-1}\int_0^\infty d\tau_n\int_{-N_nb}^{N_nb}\right]\left(...\left(\left[(2\pi N_1)^{-1}\int_0^\infty d\tau_1\int_{-N_1b}^{N_1b}\right]\right.\right. $$

$$ \left.\left. f(\tau)\exp\{-u_N(a+p,t;\zeta)\}\exp\{u_N(a+p,\tau;\zeta)\}\right)...\right)dp, $$

since the integral $\int_{U_{1,...,1}} f(\tau)\exp\{-u_N(a+p,\tau;\zeta)\}d\tau$ for any marked $0 < \delta < (a_{-1} - a_1)/3$ is uniformly converging relative to p in the domain $a_1 + \delta \le Re(p) \le a_{-1} - \delta$ in $\mathcal{A}_r$ (see also Proposition 2.18 [18]).

If take marked t_k for each $k \neq j$ and $S = N_j$ for some $j \geq 1$ in Lemma 2.17 [18] considering the variable t_j, then with a suitable ($\mathbf{R}$-linear) automorphism $\mathbf{v}$ of the Cayley-Dickson algebra $\mathcal{A}_r$ an expression for $\mathbf{v}(M(p,t;\zeta))$ simplifies like in the complex case with $\mathbf{C}_K := \mathbf{R} \oplus \mathbf{R}K$ for a purely imaginary Cayley-Dickson number K, $|K| = 1$, instead of $\mathbf{C} := \mathbf{R} \oplus \mathbf{R}i_1$, where $\mathbf{v}(x) = x$ for each real number $x \in \mathbf{R}$. But each equality $\alpha = \beta$ in $\mathcal{A}_r$ is equivalent to $\mathbf{v}(\alpha) = \mathbf{v}(\beta)$. Then

(5) $Re\,[(N_jN_q)(N_jN_l)^*] = Re(N_qN_l^*) = \delta_{q,l}$ for each q, l.

If $S^j = \sum_{0 \leq l \leq n; l \neq j} \alpha_l N_l$, $N^j = \sum_{0 \leq l \leq n; l \neq j} \beta_l N_l$ with $j \geq 1$ and real numbers $\alpha_l, \beta_l \in \mathbf{R}$ for each l, then

(6) $Re\,[(N_jS^j)(N_jN^j)^*] = Re\,[S^j(N^j)^*] = \sum_l \alpha_l\beta_l$.

The latter identity can be applied to either $S^k = M_{k+1}(p_{k+1}N_{k+1} + ... + p_nN_n, (t_{k+1}, ..., t_n); \zeta_{k+1}N_{k+1} + ... + \zeta_nN_n)$ and $N^k = M_{k+1}(p_{k+1}N_{k+1} + ... + p_nN_n, (\tau_{k+1}, ..., \tau_n); \zeta_{k+1}N_{k+1} + ... + \zeta_nN_n)$, or $S^k = (p_{k+1}t_{k+1} + \zeta_{k+1})N_{k+1} + ... + (p_nt_n + \zeta_n)N_n$ and $N^k = (p_{k+1}\tau_{k+1} + \zeta_{k+1})N_{k+1} + ... + (p_n\tau_n + \zeta_n)N_n$, where

(7) $M_{k+1}(p_{k+1}N_{k+1} + ... + p_nN_n, (t_{k+1}, ..., t_n); \zeta_{k+1}N_{k+1} + ... + \zeta_nN_n) = (p_{k+1}s_{1,k+1} + \zeta_{k+1})[N_{k+1}\cos(p_{k+2}s_{2,k+1} + \zeta_{k+2}) + ... + N_n \sin(p_{k+2}s_{2,k+1} + \zeta_{k+2}) ... \sin(p_ns_{n-k,k+1} + \zeta_n)]$,

(8) $s_{j,k+1} = s_{j,k+1}(n;t) = t_{k+j} + ... + t_n = s_{k+j}(n;t)$ for each $j = 1, ..., n-1$; $s_{n-k,k+1} = s_{n-k,k+1}(n;t) = t_n$.

We take the limit of $g_b(t)$ when b tends to the infinity. Evidently, $s_k(n;\tau) - s_j(n;\tau) = s_k(j-1;\tau) = \tau_k + ... + \tau_{j-1}$ for each $1 \leq k < j \leq n$. By our convention $s_k(n;\tau) = s_1(n;\tau)$ for $k < 1$, while $s_k(n;\tau) = 0$ for $k > n$. Put

(9) $u_{n,j}(p_0 + p_jN_j + ... + p_nN_n, (\tau_j, ..., \tau_n); \zeta_0 + \zeta_jN_j + ... + \zeta_nN_n) = \zeta_0 + p_0s_{1,j} + M_j(p_jN_j + ... + p_nN_n, (\tau_j, ..., \tau_n); \zeta_0 + \zeta_jN_j + ... + \zeta_nN_n)$
for u_N given by $2(1, 2, 2.1)$, where M_j is prescribed by (7), $s_{k,j} = s_{k,j}(n;\tau)$;

(10) $u_{n,j}(p_0 + p_jN_j + ... + p_nN_n, (\tau_j, ..., \tau_n); \zeta_0 + \zeta_jN_j + ... + \zeta_nN_n) = \zeta_0 + p_0s_{1,j} + \sum_{k=j}^n (p_k\tau_k + \zeta_k)N_k$
for $u = u_N$ given by $1(8, 8.1)$. For $j > 1$ the parameter ζ_0 for $u = u_N$ given by $1(8, 8.1)$ or $2(1, 2, 2.1)$ can be taken equal to zero.

When $t_1, ..., t_{j-1}, t_{j+1}, ..., t_n$ and $p_1, ..., p_{j-1}, p_{j+1}, ..., p_n$ variables are marked, we take the parameter
$\zeta^j := \zeta^j(p_jN_j + ... + p_nN_n, (\tau_j, ..., \tau_n); \zeta_0 + \zeta_jN_j + ... + \zeta_nN_n) := (\zeta_0 + \zeta_jN_j + ... + \zeta_nN_n) + (a + p_0)s_{j+1} + p_{j+1}s_{j+1}N_{j+1} + ... + p_ns_nN_n$ for $u(p,\tau;\zeta)$ given by Formulas $2(1, 2, 2.1)$ or
$\zeta^j := \zeta^j(p_jN_j + ... + p_nN_n), (\tau_j, ..., \tau_n); \zeta_0 + \zeta_jN_j + ... + \zeta_nN_n) := (\zeta_0 + \zeta_jN_j + ... + \zeta_nN_n) + (a + p_0)s_{j+1} + p_{j+1}\tau_{j+1}N_{j+1} + ... + p_n\tau_nN_n$ for $u(p,\tau;\zeta)$ described in $1(8, 8.1)$. Then the integral operator
$\lim_{b\to\infty}[(2\pi N_j)^{-1}\int_0^\infty d\tau_j \int_{-N_jb}^{N_jb}] ... (dp_jN_j)$ (see also Formula (4) above) applied to the function $f(t_1, ..., t_{j-1}, \tau_j, ..., \tau_n)\exp\{-u_{N,j}(a + p_0 + p_jN_j + ... + p_nN_n, (t_j, ..., t_n); \zeta_0 + \zeta_jN_j + ... + \zeta_nN_n)\}\exp\{u_{N,j}(a + p_0 + p_jN_j + ... + p_nN_n, (\tau_j, ..., \tau_n); \zeta_0 + \zeta_jN_j + ... + \zeta_nN_n)\}$ with the parameter ζ^j instead of ζ treated by Theorems 2.19 and 3.15 [18] gives the inversion formula corresponding to the real variable t_j for $f(t)$ and to the Cayley-Dickson variable $p_0N_0 + p_jN_j$ restricted on the complex plane $\mathbf{C}_{N_j} = \mathbf{R} \oplus \mathbf{R}N_j$, since $d(\tau_j + c) = d\tau_j$ for each (real) constant c. After integrations with $j = 1, ..., k$ with the help of Formulas $(6 - 10)$ and $3(1, 2)$ we get the following:

$$(11) \quad \lim_{b\to\infty} g_b(t) = Re\left[(2\pi N_n)^{-1}\int_0^\infty d\tau_n \int_{-N_n\infty}^{N_n\infty}\right]\left(...\left(\left[(2\pi N_{k+1})^{-1}\int_0^\infty d\tau_{k+1}\int_{-N_{k+1}\infty}^{N_{k+1}\infty}\right]\right.\right.$$

$$f(t_1, ..., t_k, \tau_{k+1}, ..., \tau_n)\exp\{-u_{N,k+1}((a + p_0 + p_{k+1}N_{k+1} + ... + p_nN_n), (t_{k+1}, ..., t_n);$$

$$(\zeta_0+\zeta_{k+1}N_{k+1}+...+\zeta_nN_n))\}\exp\{u_{N,k+1}((a+p_0+p_{k+1}N_{k+1}+...+p_nN_n),$$

$$(\tau_{k+1},...,\tau_n);(\zeta_0+\zeta_{k+1}N_{k+1}+...+\zeta_nN_n))\}\Big)...\Big)dp.$$

Moreover, $Re(f_q)=f_q$ for each q and in (11) the function $f=f_q$ stands for some marked q in accordance with Decompositions 3(3, 3.1) and the beginning of this proof.

Mention, that the algebra $alg_{\mathbf{R}}(N_j,N_k,N_l)$ over the real field with three generators N_j, N_k and N_l is alternative. The product N_kN_l of two generators is also the corresponding generator $(-1)^{\xi(k,l)}N_m$ with the definite number $m=m(k,l)$ and the sign multiplier $(-1)^{\xi(k,l)}$, where $\xi(k,l)\in\{0,1\}$. On the other hand, $N_{k_1}[\tilde{N}_j(N_j(N_{k_2}N_l))]=N_{k_1}(N_{k_2}N_l)$. We use decompositions $(7-10)$ and take $k_2=l$ due to Formula (11), where Re stands on the right side of the equality, since $Re(N_kN_l)=0$ and $Re[\tilde{N}_j(N_j(N_kN_l))]=0$ for each $k\ne l$. Thus the repeated application of this procedure by $j=1,2,...,n$ leads to Formula (1) of this theorem.

6.1. Corollary. *If the conditions of Theorem 6 are satisfied, then*

$$(1)\qquad f(t)=(2\pi)^{-n}\int_{\mathbf{R}^n}F^n_u(a+p;\zeta)\exp\{u(a+p,t;\zeta)\}dp_1...dp_n$$

$$=(\mathcal{F}^n)^{-1}\big({}_NF^n_u(a+p;\zeta),u,t;\zeta\big).$$

Proof. Each algebra $alg_{\mathbf{R}}(N_j,N_k,N_l)$ is alternative. Therefore, in accordance with §6 and Formulas 1(8, 8.1) and 2(1 − 4) for each non-commutative integral given by the left algorithm we get

$$(2)\qquad N_j^{-1}\int_{-N_jb}^{N_jb}[f(\tau)\exp\{-u_N(a+p,t;\zeta)\}]\exp\{u_N(a+p,\tau;\zeta)\}d(p_jN_j)$$

$$\sum_{l=0}^{2^r-1}\tilde{N}_j\left[N_j\left(\int_{-N_jb}^{N_jb}[N_lf_l(\tau)\exp\{-u_N(a+p,t;\zeta)\}]\exp\{u_N(a+p,\tau;\zeta)\}dp_j\right)\right]$$

$$=\int_{-b}^{b}[f(\tau)\exp\{-u_N(a+p,t;\zeta)\}]\exp\{u_N(a+p,\tau;\zeta)\}dp_j$$

for each $j=1,...,n$, since the real field is the center of the Cayley-Dickson algebra $\mathcal{A}_r$, while the functions sin and cos are analytic with real expansion coefficients. Thus

(3) $g_b(t)=(2\pi)^{-n}\Big[\int_0^\infty d\tau_n\int_{-b}^b\Big]\Big(...\Big(\Big[\int_0^\infty d\tau_1\int_{-b}^b\Big]f(\tau)\exp\{-u_N(a+p,t;\zeta)\}$
$\exp\{u_N(a+p,\tau;\zeta)\}\Big)...\Big)dp_1...dp_n,$

hence taking the limit with b tending to the infinity implies, that the non-commutative iterated (multiple) integral in Formula 6(1) reduces to the principal value of the usual integral by real variables $(\tau_1,...,\tau_n)$ and $(p_1,...,p_n)$ 6.1(1).

7. Theorem. *An original $f(t)$ with $f(\mathbf{R}^n)\subset\mathcal{A}_r$ over the Cayley-Dickson algebra $\mathcal{A}_r$ with $1\le r\in\mathbf{N}$ is completely defined by its image ${}_NF^n_u(p;\zeta)$ up to values at points of discontinuity, where the function $u(p,t;\zeta)$ is given by 1(8,8.1) or 2(1,2, 2.1).*

Proof. Due to Corollary 6.1 the value $f(t)$ at each point t of continuity of $f(t)$ has the expression throughout ${}_NF^n_u(p;\zeta)$ prescribed by Formula 6.1(1). Moreover, values of the original at points of discontinuity do not influence on the image ${}_NF^n_u(p;\zeta)$, since on each bounded interval in $\mathbf{R}$ by each variable t_j a number of points of discontinuity is finite and by our supposition above the original function $f(t)$ is λ_n-almost everywhere on $\mathbf{R}^n$ continuous.

8. Theorem. *Suppose that a function ${}_NF^n_u(p;\zeta)$ is analytic by the variable $p\in\mathcal{A}_r$ in a domain*

$W := \{p \in \mathcal{A}_r : a_1 < Re(p) < a_{-1}\}$, where $2 \le r \in \mathbf{N}$, $2^{r-1} \le n \le 2^r - 1$, $f(\mathbf{R}^n) \subset \mathcal{A}_r$, either $u(p,t;\zeta) =< p,t) + \zeta$ or $u(p,t;\zeta) := p_0 s_1 + M(p,t;\zeta) + \zeta_0$ (see §§1 and 2). Let ${}_NF^n_u(p;\zeta)$ be written in the form ${}_NF^n_u(p;\zeta) = {}_NF^{n,0}_u(p;\zeta) + {}_NF^{n,1}_u(p;\zeta)$, where ${}_NF^{n,0}_u(p;\zeta)$ is holomorphic by p in the domain $a_1 < Re(p)$. Let also ${}_NF^{n,1}_u(p;\zeta)$ be holomorphic by p in the domain $Re(p) < a_{-1}$. Moreover, for each $a > a_1$ and $b < a_{-1}$ there exist constants $C_a > 0$, $C_b > 0$ and $\epsilon_a > 0$ and $\epsilon_b > 0$ such that

(1) $\left| {}_NF^{n,0}_u(p;\zeta)\right| \le C_a \exp(-\epsilon_a |p|)$ *for each $p \in \mathcal{A}_r$ with $Re(p) \ge a$,*

(2) $\left| {}_NF^{n,1}_u(p;\zeta)\right| \le C_b \exp(-\epsilon_b |p|)$ *for each $p \in \mathcal{A}_r$ with $Re(p) \le b$, the integral*

(3) $\int_{-N_n\infty}^{N_n\infty} ... \int_{-N_1\infty}^{N_1\infty} {}_NF^{n,k}_u(w+p;\zeta)dp$ *converges absolutely for $k=0$ and $k=1$ and each $a_1<w<a_{-1}$.*

Then ${}_NF^n_u(w+p;\zeta)$ is the image of the function

(4)
$$f(t) = \left[(2\pi)^{-1}\tilde{N}_n \int_{-N_n\infty}^{N_n\infty}\right]\left(...\left(\left[(2\pi)^{-1}\tilde{N}_1 \int_{-N_1\infty}^{N_1\infty}\right] {}_NF^n_u(w+p;\zeta)\exp\{u(w+p,t;\zeta)\}\right)...\right)dp$$

$$= (\mathcal{F}^n)^{-1}({}_NF^n_u(w+p;\zeta), u, t;\zeta).$$

Proof. For the function ${}_NF^{n,1}_u(p;\zeta)$ we consider the substitution of the variable $p = -g$, $-a_{-1} < Re(g)$. Thus the proof reduces to the consideration of ${}_NF^{n,0}_u(w+p;\zeta)$.

An integration by dp in the iterated integral (4) is treated as in §6. Take marked values of variables $p_1,...,p_{j-1},p_{j+1},...,p_n$ and $t_1,...,t_{j-1},t_{j+1},...,t_n$, where $s_k = s_k(n;\tau)$ for each $k = 1,...,n$ (see §6 also). For a given parameter $\zeta^j := (\zeta_0 + \zeta_j N_j + ... + \zeta_n N_n) + (w+p_0)s_{j+1} + p_{j+1}s_{j+1}N_{j+1} + ... + p_n s_n N_n$ for $u(p,\tau;\zeta)$ prescribed by Formulas 2(1, 2, 2.1) or $\zeta^j := (\zeta_0 + \zeta_j N_j + ... + \zeta_n N_n) + (w+p_0)s_{j+1} + p_{j+1}\tau_{j+1}N_{j+1} + ... + p_n\tau_n N_n$ for $u(p,t;\zeta)$ given by 1(8, 8.1) instead of ζ and any non-zero Cayley-Dickson number $\beta \in \mathcal{A}_r$ we have $\lim_{\tau_j\to\infty}[\beta\tau_j + \zeta^j]/[\beta\tau_j + \zeta] = 1$.

For any locally z-analytic function $g(z)$ in a domain U satisfying conditions of §5 the homotopy theorem for a non-commutative line integral over $\mathcal{A}_r$, $2 \le r$, is satisfied (see [17, 16]). In particular if U contains the straight line $w + \mathbf{R}N_j$ and the path $\gamma_j(t_j) := \zeta^j + t_j N_j$, then $\int_{-N_j\infty}^{N_j\infty} g(z)dz = \int_{\gamma_j} g(w+z)dz$, when $\hat{g}(z) \to 0$ while $|z|$ tends to the infinity, since $|\zeta^j|$ is a finite number (see Lemma 2.23 in [18]). We apply this to the integrand in Formula (4), since ${}_NF^n_u(w+p;\zeta)$ is locally analytic by p in accordance with Theorem 4 and Conditions (1, 2) are satisfied.

Then the integral operator $\left[(2\pi N_j)^{-1}\int_{-N_j\infty}^{N_j\infty}\right]$ on the j-th step with the help of Theorems 2.22 and 3.16 [18] gives the inversion formula corresponding to the real parameter t_j for $f(t)$ and to the Cayley-Dickson variable $p_0N_0 + p_jN_j$ which is restricted on the complex plane $\mathbf{C}_{N_j} = \mathbf{R} \oplus \mathbf{R}N_j$ (see also Formulas 6(4, 11) above). Therefore, an application of this procedure by $j = 1, 2, ..., n$ as in §6 implies Formula (4) of this theorem.

Thus there exist originals f^0 and f^1 for functions ${}_NF^{n,0}_u(p;\zeta)$ and ${}_NF^{n,1}_u(p;\zeta)$ with a choice of $w \in \mathbf{R}$ in the common domain $a_1 < Re(p) < a_{-1}$. Then $f = f^0 + f^1$ is the original for ${}_NF^n_u(p;\zeta)$ due to the distributivity of the multiplication in the Cayley-Dickson algebra $\mathcal{A}_r$ leading to the additivity of the considered integral operator in Formula (4).

8.1. Corollary. *Let the conditions of Theorem 8 be satisfied, then*

(1)
$$f(t) = (2\pi)^{-n}\int_{\mathbf{R}^n} {}_NF^n_u(w+p;\zeta)\exp\{u(w+p,t;\zeta)\}dp_1...dp_n$$

$$= (\mathcal{F}^n)^{-1}({}_NF^n_u(w+p;\zeta), u, t;\zeta).$$

Proof. In accordance with §§6 and 6.1 each non-commutative integral given by the left algorithm reduces to the principal value of the usual integral by the corresponding real variable:

$$(2)\qquad (2\pi)^{-1}\tilde{N}_j\int_{-N_j\infty}^{N_j\infty} {}_NF^n_u(w+p;\zeta)\exp\{u(w+p,t;\zeta)\}d(p_jN_j)$$

$$=(2\pi)^{-1}\int_{-\infty}^{\infty} {}_NF^n_u(w+p;\zeta)\exp\{u(w+p,t;\zeta)\}dp_j$$

for each $j=1,...,n$. Thus Formula 8(4) with the non-commutative iterated (multiple) integral reduces to Formula 8.1(1) with the principal value of the usual integral by real variables $(p_1,...,p_n)$.

9. Note. In Theorem 8 Conditions $(1,2)$ can be replaced on

(1) $\lim_{n\to\infty}\sup_{p\in C_{R(n)}}\|\hat{F}(p)\|=0$,

where $C_{R(n)}:=\{z\in\mathcal{A}_r: |z|=R(n), a_1<Re(z)<a_{-1}\}$ is a sequence of intersections of spheres with a domain W, where $R(n)<R(n+1)$ for each n, $\lim_{n\to\infty}R(n)=\infty$. Indeed, this condition leads to the accomplishment of the $\mathcal{A}_r$ analog of the Jordan Lemma for each $r\ge 2$ (see also Lemma 2.23 and Remark 2.24 [18]).

Subsequent properties of quaternion, octonion and general $\mathcal{A}_r$ multiparameter non-commutative analogs of the Laplace transform are considered below. We denote by

(2) $W_f=\{p\in\mathcal{A}_r:\ a_1(f)<Re(p)<a_{-1}(f)\}$ a domain of ${}_NF^n_u(p;\zeta)$ by the p variable, where $a_1=a_1(f)$ and $a_{-1}=a_{-1}(f)$ are as in §1. For an original

(3) $f(t)\chi_{U_{1,...,1}}(t)$ we put $W_f=\{p\in\mathcal{A}_r:\ a_1(f)<Re(p)\}$,

that is $a_{-1}=\infty$. Cases may be, when either the left hyperplane $Re(p)=a_1$ or the right hyperplane $Re(p)=a_{-1}$ is (or both are) included in W_f. It may also happen that a domain reduces to the hyperplane $W_f=\{p:\ Re(p)=a_1=a_{-1}\}$.

10. Proposition. *If images ${}_NF^n_u(p;\zeta)$ and ${}_NG^n_u(p;\zeta)$ of functions-originals $f(t)$ and $g(t)$ exist in domains W_f and W_g with values in $\mathcal{A}_r$, where the function $u(p,t;\zeta)$ is given by 1$(8,8.1)$ or 2$(1,2,2.1)$, then for each $\alpha,\beta\in\mathcal{A}_r$ in the case $\mathcal{A}_2=\mathbf{H}$; as well as f and g with values in $\mathbf{R}$ and each $\alpha,\beta\in\mathcal{A}_r$ or f and g with values in $\mathcal{A}_r$ and each $\alpha,\beta\in\mathbf{R}$ in the case of $\mathcal{A}_r$ with $r\ge 3$; the function $\alpha\ {}_NF_u(p;\zeta)+\beta\ {}_NG_u(p;\zeta)$ is the image of the function $\alpha f(t)+\beta g(t)$ in a domain $W_f\cap W_g$.*

Proof. Since the transforms ${}_NF^n_u(p;\zeta)$ and ${}_NG^n_u(p;\zeta)$ exist, then the integral

$$\int_{\mathbf{R}^n}(\alpha f(t)+\beta g(t))\exp(-u(p,t;\zeta))dt=\int_{\mathbf{R}^n}\alpha f(t)\exp(-u(p,t;\zeta))dt$$

$$+\int_{\mathbf{R}^n}\beta g(t)\exp(-u(p,t;\zeta))dt$$

converges in the domain

$W_f\cap W_g=\{p\in\mathcal{A}_r:\ \max(a_1(f),a_1(g))<Re(p)<\min(a_{-1}(f),a_{-1}(g))\}$.

We have $t\in\mathbf{R}^n$, $2^{r-1}\le n\le 2^r-1$, while $\mathbf{R}$ is the center of the Cayley-Dickson algebra $\mathcal{A}_r$. The quaternion skew field $\mathbf{H}$ is associative. Thus, under the imposed conditions the constants α,β can be carried out outside integrals.

11. Theorem. *Let $\alpha=const>0$, let also $F^n(p;\zeta)$ be an image of an original function $f(t)$ with either $u=<p,t)+\zeta$ or u given by Formulas 2$(1,2)$ over the Cayley-Dickson algebra $\mathcal{A}_r$ with $2\le r<\infty$, $2^{r-1}\le n\le 2^r-1$. Then an image $F^n(p/\alpha;\zeta)/\alpha^n$ of the function $f(\alpha t)$ exists.*

Proof. Since $p_js_j+\zeta_j = p_j(s'_j/\alpha)+\zeta_j = (p_j/\alpha)s'_j+\zeta_j$ for each $j=1,...,n$, where $s_j\alpha = s'_j$, $s_j = s_j(n;t)$, $s'_j = s_j(n;\tau)$, $\tau_j = \alpha t_j$ for each $j=1,...,n$. Then changing of these variables implies:

$\int_{\mathbf{R}^n} f(\alpha t)e^{-u(p,t;\zeta)}dt = \int_{\mathbf{R}^n} f(\tau)e^{-u(p,\tau/\alpha;\zeta)}d\tau/\alpha^n = F^n(p/\alpha;\zeta)/\alpha^n$

due to the fact that the real filed $\mathbf{R}$ is the center $Z(\mathcal{A}_r)$ of the Cayley-Dickson algebra $\mathcal{A}_r$.

12. Theorem. *Let $f(t)$ be a function-original on the domain $U_{1,...,1}$ such that $\partial f(t)/\partial t_k$ also for $k=j-1$ and $k=j$ satisfies Conditions 1(1 − 4). Suppose that $u(p,t;\zeta)$ is given by 2(1, 2, 2.1) or 1(8, 8.1) over the Cayley-Dickson algebra $\mathcal{A}_r$ with $2\le r<\infty$, $2^{r-1}\le n\le 2^r-1$. Then*

$$(1)\qquad \mathcal{F}^n\Big((\partial f(t)/\partial t_j)\chi_{U_{1,...,1}}(t),u;p;\zeta\Big) = -\mathcal{F}^{n-1;t^j}\Big(f(t)\chi_{U_{1,...,1}}(t^j),u(p,t^j;\zeta);p;\zeta\Big)$$

$$+\left[p_0+\sum_{k=1}^{j}p_k\mathsf{S}_{e_k}\right]\mathcal{F}^n\Big(f(t)\chi_{U_{1,...,1}}(t),u;p;\zeta\Big)$$

in the $\mathcal{A}_r$ spherical coordinates or

$$(1.1)\qquad \mathcal{F}^n\Big((\partial f(t)/\partial t_j)\chi_{U_{1,...,1}}(t),u;p;\zeta\Big) = -\mathcal{F}^{n-1;t^j}\Big(f(t)\chi_{U_{1,...,1}}(t^j),u(p,t^j;\zeta);p;\zeta\Big)$$

$$+\big[p_0+p_j\mathsf{S}_{e_j}\big]\mathcal{F}^n\Big(f(t)\chi_{U_{1,...,1}}(t),u;p;\zeta\Big)$$

in the $\mathcal{A}_r$ Cartesian coordinates in a domain $W=\{p\in\mathcal{A}_r:\ \max(a_1(f),a_1(\partial f/\partial t_j))<Re(p)\}$, where $t^j:=(t_1,...,t_j,...,t_n:\ t_j=0)$, $\mathsf{S}_{e_k}=-\partial/\partial\zeta_k$ for each $k\ge 1$.

Proof. Certainly,

(2) $\partial f\big(t(s)\big)/\partial s_1 = \partial f(t)/\partial t_1$ and

(2.1) $\partial f(t)/\partial t_j = \sum_{k=1}^n\big(\partial f(t(s))/\partial s_k\big)(\partial s_k/\partial t_j) = \sum_{k=1}^j \partial f(t(s))/\partial s_k$

for each $j=2,...,n$, since $t_j = s_j - s_{j+1}$, $t_1 = s_1 - s_2$, where $s_j = s_j(n;t)$, $s_{n+l}=0$ for each $l\ge 1$. From Formulas 30(6, 7) [18] we have the equality in the $\mathcal{A}_r$ spherical coordinates:

(3) $\partial\exp(-u(p,t;\zeta))/\partial s_j = -p_0\delta_{1,j}\exp(-u(p,t;\zeta)) - p_j\mathsf{S}_{e_j}\exp(-u(p,t;\zeta))$,

since

$\exp(-u(p,t;\zeta)) = \exp\{-p_0s_1-\zeta_0\}\exp(-M(p,t;\zeta))$,

$\partial\exp(-p_0s_1-\zeta_0)/\partial s_j = -p_0\delta_{1,j}\exp(-p_0s_1-\zeta_0)$,

$\partial[\cos(p_js_j+\zeta_j)-\sin(p_js_j+\zeta_j)i_j]/\partial s_j = \partial\exp(-(p_js_j+\zeta_j)i_j)/\partial s_j = -p_ji_j\exp(-(p_js_j+\zeta_j)i_j)$

$= -p_j\exp(-(p_js_j+\zeta_j-\pi/2)i_j) = -p_j[\cos(p_js_j+\zeta_j-\pi/2)-\sin(p_js_j+\zeta_j-\pi/2)i_j]$

$= -p_j\mathsf{S}_{e_j}[\cos(p_js_j+\zeta_j)-\sin(p_js_j+\zeta_j)i_j]$,

since s_j and s_k are real independent variables for each $k\ne j$, where $\delta_{j,k}=0$ for $j\ne k$, while $\delta_{j,j}=1$,

(3.1) $\mathsf{S}_{e_j}[\cos(p_js_j+\zeta_j)-\sin(p_js_j+\zeta_j)i_j] =$

$-\partial[\cos(p_js_j+\zeta_j)-\sin(p_js_j+\zeta_j)i_j]/\partial\zeta_j$

$= [\cos(p_js_j+\zeta_j-\pi/2)-\sin(p_js_j+\zeta_j-\pi/2)i_j]$.

In the $\mathcal{A}_r$ Cartesian coordinates we take t_j instead of s_j in (3.1). If $\phi(z)$ is a differentiable function by z_j for each j, $\phi:\mathcal{A}_r\to\mathcal{A}_r$, $z_j = p_jt_j+\zeta_j$, then

(3.2) $\partial\exp(-\phi(z))/\partial(qt_j) = -q[d\exp(\xi)/d\xi]|_{\xi=-\phi}.(\partial\phi(z)/\partial z_j)p_j$

$= -qp_j[\sum_{n=1}^\infty\sum_{k=1}^{n-1}((\xi(z))^k(\partial\phi(z)/\partial z_j))(\xi(z))^{n-1-k}/n!]|_{\xi=-\phi}$

$= -qp_j(-\partial\exp(-\phi(z))/\partial\zeta_j) = -\,p_j\mathsf{S}_{qe_j}\exp(-\phi(z))$,

where either $q=1$ or $q=-1$, since $\partial z_j/\partial\zeta_j=1$.

That is

(3.3) $\mathsf{S}^x_{e_j} \exp(-i_k(\phi_k+\zeta_k)) = 0$ for each $j \ne k \ge 1$ and any positive number $x > 0$,

(3.4) $\mathsf{S}^x_{e_j} \exp(-i_j(\phi_j+\zeta_j)) = \exp(-i_j(\phi_j+\zeta_j - x\pi/2))$ and
$\mathsf{S}^x_{-e_j} \exp(-i_j(\phi_j+\zeta_j)) = \exp(-i_j(\phi_j+\zeta_j + x\pi/2))$

for each non-negative real number $x \ge 0$, ϕ_k and $\zeta_k \in \mathbf{R}$, where $\mathsf{S}_{e_j} = \mathsf{S}_{e_j}(\zeta_j)$, the zero power $\mathsf{S}^0_{e_j} = I$ is the unit operator;

(3.5) $\mathsf{S}_{qe_j} e^{-u(p,t;\zeta)} = e^{-p_0 s_1 - \zeta_0}$

$T^q_j\Big[i_0\delta_{j,1}\cos(p_1s_1+\zeta_1) + (1-\delta_{j,1}) i_{j-1}\sin(p_1s_1+\zeta_1)...\cos(p_js_j+\zeta_j) + \Big\{\sum_{k=j}^{2^r-2} i_k \sin(p_1s_1+\zeta_1)...\cos(p_{k+1}s_{k+1}+\zeta_{k+1})\Big\} + i_{2^r-1}\sin(p_1s_1+\zeta_1)...\sin(p_{2^r-1}s_{2^r-1}+\zeta_{2^r-1})\Big]$

in the $\mathcal{A}_r$ spherical coordinates, where either $q=1$ or $q=-1$ and

(3.6) $T^x_j\xi(\zeta_j) := \xi(\zeta_j - x\pi/2)$

for any function $\xi(\zeta_j)$ and any real number $x \in \mathbf{R}$, where $j \ge 1$. Then in accordance with Formula (3.2) we have:

(3.7) $\mathsf{S}_{qe_j}\exp(-u(p,t;\zeta)) =$
$= \Big[\sum_{n=1}^{\infty}\sum_{k=1}^{n-1}\big((\xi(z))^k q i_j\big)(\xi(z))^{n-1-k}/n!\Big]\Big|_{\xi=-u(p,t;\zeta)}$

for $u(p,t;\zeta)$ given by Formulas 1(8, 8.1) in the $\mathcal{A}_r$ Cartesian coordinates, where either $q=1$ or $q=-1$.

The integration by parts theorem (Theorem 2 in §II.2.6 on p. 228 [10]) states: if $a<b$ and two functions f and g are Riemann integrable on the segment $[a,b]$, $F(x) = A + \int_a^x f(t)dt$ and $G(x) = B + \int_a^x g(t)dt$, where A and B are two real constants, then $\int_a^b F(x)g(x)dx = F(x)G(x)\big|_a^b - \int_a^b f(x)G(x)dx$.

Therefore, the integration by parts gives

$$(4)\quad \int_0^\infty (\partial f(t)/\partial t_j)\exp(-u(p,t;\zeta))dt_j = f(t)\exp(-u(p,t;\zeta))\Big|_{t_j=0}^{t_j=\infty}$$
$$-\int_0^\infty \big[f(t)(\partial\exp(-u(p,t;\zeta))/\partial t_j)\big]dt_j.$$

Using the change of variables $t \mapsto s$ with the unit Jacobian $\partial(t_1,...,t_n)/\partial(s_1,...,s_n)$ and applying the Fubini's theorem componentwise to $f_j i_j$ we infer:

$$(5)\quad \int_{U_{1,...,1}} (\partial f(t)/\partial t_j)\exp(-u(p,t;\zeta))dt = \int_{s_1\ge s_2\ge...\ge s_n\ge 0} (\partial f(t)/\partial t_j)\exp(-u(p,t;\zeta))ds$$
$$= \int_0^\infty ... \int_0^\infty \left[\int_{s_{j+1}}^\infty (\partial f(t)/\partial t_j)\exp(-u(p,t;\zeta))ds_j\right] dt^j$$
$$= -\left[\int_0^\infty ... \int_0^\infty f(t^j)\exp(-u(p,t^j;\zeta))dt^j\right]$$
$$+\left[p_0 + \sum_{k=1}^j p_k \mathsf{S}_{e_k}\right]\int_0^\infty ... \int_0^\infty f(t)\exp(-u(p,t;\zeta))dt$$

in the $\mathcal{A}_r$ spherical coordinates, or

$$(5.1)\quad \int_{U_{1,...,1}} (\partial f(t)/\partial t_j)\exp(-u(p,t;\zeta))dt$$

$$= -\left[\int_0^\infty ... \int_0^\infty f(t^j) \exp(-u(p,t^j;\zeta)) dt^j\right]$$

$$+\left[p_0 + p_j \mathsf{S}_{e_j}\right] \int_0^\infty ... \int_0^\infty f(t) \exp(-u(p,t;\zeta)) dt$$

in the $\mathcal{A}_r$ Cartesian coordinates, since $\partial \exp(-(p_0 s_1 + \zeta_0))/\partial t_j = -p_0 \exp(-(p_0 s_1 + \zeta_0))$ for each $1 \le j \le n$. This gives Formula (1), where

(6) $$\mathcal{F}^{n-1;t^j}(f(t^j)\chi_{U_{1,\dots,1}}, u(p,t^j;\zeta); p; \zeta) = \int_0^\infty ... \int_0^\infty f(t^j) \exp(-u(p,t^j;\zeta)) dt^j$$

$$= \int_0^\infty dt_1 ... \int_0^\infty dt_{j-1} \int_0^\infty dt_{j+1} ... \int_0^\infty (dt_n) f(t^j) \exp(-u(p,t^j;\zeta))$$

is the non-commutative transform by $t^j = (t_1, ..., t_{j-1}, 0, t_{j+1}, ..., t_n)$.

12.1. Remark. Shift operators of the form $\xi(x+\phi) = \exp(\phi d/dx)\xi(x)$ in real variables are also frequently used in the class of infinite differentiable functions with converging Taylor series expansion in the corresponding domain.

It is possible to use also the following convention. One can put $\cos(\phi_1 + \zeta_1) = \cos(\phi_1 + \zeta_1)\cos(\psi_2)...\cos(\psi_{2^r-1})$,...,$\sin(\phi_1 + \zeta_1)...\cos(\phi_k + \zeta_k) = \sin(\phi_1 + \zeta_1)...\cos(\phi_k + \zeta_k)\cos(\psi_{k+1})...\cos(\psi_{2^r-1})$, where $\psi_j = 0$ for each $j \ge 1$, $2 \le k < 2^r - 1$, so that $T_j^l \cos(\phi_1 + \zeta_1) = 0$ for each $j > 1$ and $l \ge 1$, $T_j^l \sin(\phi_1 + \zeta_1)...\cos(\phi_k + \zeta_k) = 0$ for each $j > k$ and $l \ge 1$, where $T_j^l \xi = T_j^{l-1}(T_j \xi)$ is the iterated composition for $l > 1$, $l \in \mathbf{N}$. Then $T_j^l e^{-u(p,t;\zeta)}$ gives with such convention the same result as $\mathsf{S}_{e_j}^l e^{-u(p,t;\zeta)}$, so one can use the symbolic notation $T_j^l e^{-u(p,t;\zeta)} = e^{-u(p,t;\zeta - i_j \pi l/2)}$. But to avoid misunderstanding we shall use S_{e_j} and T_j in the sense of Formulas 12(3.1 − 3.7).

It is worth to mention that instead of 12(3.7) also the formulas

(1) $\exp(p_1 i_1 + ... + p_n i_n) = \cos(\phi) + M \sin(\phi)$ with $\phi := \phi(p) := \left[p_1^2 + ... + p_n^2\right]^{1/2}$ and $M = (p_1 i_1 + ... + p_n i_n)/\phi$ for $\phi \ne 0$, $e^0 = 1$;

(2) $\partial \exp(p_1 i_1 + .. + p_n i_n)/\partial p_j = \left[-\sin(\phi) + M\cos(\phi)\right] p_j/\phi + (\phi i_j - M p_j)\phi^{-2}\sin(\phi)$ and $\partial(p_j t_j + \zeta_j)/\partial \zeta_j = 1$ can be used.

13. Theorem. *Let $f(t)$ be a function-original. Suppose that $u(p,t;\zeta)$ is given by 2(1, 2, 2.1) or 1(8, 8.1) over the Cayley-Dickson algebra $\mathcal{A}_r$ with $2 \le r < \infty$. Then a (super)derivative of an image is given by the following formula:*

(1) $\left(\partial \mathcal{F}^n(f(t), u; p; \zeta)/\partial p\right).h = -\mathcal{F}^n\left(f(t)s_1, u; p; \zeta\right)h_0$
$-\mathsf{S}_{e_1}\mathcal{F}^n\left(f(t)s_1, u; p; \zeta\right)h_1 - ... - \mathsf{S}_{e_n}\mathcal{F}^n\left(f(t)s_n, u; p; \zeta\right)h_n$

in the $\mathcal{A}_r$ spherical coordinates, or

(1.1) $\left(\partial \mathcal{F}^n(f(t), u; p; \zeta)/\partial p\right).h = -\mathcal{F}^n\left(f(t)s_1, u; p; \zeta\right)h_0$
$-\mathsf{S}_{e_1}\mathcal{F}^n\left(f(t)t_1, u; p; \zeta\right)h_1 - ... - \mathsf{S}_{e_n}\mathcal{F}^n\left(f(t)t_n, u; p; \zeta\right)h_n$

in the $\mathcal{A}_r$ Cartesian coordinates for each $h = h_0 i_0 + ... + h_n i_n \in \mathcal{A}_r$, where $h_0, ..., h_n \in \mathbf{R}$, $2^{r-1} \le n \le 2^r - 1$, $p \in W_f$.

Proof. The inequalities $a_1(f) < Re(p) < a_{-1}(f)$ are equivalent to the inequalities $a_1(f(t)|t|) < Re(p) < a_{-1}(f(t)|t|)$, since $\lim_{|t|\to+\infty} \exp(-b|t|)|t| = 0$ for each $b > 0$. An image $\mathcal{F}^n(f(t), u; p; \zeta)$ is a holomorphic function by p for $a_1(f) < Re(p) < a_{-1}(f)$ by Theorem 4, also $|\int_0^\infty e^{-ct} t^n dt| < \infty$ for each $c > 0$ and $n = 0, 1, 2, ...$ Thus it is possible to differentiate under the sign of the integral:

(2) $$\left(\partial\left(\int_{\mathbf{R}^n} f(t)\exp(-u(p,t;\zeta))dt\right)/\partial p\right).h =$$

$$\sum_{v\in\{-1,1\}^n}\left(\partial\left(\int_{U_v} f(t)\exp(-u(p,t;\zeta))\chi_{U_v}dt\right)/\partial p\right).h =$$

$$= \int_{\mathbf{R}^n} f(t)(\partial\exp(-u(p,t;\zeta))/\partial p).hdt.$$

Due to Formulas 12(3, 3.2) we get:

(3) $(\partial\exp(-u(p,t;\zeta))/\partial p).h = -\exp(-u(p,t;\zeta))s_1h_0 - \mathsf{S}_{e_1}\exp(-u(p,t;\zeta))s_1h_1 - ... - \mathsf{S}_{e_n}\exp(-u(p,t;\zeta))s_nh_n$
in the $\mathcal{A}_r$ spherical coordinates, or

(4) $(\partial\exp(-u(p,t;\zeta))/\partial p).h = -\exp(-u(p,t;\zeta))s_1h_0 - \mathsf{S}_{e_1}\exp(-u(p,t;\zeta))t_1h_1 - ... - \mathsf{S}_{e_n}\exp(-u(p,t;\zeta))t_nh_n$

in the $\mathcal{A}_r$ Cartesian coordinates.
Thus from Formulas (2, 3) we deduce Formula (1).

14. Theorem. *If $f(t)$ is a function-original, then*

(1) $\mathcal{F}^n(f(t-\tau),u;p;\zeta) = \mathcal{F}^n(f(t),u;p;\zeta+ <p,\tau])$ *for either*

(*i*) $u(p,t;\zeta) = p_0s_1 + M(p,t;\zeta) + \zeta_0$ *or*

(*ii*) $u(p,t;\zeta) =<p,t)+\zeta$ *over* $\mathcal{A}_r$ *with* $2\le r<\infty$ *in a domain* $p\in W_f$, *where* $\tau\in\mathbf{R}^n$, $2^{r-1}\le n\le 2^r-1$,

(2) $<p,\tau] = p_0s_1 + p_1s_1i_1 + ... + p_ns_ni_n$ *with* $s_j = s_j(n;\tau)$ *for each* j *in the first* (*i*) *and* $<p,\tau] =<p,\tau)$ *in the second* (*ii*) *case (see also Formulas 1*(8), *2*(1, 2, 2.1)*).*

Proof. For p in the domain $Re(p) > a_1$ the identities are satisfied:

$$(3)\qquad \mathcal{F}^n\big((f\chi_{U_{1,\dots,1}})(t-\tau),u;p;\zeta\big) = \int_{\tau_1}^{\infty}...\int_{\tau_n}^{\infty} f(t-\tau)e^{-u(p,t;\zeta)}dt$$

$$= \int_{U_{1,\dots,1}} f(t)e^{-u(p,\xi;\zeta+<p,\tau])}d\xi = \mathcal{F}^n\big((f\chi_{U_{1,\dots,1}})(t),u;p;\zeta+<p,\tau]\big),$$

due to Formulas 1(7, 8) and 2(1, 2, 2.1, 4), since $p_0s_1(n;t)+\zeta_0 = p_0s_1(n;\xi)+\zeta_0+p_0s_1(n;\tau)$ and $p_jt_j+\zeta_j = p_j\xi_j + (\zeta_j+p_j\tau_j)$ and $p_js_j(n;t)+\zeta_j = p_js_j(n;\xi)+(\zeta_j+p_js_j(n;\tau))$ for each $j=1,...,2^r-1$, where $t=\xi+\tau$. Symmetrically we get (2) for U_v instead of $U_{1,\dots,1}$. Naturally, that the multiparameter non-commutative Laplace integral for an original f can be considered as the sum of 2^n integrals by the sub-domains U_v:

$$(4)\qquad \int_{\mathbf{R}^n} f(t)\exp\big(-u(p,t;\zeta)\big)dt = \sum_{v\in\{-1,1\}^n}\int_{\mathbf{R}^n} f(t)\exp\big(-u(p,t;\zeta)\big)\chi_{U_v}(t)dt.$$

The summation by all possible $v\in\{-1,1\}^n$ gives Formula (1).

15. Note. In view of the definition of the non-commutative transform $\mathcal{F}^n$ and $u(p,t;\zeta)$ and Theorem 14 the term $\zeta_1i_1+...+\zeta_{2^r-1}i_{2^r-1}$ has the natural interpretation as the initial phase of a retardation.

16. Theorem. *If $f(t)$ is a function-original with values in $\mathcal{A}_r$ for $2\le r<\infty$, $2^{r-1}\le n\le 2^r-1$, $b\in\mathbf{R}$, then*

$$(1)\qquad \mathcal{F}^n\big(e^{b(t_1+\dots+t_n)}f(t),u;p;\zeta\big) = \mathcal{F}^n\big(f(t),u;p-b;\zeta\big)$$

for each $a_{-1}+b > Re(p) > a_1+b$, where u is given by 1(8, 8.1) *or 2*(1, 2).

Proof. In accordance with Expressions 1(8, 8.1) and 2(1, 2, 2.1) one has $u(p,t;\zeta) - b(t_1 + ... + t_n) = u(p-b,t;\zeta)$. If $a_{-1} + b > Re(p) > a_1 + b$, then the integral

$$(2) \qquad \mathcal{F}^n\left(e^{b(t_1+...+t_n)} f(t)\chi_{U_v}(t), u; p; \zeta\right) = \int_{U_v} f(t)e^{b(t_1+...+t_n)} \exp(-u(p,t;\zeta))dt$$

$$= \int_{U_v} f(t)\exp(-u(p-b,t;\zeta))dt = \mathcal{F}^n(f(t)\chi_{U_v}(t), u; p-b; \zeta)$$

converges. Applying Decomposition 14(4) we deduce Formula (1).

17. Theorem. *Let a function $f(t)$ be a real valued original, $F(p;\zeta) = \mathcal{F}^n(f(t); u; p; \zeta)$, where the function $u(p,t;\zeta)$ is given by 1(8, 8.1) or 2(1, 2, 2.1). Let also $G(p;\zeta)$ and $q(p)$ be locally analytic functions such that*

(1) $\mathcal{F}^n(g(t,\tau); u; p; \zeta) = G(p;\zeta)\exp(-u(q(p),\tau;\zeta))$
for $u =< p,t) + \zeta$ or $u = p_0(t_1 + ... + t_n) + M(p,t;\zeta) + \zeta_0$, then

(2) $\mathcal{F}^n(\int_{\mathbf{R}^n} g(t,\tau)f(\tau)d\tau; u; p; \zeta) = G(p;\zeta)F(q(p);\zeta)$
for each $p \in W_g$ and $q(p) \in W_f$, where $2 \le r < \infty$, $2^{r-1} \le n \le 2^r - 1$.

Proof. If $p \in W_g$ and $q(p) \in W_f$, then in view of the Fubini's theorem and the theorem conditions a change of an integration order gives the equalities:

$$\int_{\mathbf{R}^n}\left(\int_{\mathbf{R}^n} g(t,\tau)f(\tau)d\tau\right)\exp(-u(p,t;\zeta))dt$$

$$= \int_{\mathbf{R}^n}\left(\int_{\mathbf{R}^n} g(t,\tau)\exp(-u(p,t;\zeta))dt\right)f(\tau)d\tau$$

$$= \int_{\mathbf{R}^n} G(p;\zeta)\exp(-u(q(p),\tau;\zeta))f(\tau)d\tau$$

$$= G(p;\zeta)\int_{\mathbf{R}^n} f(\tau)\exp(-u(q(p),\tau;\zeta))d\tau = G(p;\zeta)F(q(p);\zeta),$$

since $t,\tau \in \mathbf{R}^n$ and the center of the algebra $\mathcal{A}_r$ is $\mathbf{R}$.

18. Theorem. *If a function $f(t)\chi_{U_{1,...,1}}$ is original together with its derivative $\partial^n f(t)\chi_{U_{1,...,1}}(t)/\partial s_1...\partial s_n$ or $\partial^n f(t)\chi_{U_{1,...,1}}(t)/\partial t_1...\partial t_n$, where $F^n_u(p;\zeta)$ is an image function of $f(t)\chi_{U_{1,...,1}}$ over the Cayley-Dickson algebra $\mathcal{A}_r$ with $2 \le r \in \mathbf{N}$, $2^{r-1} \le n \le 2^r - 1$, for $u = p_0 s_1 + M(p,t;\zeta) + \zeta_0$ given by 2(1, 2, 2.1), then*

$$(1) \qquad \lim_{p\to\infty}\Big\{[p_0 + p_1\mathsf{S}_{e_1}]p_2\mathsf{S}_{e_2}...p_n\mathsf{S}_{e_n}F^n_u(p;\zeta) + \sum_{m=0}^{n-1}(-1)^m$$

$$\sum_{1\le j_1<...<j_{n-m}\le n;\ 1\le l_1<...<l_m\le n;\ l_\alpha\ne j_\beta\ \forall \alpha,\beta} [p_0\delta_{1,j_1} + p_{j_1}\mathsf{S}_{e_{j_1}}]p_{j_2}\mathsf{S}_{e_{j_2}}...p_{j_{n-m}}\mathsf{S}_{e_{j_{n-m}}}$$

$$F^{n-m}_u(p^{(l)};\zeta)\Big\} = (-1)^{n+1}f(0)e^{-u(0,0;\zeta)},$$

or

$$(1.1) \qquad \lim_{p\to\infty}\Big\{[p_0 + p_1\mathsf{S}_{e_1}][p_0 + p_2\mathsf{S}_{e_2}]...[p_0 + p_n\mathsf{S}_{e_n}]F^n_u(p;\zeta) + \sum_{m=0}^{n-1}(-1)^m$$

$$\sum_{1\le j_1<...<j_{n-m}\le n;\ 1\le l_1<...<l_m\le n;\ l_\alpha\ne j_\beta\ \forall \alpha,\beta} [p_0+p_{j_1}\mathsf{S}_{e_{j_1}}][p_0+p_{j_2}\mathsf{S}_{e_{j_2}}]...[p_0+p_{j_{n-m}}\mathsf{S}_{e_{j_{n-m}}}]$$

$$F^{n-m}_u(p^{(l)};\zeta)\Big\}=(-1)^{n+1}f(0)e^{-u(0,0;\zeta)}$$

for $u(p,t;\zeta)$ *given by 1(8, 8.1), where* $f(0)=\lim_{t\in U_{1,...,1};t\to 0}f(t)$, p *tends to the infinity inside the angle* $|Arg(p)|<\pi/2-\delta$ *for some* $0<\delta<\pi/2$, $1\le j\le 2^r-1$, $p^{(l)}=\sum_{j=0,j\notin(l)}^n p_j i_j$, $(l)=(l_1,...,l_m)$. *If the restriction*

$f(t)|_{t_{j_1}=0,...,t_{j_m}=0;t_k=\infty\forall k\notin\{j_1,...,j_m\}}=\lim_{t\in U_{1,...,1};t_{j_1}\to 0,...,t_{j_m}\to 0;t_k\to\infty\ \forall k\notin\{j_1,...,j_m\}}f(t)$ *exists for all* $1\le j_1<...<j_m\le n$, *then*

(2) $$\lim_{p\to 0}\Big\{[p_0+p_1\mathsf{S}_{e_1}]p_2\mathsf{S}_{e_2}...p_n\mathsf{S}_{e_n}F^n_u(p;\zeta)+\sum_{m=0}^{n-1}(-1)^m$$

$$\sum_{1\le j_1<...<j_{n-m}\le n;\ 1\le l_1<...<l_m\le n;\ l_\alpha\ne j_\beta\ \forall \alpha,\beta}[p_0\delta_{1,j_1}+p_{j_1}\mathsf{S}_{e_{j_1}}]p_{j_2}\mathsf{S}_{e_{j_2}}...p_{j_{n-m}}\mathsf{S}_{e_{j_{n-m}}}$$

$$F^{n-m}_u(p^{(l)};\zeta)\Big\}$$

$$=\sum_{m=0}^{n-1}(-1)^m\sum_{1\le j_1<...<j_m\le n}f(t)\big|_{t_{j_1}=0,...,t_{j_m}=0;t_k=\infty\forall k\notin\{j_1,...,j_m\}}e^{-u(0,0,\zeta)}$$

in the $\mathcal{A}_r$ *spherical coordinates or*

(2.1) $$\lim_{p\to 0}\Big\{[p_0+p_1\mathsf{S}_{e_1}][p_0+p_2\mathsf{S}_{e_2}]...[p_0+p_n\mathsf{S}_{e_n}]F^n_u(p;\zeta)+\sum_{m=0}^{n-1}(-1)^m$$

$$\sum_{1\le j_1<...<j_{n-m}\le n;\ 1\le l_1<...<l_m\le n;\ l_\alpha\ne j_\beta\ \forall \alpha,\beta}[p_0+p_{j_1}\mathsf{S}_{e_{j_1}}][p_0+p_{j_2}\mathsf{S}_{e_{j_2}}]...[p_0+p_{j_{n-m}}\mathsf{S}_{e_{j_{n-m}}}]$$

$$F^{n-m}_u(p^{(l)};\zeta)\Big\}$$

$$=\sum_{m=0}^{n-1}(-1)^m\sum_{1\le j_1<...<j_m\le n}f(t)\big|_{t_{j_1}=0,...,t_{j_m}=0;t_k=\infty\forall k\notin\{j_1,...,j_m\}}e^{-u(0,0,\zeta)}$$

in the $\mathcal{A}_r$ *Cartesian coordinates, where* $p\to 0$ *inside the same angle.*

Proof. In accordance with Theorem 12 the equality follows:

(3) $$\mathcal{F}^n((\partial f(t)/\partial s_j)\chi_{U_{1,...,1}}(t),u;p;\zeta)=[p_0\delta_{1,j}+p_j\mathsf{S}_{e_j}]\mathcal{F}^n(f(t)\chi_{U_{1,...,1}}(t),u(p,t;\zeta),p;\zeta)$$

$$-\mathcal{F}^{n-1;t^j}(f(t^j)\chi_{U_{1,...,1}},u(p,t^j;\zeta);p;\zeta)$$

for $u=u(p,t;\zeta)=p_0s_1+M(p,t;\zeta)+\zeta_0$ in the $\mathcal{A}_r$ spherical coordinates, or

(3.1) $$\mathcal{F}^n((\partial f(t)/\partial t_j)\chi_{U_{1,...,1}}(t),u;p;\zeta)=[p_0+p_j\mathsf{S}_{e_j}]\mathcal{F}^n(f(t)\chi_{U_{1,...,1}}(t),u(p,t;\zeta),p;\zeta)$$

$$-\mathcal{F}^{n-1;t^j}(f(t^j)\chi_{U_{1,...,1}},u(p,t^j;\zeta);p;\zeta)$$

in the $\mathcal{A}_r$ Cartesian coordinates, since

(3.2) $\partial f(t(s))/\partial s_j=-\partial f(t)/\partial t_{j-1}+\partial f(t)/\partial t_j$ for each $j\ge 2$, $\partial f(t(s))/\partial s_1=\partial f(t)/\partial t_1$, where $p=p_0+p_1i_1+...+p_{2^r-1}i_{2^r-1}\in\mathcal{A}_r$, $p_0,...,p_{2^r-1}\in\mathbf{R}$, $\{i_0,...,i_{2^r-1}\}$ are the generators of the Cayley-Dickson algebra $\mathcal{A}_r$, $s_{n+l}=0$ for each $l\ge 1$, the zero power $\mathsf{S}^0_{e_j}=I$ is the unit operator. For short we write f instead of $f\chi_{U_{1,...,1}}$. Thus the limit exists:

$$(4)\qquad \mathcal{F}^{n-1;t^j}(f(t^j),u(p,t^j;\zeta);p;\zeta)=$$

$$\lim_{t_j\to +0}\int_0^\infty dt_1...\int_0^\infty dt_{j-1}\int_0^\infty dt_{j+1}...\int_0^\infty (dt_n)f(t)\exp(-u(p,t;\zeta)).$$

Mention, that $(...((t^1)^2)...)^j=(0,...,0,t_j,...,t_n:t_j=0)$ for every $1\le j\le n$, since $t_k=s_k-s_{k+1}$ for each $1\le k\le n$. We apply these Formulas $(3,4)$ by induction $j=1,...,n$, $2^{r-1}\le n\le 2^r-1$, to $\partial^n f(t)/\partial s_1...\partial s_n$,...,$\partial^{n-j+1}f(t)/\partial s_j...\partial s_n$, ...,$\partial f(t)/\partial s_n$ instead of $\partial f(t)/\partial s_j$.

From Note 8 [18] it follows, that in the $\mathcal{A}_r$ spherical coordinates

$$\lim_{p\to\infty,|Arg(p)|<\pi/2-\delta}\mathcal{F}^n((\partial^n f(t)/\partial s_1...\partial s_n)\chi_{U_{1,..,1}},u;p;\zeta)=0,$$

also in the $\mathcal{A}_r$ Cartesian coordinates

$$\lim_{p\to\infty,|Arg(p)|<\pi/2-\delta}\mathcal{F}^n((\partial^n f(t)/\partial t_1...\partial t_n)\chi_{U_{1,..,1}},u;p;\zeta)=0,$$

which gives the first statement of this theorem, since $u(p,0,\zeta)=u(0,t;\zeta)=u(0,0,\zeta)$ and $F^0_u(p^{(1,...,1)};\zeta)=f(0)e^{-u(0,0,\zeta)}$, while $F^n_u(p;\zeta)$ is defined for each $Re(p)>0$.

If the limit $f(t^{<j>})$ exists, where $t^{<j>}:=(t_1,...,t_j,...,t_n:\ t_j=\infty)$, then

$$(5)\qquad \lim_{t_j\to\infty}\int_0^\infty dt_1...\int_0^\infty dt_{j-1}\int_0^\infty dt_{j+1}...\int_0^\infty (dt_n)f(t)\exp(-u(p,t;\zeta))$$

$$=:\mathcal{F}^{n-1;<t^j>}(f(t^{<j>}),u(p,t^{<j>};\zeta);p;\zeta).$$

Certainly, $(...((t^{<1>})^{<2>})...)^{<j>}=(t_1,...,t_n:t_1=\infty,...,t_j=\infty)$ for each $1\le j\le n$. Therefore, the limit exists:

$$\lim_{p\to 0,|Arg(p)|<\pi/2-\delta}\int_{U_{1,...,1}}(\partial^n f(t)/\partial s_1...\partial s_n)\exp(-p_0s_1-\zeta_0-M(p,t;\zeta))$$

$$=\int_{U_{1,...,1}}(\partial^n f(t)/\partial s_1...\partial s_n)e^{-u(0,0;\zeta)}dt$$

$$=\sum_{m=0}^n(-1)^m\sum_{1\le j_1<...<j_m\le n}f(t)|_{t_{j_1}=0,...,t_{j_m}=0;t_k=\infty\ \forall k\notin\{j_1,...,j_m\}}$$

$$=\lim_{p\to 0,|Arg(p)|<\pi/2-\delta}\Big\{[p_0+p_1\mathsf{S}_{e_1}]p_2\mathsf{S}_{e_2}...p_n\mathsf{S}_{e_n}F^n_u(p;\zeta)$$

$$+\sum_{m=0}^{n-1}(-1)^m$$

$$\sum_{1\le j_1<...<j_{n-m}\le n;\ 1\le l_1<...<l_m\le n;\ l_\alpha\ne j_\beta\ \forall\alpha,\beta}[p_0\delta_{1,j_1}+p_{j_1}\mathsf{S}_{e_{j_1}}]p_{j_2}\mathsf{S}_{e_{j_2}}...p_{j_{n-m}}\mathsf{S}_{e_{j_{n-m}}}$$

$$F^{n-m}_u(p^{(l)};\zeta)+(-1)^nf(0)e^{-u(0,0,\zeta)}\Big\},$$

from which the second statement of this theorem follows in the $\mathcal{A}_r$ spherical coordinates and analogously in the $\mathcal{A}_r$ Cartesian coordinates using Formula (3.1).

19. Definitions. Let X and Y be two $\mathbf{R}$ linear normed spaces which are also left and right $\mathcal{A}_r$

modules, where $1 \le r$. Let Y be complete relative to its norm. We put $X^{\otimes k} := X \otimes_{\mathbf{R}} ... \otimes_{\mathbf{R}} X$ is the k times ordered tensor product over $\mathbf{R}$ of X. By $L_{q,k}(X^{\otimes k}, Y)$ we denote a family of all continuous k times $\mathbf{R}$ poly-linear and $\mathcal{A}_r$ additive operators from $X^{\otimes k}$ into Y. Then $L_{q,k}(X^{\otimes k}, Y)$ is also a normed $\mathbf{R}$ linear and left and right $\mathcal{A}_r$ module complete relative to its norm. In particular, $L_{q,1}(X, Y)$ is denoted also by $L_q(X, Y)$.

We present X as the direct sum $X = X_0 i_0 \oplus ... \oplus X_{2^r-1} i_{2^r-1}$, where X_0,...,X_{2^r-1} are pairwise isomorphic real normed spaces. If $A \in L_q(X, Y)$ and $A(xb) = (Ax)b$ or $A(bx) = b(Ax)$ for each $x \in X_0$ and $b \in \mathcal{A}_r$, then an operator A we call right or left $\mathcal{A}_r$-linear respectively.

An $\mathbf{R}$ linear space of left (or right) k times $\mathcal{A}_r$ poly-linear operators is denoted by $L_{l,k}(X^{\otimes k}, Y)$ (or $L_{r,k}(X^{\otimes k}, Y)$ respectively).

We consider a space of test function $\mathcal{D} := \mathcal{D}(\mathbf{R}^n, Y)$ consisting of all infinite differentiable functions $f : \mathbf{R}^n \to Y$ on $\mathbf{R}^n$ with compact supports. A sequence of functions $f_n \in \mathcal{D}$ tends to zero, if all f_n are zero outside some compact subset K in the Euclidean space $\mathbf{R}^n$, while on it for each $k = 0, 1, 2, ...$ the sequence $\{f_n^{(k)} : n \in \mathbf{N}\}$ converges to zero uniformly. Here as usually $f^{(k)}(t)$ denotes the k-th derivative of f, which is a k times $\mathbf{R}$ poly-linear symmetric operator from $(\mathbf{R}^n)^{\otimes k}$ to Y, that is $f^{(k)}(t).(h_1, ..., h_k) = f^{(k)}(t).(h_{\sigma(1)}, ..., h_{\sigma(k)}) \in Y$ for each $h_1, ..., h_k \in \mathbf{R}^n$ and every transposition $\sigma : \{1, ..., k\} \to \{1, ..., k\}$, σ is an element of the symmetric group S_k, $t \in \mathbf{R}^n$. For convenience one puts $f^{(0)} = f$. In particular, $f^{(k)}(t).(e_{j_1}, ..., e_{j_k}) = \partial^k f(t)/\partial t_{j_1} ... \partial t_{j_k}$ for all $1 \le j_1, ..., j_k \le n$, where $e_j = (0, ..., 0, 1, 0, ..., 0) \in \mathbf{R}^n$ with 1 on the j-th place.

Such convergence in $\mathcal{D}$ defines closed subsets in this space $\mathcal{D}$, their complements by the definition are open, that gives the topology on $\mathcal{D}$. The space $\mathcal{D}$ is $\mathbf{R}$ linear and right and left $\mathcal{A}_r$ module.

By a generalized function of class $\mathcal{D}' := [\mathcal{D}(\mathbf{R}^n, Y)]'$ is called a continuous $\mathbf{R}$-linear $\mathcal{A}_r$-additive function $g : \mathcal{D} \to \mathcal{A}_r$. The set of all such functionals is denoted by $\mathcal{D}'$. That is, g is continuous, if for each sequence $f_n \in \mathcal{D}$, converging to zero, a sequence of numbers $g(f_n) =: [g, f_n) \in \mathcal{A}_r$ converges to zero for n tending to the infinity.

A generalized function g is zero on an open subset V in $\mathbf{R}^n$, if $[g, f) = 0$ for each $f \in \mathcal{D}$ equal to zero outside V. By a support of a generalized function g is called the family, denoted by $supp(g)$, of all points $t \in \mathbf{R}^n$ such that in each neighborhood of each point $t \in supp(g)$ the functional g is different from zero. The addition of generalized functions g, h is given by the formula:

(1) $[g + h, f) := [g, f) + [h, f)$.

The multiplication $g \in \mathcal{D}'$ on an infinite differentiable function w is given by the equality:

(2) $[gw, f) = [g, wf)$ either for $w : \mathbf{R}^n \to \mathcal{A}_r$ and each test function $f \in \mathcal{D}$ with a real image $f(\mathbf{R}^n) \subset \mathbf{R}$, where $\mathbf{R}$ is embedded into Y; or $w : \mathbf{R}^n \to \mathbf{R}$ and $f : \mathbf{R}^n \to Y$.

A generalized function g' prescribed by the equation:

(3) $[g', f) := -[g, f')$ is called a derivative g' of a generalized function g, where $f' \in \mathcal{D}(\mathbf{R}^n, L_q(\mathbf{R}^n, Y))$, $g' \in [\mathcal{D}(\mathbf{R}^n, L_q(\mathbf{R}^n, Y))]'$.

Another space $\mathcal{B} := \mathcal{B}(\mathbf{R}^n, Y)$ of test functions consists of all infinite differentiable functions $f : \mathbf{R}^n \to Y$ such that the limit $\lim_{|t| \to +\infty} |t|^m f^{(j)}(t) = 0$ exists for each $m = 0, 1, 2, ...$, $j = 0, 1, 2, ...$. A sequence $f_n \in \mathcal{B}$ is called converging to zero, if the sequence $|t|^m f_n^{(j)}(t)$ converges to zero uniformly on $\mathbf{R}^n \setminus B(\mathbf{R}^n, 0, R)$ for each $m, j = 0, 1, 2, ...$ and each $0 < R < +\infty$, where $B(Z, z, R) := \{y \in Z : \rho(y, z) \le R\}$ denotes a ball with center at z of radius R in a metric space Z with a metric ρ. The family of all $\mathbf{R}$-linear and $\mathcal{A}_r$-additive functionals on $\mathcal{B}$ is denoted by $\mathcal{B}'$.

In particular we can take $X = \mathcal{A}_r^{\alpha}$, $Y = \mathcal{A}_r^{\beta}$ with $1 \le \alpha, \beta \in \mathbf{Z}$. Analogously spaces $\mathcal{D}(U, Y)$, $[\mathcal{D}(U, Y)]'$, $\mathcal{B}(U, Y)$ and $[\mathcal{B}(U, Y)]'$ are defined for domains U in $\mathbf{R}^n$, for example, $U = U_v$ (see also §1).

A generalized function $f \in \mathcal{B}'$ we call a generalized original, if there exist real numbers $a_1 < a_{-1}$ such that for each $a_1 < w_{-1}, w_1, ..., w_{-n}, w_n < a_{-1}$ the generalized function

(4) $f(t)\exp(-(q_v, t))\chi_{U_v}$ is in $[\mathcal{B}(U_v, Y)]'$ for all $v = (v_1, ..., v_n)$, $v_j \in \{-1, 1\}$ for every $j = 1, ..., n$ for each $t \in \mathbf{R}^n$ with $t_j v_j \ge 0$ for each $j = 1, ..., n$, where $q_v = (v_1 w_{v_1 1}, ..., v_n w_{v_n n})$.

By an image of such original we call a function

(5) $\mathcal{F}^n(f, u; p; \zeta) := [f, \exp(-u(p, t; \zeta)))$ of the variable $p \in \mathcal{A}_r$ with the parameter $\zeta \in \mathcal{A}_r$, defined in the domain $W_f = \{p \in \mathcal{A}_r : \ a_1 < Re(p) < a_{-1}\}$ by the following rule. For a given $p \in W_f$ choose $a_1 < w_1, ..., w_n < Re(p) < w_{-1}, ..., w_{-n} < a_{-1}$, then

(6) $[f, \exp(-u(p, t; \zeta)) := \sum_v [f\exp(-(q_v, t)), \exp\{-[u(p, t; \zeta) - (q_v, t)]\}\chi_{U_v})$,
since $\exp\{-[u(p, t; \zeta) - (q_v, t)]\} \in \mathcal{B}(U_v, Y)$, where in each term
$[f\exp(-(q_v, t)), \exp\{-[u(p, t; \zeta) - (q_v, t)]\}\chi_{U_v})$ the generalized function belongs to $[\mathcal{B}(U_v, Y)]'$ by Condition (4), while the sum in (6) is by all admissible vectors $v \in \{-1, 1\}^n$.

20. Note and Examples. Evidently the transform $\mathcal{F}^n(f, u; p; \zeta)$ does not depend on a choice of $\{w_{-1}, w_1, ..., w_{-n}, w_n\}$, since

$[f\exp(-(q_v, t), \exp(-[u(p, t; \zeta) - (q_v, t)])\chi_{U_v}) =$
$[f\exp(-(q_v, t) - (b_v, t)), \exp(-[u(p, t; \zeta) - (q_v, t) - (b_v, t)])\chi_{U_v})$

for each $b \in \mathbf{R}^n$ such that $a_1 < w_j + b_j < Re(p) < w_{-j} + b_{-j} < a_{-1}$ for each $j = 1, ..., n$, because $\exp(-(b_v, t)) \in \mathbf{R}$. At the same time the real field $\mathbf{R}$ is the center of the Cayley-Dickson algebra $\mathcal{A}_r$, where $2 \le r \in \mathbf{N}$.

Let δ be the Dirac delta function, defined by the equation

(DF) $[\delta(t), \phi(t)) := \phi(0)$ for each $\phi \in \mathcal{B}$. Then

(1) $\mathcal{F}^n(\delta^{(j)}(t-\tau), u; p; \zeta) = \sum_{v \in \{-1,1\}^n} [\delta^{(j)}(t-\tau)\exp(-(q_v, t)), \exp(-[u(p, t; \zeta) - (q_v, t)])\chi_{U_v})$
$= (-1)^j \partial_t^j \exp(-[u(p, t; \zeta)])|_{t=\tau}$,

since it is possible to take $-\infty < a_1 < 0 < a_{-1} < \infty$ and $w_k = 0$ for each $k \in \{-1, 1, -2, 2, ..., -n, n\}$, where $\tau \in \mathbf{R}^n$ is the parameter, $\partial_t^j := \partial^{|j|}/\partial t_1^{j_1} ... \partial t_1^{j_1}$. In particular, for $j = 0$ we have

(2) $\mathcal{F}^n(\delta(t - \tau), u; p; \zeta) = \exp(-u(p, \tau; \zeta))$.

In the general case:

(3) $\mathcal{F}^n(\partial^{|j|}\delta(t)/\partial s_1^{j_1} ... \partial s_n^{j_n}, u; p; \zeta) =$
$\sum_{0 \le k_1 \le j_1} \binom{j_1}{k_1} p_0^{j_1 - k_1} (p_1 \mathsf{S}_{e_1})^{k_1} (p_2 \mathsf{S}_{e_2})^{j_2} ... (p_n \mathsf{S}_{e_n})^{j_n} \exp(-\zeta_0 - M(p, 0; \zeta))$

in the $\mathcal{A}_r$ spherical coordinates, or

(3.1) $\mathcal{F}^n(\partial^{|j|}\delta(t)/\partial t_1^{j_1} ... \partial t_n^{j_n}, u; p; \zeta) =$
$(p_0 + p_1 \mathsf{S}_{e_1})^{j_1} (p_0 + p_2 \mathsf{S}_{e_2})^{j_2} ... (p_0 + p_n \mathsf{S}_{e_n})^{j_n} \exp(-u(p, 0; \zeta))$

in the $\mathcal{A}_r$ Cartesian coordinates, where $j_1 + ... + j_n = |j|$, $k_1, j_1, ..., j_n$ are nonnegative integers, $2^{r-1} \le n \le 2^r - 1$, $\binom{l}{m} := l!/[m!(l-m)!]$ denotes the binomial coefficient, $0! = 1$, $1! = 1$, $2! = 2$; $l! = 1 \cdot 2 \cdot ... \cdot l$ for each $l \ge 3$, $s_j = s_j(n; t)$.

The transform $\mathcal{F}^n(f)$ of any generalized function f is the holomorphic function by $p \in W_f$ and by $\zeta \in \mathcal{A}_r$, since the right side of Equation 19(5) is holomorphic by p in W_f and by ζ in view of Theorem 4. Equation 19(5) implies, that Theorems 11 - 13 are accomplished also for generalized functions.

For $a_1 = a_{-1}$ the region of convergence reduces to the vertical hyperplane in $\mathcal{A}_r$ over $\mathbf{R}$. For $a_{-1} < a_1$ there is no any common domain of convergence and $f(t)$ can not be transformed.

21. Theorem. *If $f(t)$ is an original function on $\mathbf{R}^n$, $F^n(p; \zeta)$ is its image, $\partial^{|j|} f(t)/\partial s_1^{j_1} ... \partial s_n^{j_n}$*

or $\partial^{|j|}f(t)/\partial t_1^{j_1}...\partial t_n^{j_n}$ *is an original,* $|j| = j_1 + ... + j_n$, $0 \le j_1, ..., j_n \in \mathbf{Z}$, $2^{r-1} \le n \le 2^r - 1$; *then*

(1) $$\mathcal{F}^n\left(\partial^{|j|}f(t)/\partial s_1^{j_1}...\partial s_n^{j_n}, u; p; \zeta\right)$$

$$= \sum_{0\le k_1\le j_1} \binom{j_1}{k_1} p_0^{j_1-k_1}(p_1\mathsf{S}_{e_1})^{k_1}(p_2\mathsf{S}_{e_2})^{j_2}...(p_n\mathsf{S}_{e_n})^{j_n}\mathcal{F}^n(f(t), u; p; \zeta)$$

for $u(p,t;\zeta) := p_0s_1 + M(p,t;\zeta) + \zeta_0$ *given by* 2$(1, 2, 2.1)$*, or*

(1.1) $$\mathcal{F}^n\left(\partial^{|j|}f(t)/\partial t_1^{j_1}...\partial t_n^{j_n}, u; p; \zeta\right) =$$

$$(p_0 + p_1\mathsf{S}_{e_1})^{j_1}(p_0 + p_2\mathsf{S}_{e_2})^{j_2}...(p_0 + p_n\mathsf{S}_{e_n})^{j_n}\mathcal{F}^n(f(t), u; p; \zeta)$$

for $u(p,t;\zeta)$ *given by* 1$(8, 8.1)$ *over the Cayley-Dickson algebra* $\mathcal{A}_r$ *with* $2 \le r < \infty$*. Domains, where Formulas* $(1, 1.1)$ *are true may be different from a domain of the multiparameter noncommutative transform for* f*, but they are satisfied in the domain* $a_1 < Re(p) < a_{-1}$*, where*

$a_{-1} = \min(a_{-1}(f), a_{-1}(\partial^{|m|}f(t)/\partial\phi_1^{m_1}...\partial\phi_n^{m_n}) : |m| \le |j|, 0 \le m_l \le j_l\forall l)$;

$a_1 = \max(a_1(f), a_1(\partial^{|m|}f(t)/\partial\phi_1^{m_1}...\partial\phi_n^{m_n}) : |m| \le |k|, 0 \le m_l \le j_l\forall l)$*, if* $a_1 < a_{-1}$*, where* $\phi_j = s_j$ *or* $\phi_j = t_j$ *for each* j *correspondingly.*

Proof. To each domain U_v the domain U_{-v} symmetrically corresponds. The number of different vectors $v \in \{-1, 1\}^n$ is even 2^n. Therefore, for $u = p_0t + \zeta_0 + M(p,t;\zeta)$ due to Theorem 12 the equality

(2) $$\int_{\mathbf{R}^n}(\partial f(t)/\partial s_j)e^{-u(p,t;\zeta)}ds = \int_{\mathbf{R}^n}(\partial f(t)/\partial s_j)e^{-u(p,t;\zeta)}dt =$$

$$\int_{\mathbf{R}^{n-1}}(dt^j)\left[f(t)e^{-u(p,t;\zeta)}\right]\Big|_{-\infty}^{\infty} - \int_{\mathbf{R}^{n-1}}(dt^j)\left(\int_{-\infty}^{\infty}f(t)[\partial e^{-u(p,t;\zeta)}/\partial s_j]ds_j\right)$$

is satisfied in the $\mathcal{A}_r$ spherical coordinates, since the absolute value of the Jacobian $\partial t/\partial(t^j, s_j)$ is unit. Since for $a_1 < Re(p) < a_{-1}$ the first additive is zero, while the second integral converts with the help of Formulas 12$(2, 2.1)$, Formula (1) follows for $k = 1$:

(3) $\mathcal{F}^n(\partial f(t)/\partial s_j, u; p; \zeta) = p_0\delta_{1,j}\mathcal{F}^n(f(t), u; p; \zeta) + p_j\mathsf{S}_{e_j}\mathcal{F}^n(f(t), u; p; \zeta)$.

To accomplish the derivation we use Theorem 14 so that

$$\lim_{\tau\to 0}\Big[\mathcal{F}^n(f(t), u; p; \zeta) - \mathcal{F}^n(f(t - \tau e_j), u; p; \zeta)\Big]/\tau$$

$$= \lim_{\tau\to 0}\Big[\mathcal{F}^n(f(t), u; p; \zeta) - \mathcal{F}^n(f(t), u; p; \zeta + \tau(p_0 + p_1i_1 + ... + p_ji_j))\Big]/\tau$$

$$= \lim_{\tau\to 0}\int_{\mathbf{R}^n}f(t)\Big[e^{-u(p,t;\zeta)} - e^{-u(p,t;\zeta+\tau(p_0+p_1i_1+...+p_ji_j))}\Big]\tau^{-1}dt,$$

where $e_j = (0, ..., 0, 1, 0, .., 0) \in \mathbf{R}^n$ with 1 on the j-th place. If the original $\partial^{|j|}f(t)/\partial s_1^{j_1}...\partial s_n^{j_n}$ exists, then $\partial^{|m|}f(t)/\partial s_1^{m_1}...\partial s_n^{m_n}$ is continuous for $0 \le |m| \le |j| - 1$ with $0 \le m_l \le j_l$ for each $l = 1, ..., n$, where $f^0 := f$. The interchanging of $\lim_{\tau\to 0}$ and $\int_{\mathbf{R}^n}$ may change a domain of convergence, but in the indicated in the theorem domain $a_1 < Re(p) < a_{-1}$, when it is non void, Formula (3) is valid. Applying Formula (3) in the $\mathcal{A}_r$ spherical coordinates by induction to $(\partial^{|m|}f(t)/\partial s_1^{m_1}...\partial s_n^{m_n}) : |m| \le |j|, 0 \le m_l \le j_l\forall l)$ with the corresponding order subordinated to $\partial^{|j|}f(t)/\partial s_1^{j_1}...\partial s_n^{j_n}$, or in the $\mathcal{A}_r$ Cartesian coordinates using Formula 12(1.1) for the partial derivatives $(\partial^{|m|}f(t)/\partial t_1^{m_1}...\partial t_n^{m_n}) : |m| \le |j|, 0 \le m_l \le j_l\forall l)$ with the corresponding order subordinated to $\partial^{|j|}f(t)/\partial t_1^{j_1}...\partial t_n^{j_n}$ we deduce Expressions (1) and (1.1) with the help of Statement 6 from §XVII.2.3 [30] about the differentiation of an improper integral by a parameter and §2.

22. Remarks. For the entire Euclidean space $\mathbf{R}^n$ Theorem 21 for $\partial f(t)/\partial s_j$ gives only one or two additives on the right side of 21(1) in accordance with 21(3).

Evidently Theorems 4, 11 and Proposition 10 are accomplished for $\mathcal{F}^{k;t_{j(1)},...,t_{j(k)}}(f,u;p;\zeta)$ also.

Theorem 12 is satisfied for $\mathcal{F}^{k;t_{j(1)},...,t_{j(k)}}$ and any $j\in\{j(1),...,j(k)\}$, so that $s_l=s_l(k;t)=t_{j(l)}+...+t_{j(k)}$ for each $1\le l\le k$, $p_m=0$ and $\zeta_m=0$ for each $1\le m\notin\{j(1),...,j(k)\}$ (the same convention is in 13, 14, 17, 21, see also below). For $\mathcal{F}^{k;t_{j(1)},...,t_{j(k)}}$ in Theorem 13 in Formula 13(1) it is natural to put $t_m=0$ and $h_m=0$ for each $1\le m\notin\{j(1),...,j(k)\}$, so that only $(k+1)$ additives with h_0, $h_{j(1)}$,...,$h_{j(k)}$ on the right side generally may remain. Theorems 14 and 17 and 21 modify for $\mathcal{F}^{k;t_{j(1)},...,t_{j(k)}}$ putting in 14(1) and 17(1, 2) and 21(1) $t_j=0$ and $\tau_j=0$ respectively for each $j\notin\{j(1),...,j(k)\}$.

To take into account boundary conditions for domains different from U_v, for example, for bounded domains V in $\mathbf{R}^n$ we consider a bounded noncommutative multiparameter transform

(1) $\mathcal{F}^n(f(t)\chi_V,u;p;\zeta)=:\mathcal{F}^n_V(f(t),u;p;\zeta)$.

For it evidently Theorems 4, 6-8, 11, 13, 14, 16, 17, Proposition 10 and Corollary 4.1 are satisfied as well taking specific originals f with supports in V.

At first take domains W which are quadrants, that is canonical closed subsets affine diffeomorphic with $Q^n=\prod_{j=1}^n[\mathsf{a}_j,b_j]$, where $-\infty\le\mathsf{a}_j<b_j\le\infty$, $[\mathsf{a}_j,b_j]:=\{x\in\mathbf{R}:\ \mathsf{a}_j\le x\le b_j\}$ denotes the segment in $\mathbf{R}$. This means that there exists a vector $w\in\mathbf{R}^n$ and a linear invertible mapping C on $\mathbf{R}^n$ so that $C(W)-w=Q$. We put $t^{j,1}:=(t_1,...,t_j,...,t_n:\ t_j=\mathsf{a}_j)$, $t^{j,2}:=(t_1,...,t_j,...,t_n:\ t_j=b_j)$. Consider $t=(t_1,...,t_n)\in Q^n$.

23. Theorem. *Let $f(t)$ be a function-original with a support by t variables in Q^n and zero outside Q^n such that $\partial f(t)/\partial t_j$ also satisfies Conditions 1$(1-4)$. Suppose that $u(p,t;\zeta)$ is given by 2$(1,2,2.1)$ or 1$(8,8.1)$ over $\mathcal{A}_r$ with $2\le r<\infty$, $2^{r-1}\le n\le 2^r-1$. Then*

$$\mathcal{F}^n\big((\partial f(t)/\partial t_j)\chi_{Q^n}(t),u;p;\zeta\big)=\tag{1}$$
$$\mathcal{F}^{n-1;t^{j,2}}\big(f(t^{j,2})\chi_{Q^n}(t^{j,2}),u;p;\zeta\big)-\mathcal{F}^{n-1;t^{j,1}}\big(f(t^{j,1})\chi_{Q^n}(t^{j,1}),u;p;\zeta\big)$$
$$+\Big[p_0+\sum_{k=1}^{j}p_k\mathsf{S}_{e_k}\Big]\mathcal{F}^n\big(f(t)\chi_{Q^n}(t),u;p;\zeta\big)$$

in the $\mathcal{A}_r$ spherical coordinates, or

$$\mathcal{F}^n\big((\partial f(t)/\partial t_j)\chi_{Q^n}(t),u;p;\zeta\big)=\tag{1.1}$$
$$\mathcal{F}^{n-1;t^{j,2}}\big(f(t^{j,2})\chi_{Q^n}(t^{j,2}),u;p;\zeta\big)-\mathcal{F}^{n-1;t^{j,1}}\big(f(t^{j,1})\chi_{Q^n}(t^{j,1}),u;p;\zeta\big)$$
$$+\big[p_0+p_j\mathsf{S}_{e_j}\big]\mathcal{F}^n\big(f(t)\chi_{Q^n}(t),u;p;\zeta\big)$$

in the $\mathcal{A}_r$ Cartesian coordinates in a domain $W\subset\mathcal{A}_r$; if $\mathsf{a}_j=-\infty$ or $b_j=+\infty$, then the addendum with $t^{j,1}$ or $t^{j,2}$ correspondingly is zero.

Proof. Here the domain Q^n is bounded and f is almost everywhere continuous and satisfies Conditions 1$(1-4)$, hence $f(t)\exp(-u(p,t;\zeta))\in L^1(\mathbf{R}^n,\lambda_n,\mathcal{A}_r)$ for each $p\in\mathcal{A}_r$, since $\exp(-u(p,t;\zeta))$ is continuous and $supp(f(t))\subset Q^n$.

Analogously to §12 the integration by parts gives

$$\int_{\mathsf{a}_j}^{b_j}(\partial f(t)/\partial t_j)\exp(-u(p,t;\zeta))dt_j=f(t)\exp(-u(p,t;\zeta))\Big|_{t_j=\mathsf{a}_j}^{t_j=b_j}\tag{2}$$

$$-\int_{\mathsf{a}_j}^{b_j}\big[f(t)(\partial\exp(-u(p,t;\zeta))/\partial t_j)\big]dt_j,$$

where $t=(t_1,...,t_n)$. Then the Fubini's theorem implies:

$$(3)\qquad \int_{Q^n}(\partial f(t)/\partial t_j)\exp(-u(p,t;\zeta))dt=$$

$$\int_{\mathsf{a}_1}^{b_1}...\int_{\mathsf{a}_{j-1}}^{b_{j-1}}\int_{\mathsf{a}_{j+1}}^{b_{j+1}}\int_{\mathsf{a}_n}^{b_n}\left[\int_{\mathsf{a}_j}^{b_j}(\partial f(t)/\partial t_j)\exp(-u(p,t;\zeta))dt_j\right]dt^j$$

$$=\left[\int_{t\in Q^n,\ t_j=b_j}f(t^{j,2})\exp(-u(p,t^{j,2};\zeta))dt^j\right]-\left[\int_{t\in Q^n,\ t_j=\mathsf{a}_j}f(t^{j,1})\exp(-u(p,t^{j,1};\zeta))dt^j\right]$$

$$+\left[p_0+\sum_{k=1}^{j}p_k\mathsf{S}_{e_k}\right]\int_{\mathsf{a}_1}^{b_1}...\int_{\mathsf{a}_n}^{b_n}f(t)\exp(-u(p,t;\zeta))dt$$

in the $\mathcal{A}_r$ spherical coordinates or

$$(3.1)\qquad \int_{Q^n}(\partial f(t)/\partial t_j)\exp(-u(p,t;\zeta))dt$$

$$=\left[\int_{t\in Q^n,\ t_j=b_j}f(t^{j,2})\exp(-u(p,t^{j,2};\zeta))dt^j\right]-\left[\int_{t\in Q^n,\ t_j=\mathsf{a}_j}f(t^{j,1})\exp(-u(p,t^{j,1};\zeta))dt^j\right]$$

$$+[p_0+p_j\mathsf{S}_{e_j}]\int_{\mathsf{a}_1}^{b_1}...\int_{\mathsf{a}_n}^{b_n}f(t)\exp(-u(p,t;\zeta))dt$$

in the $\mathcal{A}_r$ Cartesian coordinates, where as usually $t^j=(t_1,...,t_{j-1},0,t_{j+1},...,t_n)$, $dt^j=dt_1...dt_{j-1}dt_{j+1}...dt_n$. This gives Formulas $(1,1.1)$, where

$$(4)\qquad \mathcal{F}^{n-1;t^{j,k}}(f(t^{j,k})\chi_{Q^n}(t^{j,k}),u(p,t^{j,k};\zeta);p;\zeta)=$$

$$\int_{\mathsf{a}_1}^{b_1}...\int_{\mathsf{a}_{j-1}}^{b_{j-1}}\int_{\mathsf{a}_{j+1}}^{b_{j+1}}\int_{\mathsf{a}_n}^{b_n}f(t^{j,k})\exp(-u(p,t^{j,k};\zeta))dt^{j,k}$$

is the non-commutative transform by $t^{j,k}$, $2^{r-1}\le n\le 2^r-1$, $dt^{j,k}$ is the Lebesgue volume element on $\mathbf{R}^{n-1}$.

24. Theorem. *If a function $f(t)\chi_{Q^n}(t)$ is original together with its derivative $\partial^n f(t)\chi_{Q^n}(t)/\partial s_1...\partial s_n$ or $\partial^n f(t)\chi_{Q^n}(t)/\partial t_1...\partial t_n$, where $F^n_u(p;\zeta)$ is an image function of $f(t)\chi_{Q^n}(t)$ over the Cayley-Dickson algebra $\mathcal{A}_r$ with $2\le r\in\mathbf{N}$, $2^{r-1}\le n\le 2^r-1$, for the function $u(p,t;\zeta)$ given by 2$(1,2,2.1)$ or 1$(8,8.1)$, $Q^n=\prod_{j=1}^n[0,b_j]$, $b_j>0$ for each j, then*

$$(1)\qquad \lim_{p\to\infty}\Big\{[p_0+p_1\mathsf{S}_{e_1}]p_2\mathsf{S}_{e_2}...p_n\mathsf{S}_{e_n}F^n_u(p;\zeta)+\sum_{m=0}^{n-1}(-1)^m$$

$$\sum_{1\le j_1<...<j_{n-m}\le n;\ 1\le l_1<...<l_m\le n;\ l_\alpha\ne j_\beta\ \forall\alpha,\beta}[p_0\delta_{1,j_1}+p_{j_1}\mathsf{S}_{e_{j_1}}]p_{j_2}\mathsf{S}_{e_{j_2}}...p_{j_{n-m}}\mathsf{S}_{e_{j_{n-m}}}$$

$$F^{n-m}_u(p^{(l)};\zeta)\Big\}=(-1)^{n+1}f(0)e^{-u(0,0;\zeta)}$$

in the $\mathcal{A}_r$ spherical coordinates, or

$$(1.1)\quad \lim_{p\to\infty}\Big\{[p_0+p_1\mathsf{S}_{e_1}][p_0+p_2\mathsf{S}_{e_2}]...[p_0+p_n\mathsf{S}_{e_n}]F^n_u(p;\zeta)+\sum_{m=0}^{n-1}(-1)^m$$

$$\sum_{1\le j_1<...<j_{n-m}\le n;\ 1\le l_1<...<l_m\le n;\ l_\alpha\ne j_\beta\ \forall\alpha,\beta}[p_0+p_{j_1}\mathsf{S}_{e_{j_1}}][p_0+p_{j_2}\mathsf{S}_{e_{j_2}}]...[p_0+p_{j_{n-m}}\mathsf{S}_{e_{j_{n-m}}}]$$

$$F^{n-m}_u(p^{(l)};\zeta)\Big\}=(-1)^{n+1}f(0)e^{-u(0,0;\zeta)}$$

in the $\mathcal{A}_r$ Cartesian coordinates, where $f(0)=\lim_{t\in Q^n,\ t\to 0}f(t)$, p tends to the infinity inside the angle $|Arg(p)|<\pi/2-\delta$ for some $0<\delta<\pi/2$.

Proof. In accordance with Theorem 23 we have Equalities 23(1, 1.1). Therefore we infer that

$$(2)\quad \mathcal{F}^{n-1;t^{j,k}}\big(f(t^{j,k})\chi_{Q^n}(t^{j,k}),u(p,t^{j,k};\zeta);p;\zeta\big)=$$

$$\lim_{t_j\to\beta_{j,k}+0}\int_{\mathsf{a}_1}^{b_1}dt_1...\int_{\mathsf{a}_{j-1}}^{b_{j-1}}dt_{j-1}\int_{\mathsf{a}_{j+1}}^{b_{j+1}}dt_{j+1}...\int_{\mathsf{a}_n}^{b_n}(dt_n)f(t)\exp(-u(p,t;\zeta)),$$

where $\beta_{j,1}=\mathsf{a}_j=0$, $\beta_{j,2}=b_j>0$, $k=1,2$. Mention, that $(...((t^{1,l_1})^{2,l_2})...)^{j,l_j}=(t:\ t_1=\beta_{1,l_1},...,t_j=\beta_{j,l_j})$ for every $1\le j\le n$. Analogously to §12 we apply Formula (2) by induction $j=1,...,n$, $2^{r-1}\le n\le 2^r-1$, to

$\partial^n f(t(s))/\partial s_1...\partial s_n$,...,$\partial^{n-j+1}f(t(s))/\partial s_j...\partial s_n$,...,$\partial f(t(s))/\partial s_n$

instead of $\partial f(t(s))/\partial s_j$, $s_j=s_j(n;t)$ as in §2, or applying to the partial derivatives

$\partial^n f(t)/\partial t_1...\partial t_n$,...,$\partial^{n-j+1}f(t)/\partial t_j...\partial t_n$,...,$\partial f(t)/\partial t_n$

instead of $\partial f(t)/\partial t_j$ correspondingly. If $s_j>0$ for some $j\ge 1$, then $s_1>0$ for Q^n and $\lim_{p\to\infty}e^{-u(p,t^{(l)};\zeta)}=0$ for such $t^{(l)}$, where $t=(t_1,...,t_n)$, $(l)=(l_1,...,l_n)$, $|l|=l_1+...+l_n$, $t^{(l)}=(t^{(l)}_1,...,t^{(l)}_n)$, $t^{(l)}_j=\mathsf{a}_j$ for $l_j=1$ and $t^{(l)}_j=b_j$ for $l_j=2$, $1\le j\le 2^r-1$. Therefore,

$$\lim_{p\to\infty}\sum_{l_j\in\{1,2\};\ j=1,...,n}(-1)^{|l|}f(t^{(l)})e^{-u(p,t^{(l)};\zeta)}=(-1)^n f(0)e^{-u(0,0;\zeta)},$$

since $u(p,0;\zeta)=u(0,0;\zeta)$, where $f(^{(l)})=\lim_{t\in Q^n;t\to t^{(l)}}f(t)$.

In accordance with Note 8 [18]

$$\lim_{p\to\infty,|Arg(p)|<\pi/2-\delta}\mathcal{F}^n((\partial^n f(t)/\partial s_1...\partial s_n)\chi_{Q^n}(t),u(p,t;\zeta);p;\zeta)=0$$

in the $\mathcal{A}_r$ spherical coordinates and

$$\lim_{p\to\infty,|Arg(p)|<\pi/2-\delta}\mathcal{F}^n((\partial^n f(t)/\partial t_1...\partial t_n)\chi_{Q^n}(t),u(p,t;\zeta);p;\zeta)=0$$

in the $\mathcal{A}_r$ Cartesian coordinates, which gives the statement of this theorem.

25. Suppose that $f(t)\chi_{Q^n}(t)$ is an original function, $F^n(p;\zeta)$ is its image, $\partial^{|j|}f(t)\chi_{Q^n}(t)/\partial t_1^{j_1}...\partial t_n^{j_n}$ is an original, $|j|=j_1+...+j_n$, $0\le j_1,...,j_n\in\mathbf{Z}$, $2^{r-1}\le n\le 2^r-1$, $-\infty\le\mathsf{a}_k<b_k\le\infty$ for each $k=1,...,n$, $(l)=(l_1,...,l_n)$, $l_k\in\{0,1,2\}$, $W=\mathcal{A}_r$ for bounded Q^n. Let $W=\{p\in\mathcal{A}_r:\ a_1<Re(p)\}$ for $b_k=\infty$ for some k and finite a_k for each k; $W=\{p\in\mathcal{A}_r:\ Re(p)<a_{-1}\}$ for $\mathsf{a}_k=-\infty$ for some k and finite b_k for each k; $W=\{p\in\mathcal{A}_r:\ a_1<Re(p)<a_{-1}\}$ when $\mathsf{a}_k=-\infty$ and $b_l=+\infty$ for some k and l; $t^{(l)}=(t^{(l)}_1,...,t^{(l)}_n)$.

We put $t_k^{(l)} = t_k$ and $q_k = 0$ for $l_k = 0$, $t_k^{(l)} = \mathsf{a}_k$ for $l_k = 1$, $t_k^{(l)} = b_k$ for $l_k = 2$, $(q) = (q_1, ..., q_n)$, $|q| = q_1 + ... + q_n$,

$a_1 = \max(a_1(f), a_1(\partial^{|m|} f(t)/\partial t_1^{m_1}...\partial t_n^{m_n}) : |m| \le |j|, 0 \le m_k \le j_k \forall k)$,

$a_{-1} = \min(a_{-1}(f), a_{-1}(\partial^{|m|} f(t)/\partial t_1^{m_1}...\partial t_n^{m_n}) : |m| \le |j|, 0 \le m_k \le j_k \ \forall k)$ if $a_1 < a_{-1}$.

If $\mathsf{a}_k = -\infty$ and $b_k = +\infty$ for Q^n with a given k, then $l_k = 0$. If either $\mathsf{a}_k > -\infty$ or $b_k < +\infty$ for a marked k, then $l_k \in \{0, 1, 2\}$. We also put $h_k = h_k(l) = sign(l_k)$ for each k, where $sign(x) = -1$ for $x < 0$, $sign(0) = 0$, $sign(x) = 1$ for $x > 0$, $h = h(l)$, $|h| = |h_1| + ... + |h_n|$,

$(lj) := (l_1 sign(j_1), ..., l_n sign(j_n))$.

Let the vector (l) enumerate faces $\partial Q^n_{(l)}$ in ∂Q^n_{k-1} for $|h(l)| = k \ge 1$, so that $\partial Q^n_{k-1} = \bigcup_{|h(l)|=k} Q^n_{(l)}$, $\partial Q^n_{(l)} \cap \partial Q^n_{(m)} = \emptyset$ for each $(l) \ne (m)$ (see also more detailed notations in §28).

Let the shift operator be defined:

$T_{(m)} F(p; \zeta) := F(p; \zeta - (i_1 m_1 + ... + i_n m_n)\pi/2)$, also the operator

(SO) $\mathsf{S}_{(m)} F(p; \zeta) := \mathsf{S}_{e_1}^{m_1} ... \mathsf{S}_{e_n}^{m_n} F(p; \zeta)$,

where $(m) = (m_1, ..., m_n) \in [0, \infty)^n \subset \mathbf{R}^n$, $\mathsf{S}^k_{(m)} = \mathsf{S}_{k(m)}$ for each positive number $0 < k \in \mathbf{R}$, $\mathsf{S}_0 = I$ is the unit operator for $(m) = 0$ (see also Formulas 12(3.1 − 3.7)). As usually let $e_1 = (1, 0, ..., 0)$,...,$e_n = (0, ..., 0, 1)$ be the standard orthonormal basis in $\mathbf{R}^n$ so that $(m) = m_1 e_1 + ... + m_n e_n$.

Theorem. *Then*

$$(1) \quad \mathcal{F}^n\Big(\partial^{|j|} f(t)\chi_{Q^n}(t)/\partial t_1^{j_1}...\partial t_n^{j_n}, u(p, t; \zeta); p; \zeta\Big) =$$

$$\mathsf{R}_{e_1}^{j_1} \mathsf{R}_{e_2}^{j_2} ... \mathsf{R}_{e_n}^{j_n} \mathcal{F}^n(f(t)\chi_{Q^n}(t), u; p; \zeta)$$

$$+ \sum_{1 \le |(lj)|;\ m_k + q_k + h_k = j_k; 0 \le m_k;\ 0 \le q_k;\ h_k = sign(l_k j_k);\ q_k = 0 \text{ for } l_k j_k = 0, \text{ for each } k = 1, ..., n;\ (l) \in \{0,1,2\}^n}$$

$$(-1)^{|(lj)|} \mathsf{R}_{e_1}^{m_1} \mathsf{R}_{e_2}^{m_2} ... \mathsf{R}_{e_n}^{m_n} \mathcal{F}^{n-|h(lj)|}(\partial^{|q|} f(t^{(lj)})\chi_{\partial Q^n_{(lj)}}(t^{(lj)})/\partial t_1^{q_1}...\partial t_n^{q_n}, u; p; \zeta)$$

for $u(p, t; \zeta)$ in the $\mathcal{A}_r$ spherical coordinates or the $\mathcal{A}_r$ Cartesian coordinates over the Cayley-Dickson algebra $\mathcal{A}_r$ with $2 \le r < \infty$, where

(1.1) $\mathsf{R}_{e_1} := p_0 + p_1 \mathsf{S}_{e_1}$, $\mathsf{R}_{e_2} := p_0 + p_1 \mathsf{S}_{e_1} + p_2 \mathsf{S}_{e_2}$,..., $\mathsf{R}_{e_n} := p_0 + p_1 \mathsf{S}_{e_1} + p_2 \mathsf{S}_{e_2} + ... + p_n \mathsf{S}_{e_n}$ *in the $\mathcal{A}_r$ spherical coordinates, while*

(1.2) $\mathsf{R}_{e_1} := p_0 + p_1 \mathsf{S}_{e_1}$, $\mathsf{R}_{e_2} := p_0 + p_2 \mathsf{S}_{e_2}$,..., $\mathsf{R}_{e_n} := p_0 + p_n \mathsf{S}_{e_n}$ *in the $\mathcal{A}_r$ Cartesian coordinates, i.e. $\mathsf{R}_{e_j} = \mathsf{R}_{e_j}(p)$ are operators depending on the parameter p. If $t_j^{(l)} = \infty$ for some $1 \le j \le n$, then the corresponding addendum on the right of (1) is zero.*

Proof. In view of Theorem 23 we get the equality

$$(2) \quad \int_{Q^n} \Big[(\partial^{|m|+1} f(t)/\partial t_1^{m_1}...\partial t_{k-1}^{m_{k-1}} \partial t_k^{m_k+1} \partial t_{k+1}^{m_{k+1}}...\partial t_n^{m_n}) e^{-u(p,t;\zeta)}\Big] dt =$$

$$\int_{\mathbf{R}^{n-1} \cap Q^n} (dt^k) \Big[(\partial^{|m|} f(t)/\partial t_1^{m_1}...\partial t_n^{m_n}) e^{-u(p,t;\zeta)}\Big]\Big|_{\mathsf{a}_k}^{b_k}$$

$$- \int_{\mathbf{R}^{n-1} \cap Q^n} (dt^k) \bigg(\int_{\mathsf{a}_k}^{b_k} \Big(\partial^{|m|} f(t)/\partial t_1^{m_1}...\partial t_n^{m_n}\Big) \Big[\partial e^{-u(p,t;\zeta)}/\partial t_k\Big] dt_k\bigg)$$

is satisfied for $0 \le m_k \le j_k$ for each $k = 1, ..., n$ with $|m| < |j|$. On the other hand, for $p \in W$ additives on the right of (2) convert with the help of Formula 23(1). Each term of the form

$$\int_{\mathbf{R}^{n-|h(l)|} \cap Q^n} (dt^{(l)}) \Big[(\partial^{|q|} f(t^{(l)})\chi_{\partial Q^n_{(l)}}(t^{(l)})/\partial t_1^{q_1}...\partial t_n^{q_n} e^{-u(p,t;\zeta)}\Big]$$

can be further transformed with the help of (2) by the considered variable t_k only in the case $l_k = 0$. Applying Formula (2) by induction to partial derivatives $\partial^{|j|} f/\partial t_1^{j_1}...\partial t_n^{j_n}$, $\partial^{|j|-j_1} f/\partial t_2^{j_2}...\partial t_n^{j_n}$,...,$\partial^{j_n} f/\partial t_n^{j_n}$,...,$\partial f/\partial t_n$ as in §21 and using Theorem 14 and Remarks 22 we deduce (1).

26. Theorem. *Let $f(t)\chi_{U_{1,...,1}}(t)$ be a function-original with values in $\mathcal{A}_r$ with $2 \le r < \infty$, $2^{r-1} \le n \le 2^r - 1$, u is given by 2(1, 2, 2.1) or 1(8, 8.1),*

$$(1) \qquad g(t) := \int_0^{t_1} ... \int_0^{t_n} f(x)dx, \text{ then}$$

$$(2) \qquad \mathcal{F}^n(f\chi_{U_{1,...,1}}(t), u; p; \zeta) = \mathsf{R}_{e_1}\mathsf{R}_{e_2}...\mathsf{R}_{e_n}\mathcal{F}^n(g(t)\chi_{U_{1,...,1}}(t), u; p; \zeta)$$

in the domain $Re(p) > \max(a_1, 0)$, where the operators R_{e_j} are given by Formulas 25(1.1, 1.2).

Proof. In view of Theorem 25 the equation

$$(3) \qquad \mathcal{F}^n(f\chi_{U_{1,...,1}}(t), u; p; \zeta) =$$

$$\mathsf{R}_{e_1}\mathsf{R}_{e_2}...\mathsf{R}_{e_n}\mathcal{F}^n(g(t), u; p; \zeta)$$

$$+ \sum_{1\le|l|;\ 0\le m_k\le 1;\ m_k+h_k=1;\ h_k=sign(l_k);\ \text{for each } k=1,...,n;\ q_1=0,...,q_n=0}$$

$$(-1)^{|(l)|}\mathsf{R}_{e_1}^{m_1}\mathsf{R}_{e_2}^{m_2}...\mathsf{R}_{e_n}^{m_n}\mathcal{F}^{n-|h(l)|}(g(t^{(l)}), u; p; \zeta),$$

is satisfied, since $\partial^n g(t)/\partial t_1...\partial t_n = (f\chi_{U_{1,...,1}})(t)$, where $j_1 = 1,...,j_n = 1$, $l_j = 1$ for each $j = 1, ..., n$. Equation (3) is accomplished in the same domain $Re(p) > \max(a_1, 0)$, since $g(0) = 0$ and $g(t)$ also fulfills conditions of Definition 1, while $a_1(g) < \max(a_1(f), 0) + b$ for each $b > 0$, where $a_1 \in \mathbf{R}$. On the other hand, $g(t)$ is equal to zero on $\partial U_{1,...,1}$ and outside $U_{1,...,1}$ in accordance with formula (1), hence all terms on the right side of Equation (3) with $|l| > 0$ disappear and $supp(g(t)) \subset U_{1,...,1}$. Thus we get Equation (2).

27. Theorem. *Suppose that $F^k(p;\zeta)$ is an image $\mathcal{F}^{k;t_1,...,t_k}(f(t)\chi_{U_{1,...,1}}(t), u; p; \zeta)$ of an original function $f(t)$ for u given by 2(1, 2, 2.1) in the half space $W := \{p \in \mathcal{A}_r : Re(p) > a_1\}$ with $2 \le r < \infty$, $p_1 = 0,...,p_{j-1} = 0$; $\zeta_1 = \pi/2,...,\zeta_{j-1} = \pi/2$ for each $j \ge 2$ in the $\mathcal{A}_r$ spherical coordinates or $\zeta_1 = 0,...,\zeta_{j-1} = 0$ for each $j \ge 2$ in the $\mathcal{A}_r$ Cartesian coordinates;*

(1) *the integral $\int_{p_j i_j}^{\infty i_j} F^k(p_0 + z; \zeta)dz$ converges, where $p = p_0 + p_1 i_1 + ... + p_k i_k \in \mathcal{A}_r$, $p_j \in \mathbf{R}$ for each $j = 0, ..., 2^r - 1$, $2^{r-1} \le k \le 2^r - 1$, $U_{1,...,1} := \{(t_1, ..., t_k) \in \mathbf{R}^k :\ t_1 \ge 0, ..., t_k \ge 0\}$. Let also*

(2) *the function $F^k(p;\zeta)$ be continuous by the variable $p \in \mathcal{A}_r$ on the open domain W, moreover, for each $w > a_1$ there exist constants $C_w{}' > 0$ and $\epsilon_w > 0$ such that*

(3) *$|F^k(p;\zeta)| \le C_w{}' \exp(-\epsilon_w|p|)$ for each $p \in S_{R(n)}$, $S_R := \{z \in \mathcal{A}_r :\ Re(z) \ge w\}$, $0 < R(n) < R(n+1)$ for each $n \in \mathbf{N}$, $\lim_{n\to\infty} R(n) = \infty$, where a_1 is fixed, $\zeta = \zeta_0 i_0 + ... + \zeta_k i_k \in \mathcal{A}_r$ is marked, $\zeta_j \in \mathbf{R}$ for each $j = 0, ..., k$. Then*

$$(4) \qquad \int_{p_j i_j}^{\infty i_j} F^k(p_0 + z; \zeta)dz = \mathcal{S}_{-e_j}\mathcal{F}^{k;t_1,...,t_k}(f(t)\chi_{U_{1,...,1}}(t)/\xi_j, u; p; \zeta),$$

where $p_1 = 0,...,p_{j-1} = 0$ for each $j \ge 2$; $\zeta_1 = \pi/2,...,\zeta_{j-1} = \pi/2$ and $\xi_j = s_j(k;t)$ in the $\mathcal{A}_r$ spherical coordinates, while $\zeta_1 = 0,...,\zeta_{j-1} = 0$ and $\xi_j = t_j$ in the $\mathcal{A}_r$ Cartesian coordinates correspondingly for each $j \ge 1$.

Proof. Take a path of an integration belonging to the half space $Re(p) \geq w$ for some constant $w > a_1$. Then

$$\left| \int_{U_{1,...,1}} f(t) \exp(-u(p,t;\zeta))dt \right| \leq C \int_{U_{1,...,1}} \exp(-(p_0 - a_1)(t_1 + ... + t_k))dt < \infty$$

converges, where $C = const > 0$, $p_0 \geq w$. For $t_j > 0$ for each $j = 1, ..., k$ conditions of Lemma 2.23 [18] (that is of the noncommutative analog over $\mathcal{A}_r$ of Jordan's lemma) are satisfied. If $t_j \to \infty$, then $s_j \to \infty$, since all t_1,...,t_k are non-negative. Up to a set $\partial U_{1,...,1}$ of λ_k Lebesgue measure zero we can consider that $t_1 > 0$,...,$t_k > 0$. If $s_j \to \infty$, then also $s_1 \to \infty$. The converging integral can be written as the following limit:

$$(5) \qquad \int_{p_j i_j}^{\infty i_j} F^k(p_0 + z; \zeta)dz = \lim_{0<\kappa\to 0} \int_{p_j i_j}^{\infty i_j} F^k(p_0 + z; \zeta) \exp(-\kappa|z|)dz$$

for $1 \leq j \leq k$, since the integral $\int_{-S\infty}^{S\infty} [F^k(w + z; \zeta)]dz$ is absolutely converging and the limit $\lim_{\kappa\to 0} \exp(-\kappa|z|) = 1$ uniformly by z on each compact subset in $\mathcal{A}_r$, where S is a purely imaginary marked Cayley-Dickson number with $|S| = 1$. Therefore, in the integral

$$(6) \qquad \int_{p_j i_j}^{\infty i_j} F^k(p_0 + z; \zeta)dz = \int_{p_j i_j}^{\infty i_j} \left(\int_{U_{1,...,1}} f(t) \left[\exp(-u(p_0 + z, t; \zeta)) \right] dt \right) dz$$

the order of the integration can be changed in accordance with the Fubini's theorem applied componentwise to an integrand $g = g_0 i_0 + ... + g_n i_n$ with $g_l \in \mathbf{R}$ for each $l = 0, ..., n$:

$$(7) \qquad \int_{p_j i_j}^{\infty i_j} F^k(p_0 + z; \zeta)dz = \int_{U_{1,...,1}} dt \left(\int_{p_j i_j}^{\infty i_j} f(t) \exp(-u(p_0 + z, t; \zeta))dz \right)$$

$$= \int_{U_{1,...,1}} f(t) \left\{ \int_{p_j i_j}^{\infty i_j} \left[e^{-u(p_0+z,t;\zeta)} \right] dz \right\} dt.$$

Generally, the condition $p_1 = 0$,...,$p_{j-1} = 0$ and $\zeta_1 = \pi/2$,...,$\zeta_{j-1} = \pi/2$ in the $\mathcal{A}_r$ spherical coordinates or $\zeta_1 = 0$,...,$\zeta_{j-1} = 0$ in the $\mathcal{A}_r$ Cartesian coordinates for each $j \geq 2$ is essential for the convergence of such integral. We certainly have

$$(8) \quad \int_{p_j i_j}^{b_j i_j} \cos(i_j^* z \xi_j + \zeta_j)dz = \left[\sin(\theta_j \xi_j + \zeta_j)/\xi_j \right] \Big|_{\theta_j = p_j}^{\theta_j = b_j} = \left[-\cos(\theta_j \xi_j + \zeta_j + \pi/2)/\xi_j \right] \Big|_{\theta_j = p_j}^{\theta_j = b_j}$$

and

$$(9) \quad \int_{p_j i_j}^{b_j i_j} \sin(i_j^* z \xi_j + \zeta_j)dz_j = \left[-\cos(\theta_j \xi_j + \zeta_j)/\xi_j \right] \Big|_{\theta_j = p_j}^{\theta_j = b_j} = \left[-\sin(\theta_j \xi_j + \zeta_j + \pi/2)/\xi_j \right] \Big|_{\theta_j = p_j}^{\theta_j = b_j}$$

for each $\xi_j > 0$ and $-\infty < p_j < b_j < \infty$ and $j = 1, ..., k$. Applying Formulas $(5 - 9)$ and $2(1, 2, 2.1)$ or $1(8, 8.1)$ and $12(3.1 - 3.7)$ we deduce that:

$$\int_{p_j i_j}^{\infty i_j} \left[F^k(p_0 + z; \zeta) \right] dz = \mathsf{S}_{-e_j} \int_{U_{1,...,1}} [f(t)/\xi_j] \exp\{-u(p,t;\zeta)\}dt$$

$$= \mathsf{S}_{-e_j} \mathcal{F}^{k;t_1,...,t_k}(f(t)\chi_{U_{1,...,1}}(t)/\xi_j, u; p; \zeta),$$

where $t = (t_1, ..., t_k)$, $s_j = t_j + ... + t_k$ for each $1 \le j < k$, $s_k = t_k$, $\xi_j = s_j$ in the $\mathcal{A}_r$ spherical coordinates or $\xi_j = t_j$ in the $\mathcal{A}_r$ Cartesian coordinates.

28. Application of the noncommutative multiparameter transform to partial differential equations.

Consider a partial differential equation of the form:

(1) $A[f](t) = g(t)$, where

(2) $A[f](t) := \sum_{|j|\le\alpha} \mathbf{a}_j(t)(\partial^{|j|} f(t)/\partial t_1^{j_1}...\partial t_n^{j_n})$,

$\mathbf{a}_j(t) \in \mathcal{A}_\kappa$ are continuous functions, where $0 \le \kappa \in \mathbf{Z}$, $j = (j_1, ..., j_n)$, $|j| := j_1 + ... + j_n$, $0 \le j_k \in \mathbf{Z}$, α is a natural order of a differential operator A, $2 \le r$, $2^{r-1} \le n \le 2^r - 1$. Since $s_k = s_k(n;t) = t_k + ... + t_n$ for each $k = 1, ..., n$, the operator A can be rewritten in s coordinates as

(2.1) $A[f](t(s)) := \sum_{|j|\le\alpha} \mathbf{b}_j(t)(\partial^{|j|} f(t(s))/\partial s_1^{j_1}...\partial s_n^{j_n})$.

That is, there exists $\mathbf{b}_j \ne 0$ for some j with $|j| = \alpha$ and $\mathbf{b}_j = 0$ for $|j| > \alpha$, while a function $\sum_{j,|j|=\alpha} \mathbf{b}_j(t(s))s_1^{j_1}...s_n^{j_n}$ is not zero identically on the corresponding domain V. We consider that

($D1$) U is a canonical closed subset in the Euclidean space $\mathbf{R}^n$, that is $U = cl(Int(U))$, where $Int(U)$ denotes the interior of U and $cl(U)$ denotes the closure of U.

Particularly, the entire space $\mathbf{R}^n$ may also be taken. Under the linear mapping $(t_1, ..., t_n) \mapsto (s_1, ..., s_n)$ the domain U transforms onto V.

We consider a manifold W satisfying the following conditions $(i - v)$.

(i). The manifold W is continuous and piecewise C^α, where C^l denotes the family of l times continuously differentiable functions. This means by the definition that W as the manifold is of class $C^0 \cap C^\alpha_{loc}$. That is W is of class C^α on open subsets $W_{0,j}$ in W and $W \setminus (\bigcup_j W_{0,j})$ has a codimension not less than one in W.

(ii). $W = \bigcup_{j=0}^m W_j$, where $W_0 = \bigcup_k W_{0,k}$, $W_j \cap W_k = \emptyset$ for each $k \ne j$, $m = dim_{\mathbf{R}} W$, $dim_{\mathbf{R}} W_j = m - j$, $W_{j+1} \subset \partial W_j$.

(iii). Each W_j with $j = 0, ..., m-1$ is an oriented C^α-manifold, W_j is open in $\bigcup_{k=j}^m W_k$. An orientation of W_{j+1} is consistent with that of ∂W_j for each $j = 0, 1, ..., m-2$. For $j > 0$ the set W_j is allowed to be void or non-void.

(iv). A sequence W^k of C^α orientable manifolds embedded into $\mathbf{R}^n$, $\alpha \ge 1$, exists such that W^k uniformly converges to W on each compact subset in $\mathbf{R}^n$ relative to the metric $dist$.

For two subsets B and E in a metric space X with a metric ρ we put

(3) $dist(B, E) := \max\{\sup_{b\in B} dist(\{b\}, E), \sup_{e\in E} dist(B, \{e\})\}$, where
$dist(\{b\}, E) := \inf_{e\in E} \rho(b, e)$, $dist(B, \{e\}) := \inf_{b\in B} \rho(b, e)$, $b \in B$, $e \in E$.

Generally, $dim_{\mathbf{R}} W = m \le n$. Let $(e_1^k(x), ..., e_m^k(x))$ be a basis in the tangent space $T_x W^k$ at $x \in W^k$ consistent with the orientation of W^k, $k \in \mathbf{N}$.

We suppose that the sequence of orientation frames $(e_1^k(x_k), ..., e_m^k(x_k))$ of W^k at x_k converges to $(e_1(x), ..., e_m(x))$ for each $x \in W_0$, where $\lim_k x_k = x \in W_0$, while $e_1(x)$,...,$e_m(x)$ are linearly independent vectors in $\mathbf{R}^n$.

(v). Let a sequence of Riemann volume elements λ_k on W^k (see §XIII.2 [30]) induce a limit volume element λ on W, that is, $\lambda(B \cap W) = \lim_{k\to\infty}(B \cap W^k)$ for each compact canonical closed subset B in $\mathbf{R}^n$, consequently, $\lambda(W \setminus W_0) = 0$. We shall consider surface integrals of the second kind, i.e. by the oriented surface W (see (iv)), where each W_j, $j = 0, ..., m-1$ is oriented (see also §XIII.2.5 [30]).

Recall, that a subset V in $\mathbf{R}^n$ is called convex, if from $a, b \in V$ it follows that $\epsilon a + (1-\epsilon)b \in V$

for each $\epsilon \in [0,1]$.

(vi). Let a vector $w \in Int(U)$ exist so that $U - w$ is convex in $\mathbf{R}^n$ and let ∂U be connected. Suppose that a boundary ∂U of U satisfies Conditions $(i - v)$ and

(vii) let the orientations of ∂U^k and U^k be consistent for each $k \in \mathbf{N}$ (see Proposition 2 and Definition 3 [30]).

Particularly, the Riemann volume element λ_k on ∂U^k is consistent with the Lebesgue measure on U^k induced from $\mathbf{R}^n$ for each k. This induces the measure λ on ∂U as in (v).

Also the boundary conditions are imposed:

(4) $f(t)|_{\partial U} = f_0(t')$, $(\partial^{|q|} f(t)/\partial s_1^{q_1}...\partial s_n^{q_n})|_{\partial U} = f_{(q)}(t')$ for $|q| \le \alpha - 1$, where $s = (s_1, ..., s_n) \in \mathbf{R}^n$, $(q) = (q_1, ..., q_n)$, $|q| = q_1 + ... + q_n$, $0 \le q_k \in \mathbf{Z}$ for each k, $t \in \partial U$ is denoted by t', f_0, $f_{(q)}$ are given functions. Generally these conditions may be excessive, so one uses some of them or their linear combinations (see (5.1) below). Frequently, the boundary conditions

(5) $f(t)|_{\partial U} = f_0(t')$, $(\partial^l f(t)/\partial \nu^l)|_{\partial U} = f_l(t')$ for $1 \le l \le \alpha - 1$ are also used, where ν denotes a real variable along a unit external normal to the boundary ∂U at a point $t' \in \partial U_0$. Using partial differentiation in local coordinates on ∂U and (5) one can calculate in principle all other boundary conditions in (4) almost everywhere on ∂U.

Suppose that a domain U_1 and its boundary ∂U_1 satisfy Conditions $(D1, i - vii)$ and $g_1 = g\chi_{U_1}$ is an original on $\mathbf{R}^n$ with its support in U_1. Then any original g on $\mathbf{R}^n$ gives the original $g\chi_{U_2} =: g_2$ on $\mathbf{R}^n$, where $U_2 = \mathbf{R}^n \setminus U_1$. Therefore, $g_1 + g_2$ is the original on $\mathbf{R}^n$, when g_1 and g_2 are two originals with their supports contained in U_1 and U_2 correspondingly. Take now new domain U satisfying Conditions $(D1, i - vii)$ and $(D2 - D5)$:

$(D2)$ $U \supset U_1$ and $\partial U \subset \partial U_1$;

$(D3)$ if a straight line ξ containing a point w_1 (see (vi)) intersects ∂U_1 at two points y_1 and y_2, then only one point either y_1 or y_2 belongs to ∂U, where $w_1 \in U_1$, $U - w_1$ and $U_1 - w_1$ are convex; if ξ intersects ∂U_1 only at one point, then it intersects ∂U at the same point. That is,

$(D4)$ any straight line ξ through the point w_1 either does not intersect ∂U or intersects the boundary ∂U only at one point.

Take now g with $supp(g) \subset U$, then $supp(g\chi_{U_1}) \subset U_1$. Therefore, any problem (1) on U_1 can be considered as the restriction of the problem (1) defined on U, satisfying $(D1 - D4, i - vii)$. Any solution f of (1) on U with the boundary conditions on ∂U gives the solution as the restriction $f|_{U_1}$ on U_1 with the boundary conditions on ∂U.

Henceforward, we suppose that the domain U satisfies Conditions $(D1, D4, i - vii)$, which are rather mild and natural. In particular, for Q^n this means that either $a_k = -\infty$ or $b_k = +\infty$ for each k. Another example is: U_1 is a ball in $\mathbf{R}^n$ with the center at zero, $U = U_1 \cup (\mathbf{R}^n \setminus U_{1,...,1})$, $w_1 = 0$; or $U = U_1 \cup \{t \in \mathbf{R}^n : \ t_n \ge -\epsilon\}$ with a marked number $0 < \epsilon < 1/2$. But subsets $\partial U_{(l)}$ in ∂U can also be specified, if the boundary conditions demand it.

The complex field has the natural realization by 2×2 real matrices so that $\mathbf{i} = \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix}$, $\mathbf{i}^2 = -\begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$. The quaternion skew field, as it is well-known, can be realized with the help of 2×2 complex matrices with the generators $I = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$, $J = \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix}$, $K = \begin{pmatrix} \mathbf{i} & 0 \\ 0 & -\mathbf{i} \end{pmatrix}$, $L = \begin{pmatrix} 0 & -\mathbf{i} \\ -\mathbf{i} & 0 \end{pmatrix}$, or equivalently by 4×4 real matrices. Considering matrices with entries in the Cayley-Dickson algebra $\mathcal{A}_v$ one gets the complexified or quaternionified Cayley-Dickson algebras $(\mathcal{A}_v)_{\mathbf{C}}$ or $(\mathcal{A}_v)_{\mathbf{H}}$ with elements $z = aI + b\mathbf{i}$ or $z = aI + bJ + cK + eL$, where $a, b, c, e \in \mathcal{A}_v$, such that each $a \in \mathcal{A}_v$ commutes with the generators $\mathbf{i}$, I, J, K and L.

When $r = 2$, f and g have values in $\mathcal{A}_2 = \mathbf{H}$ and $2 \le n \le 4$ and coefficients of differential operators belong to $\mathcal{A}_2$, then the multiparameter noncommutative transform operates with the associative case so that

$\mathcal{F}^n(af) = a\mathcal{F}^n(f)$

for each $a \in \mathbf{H}$. The left linearity property $\mathcal{F}^n(af) = a\mathcal{F}^n(f)$ for any $a \in \mathbf{H}_{J,K,L}$ is also accomplished for either operators with coefficients in $\mathbf{R}$ or $\mathbf{C_i} = I\mathbf{R} \oplus \mathbf{iR}$ or $\mathbf{H}_{J,K,L} = I\mathbf{R} \oplus J\mathbf{R} \oplus K\mathbf{R} \oplus L\mathbf{R}$ and f with values in $\mathcal{A}_v$ with $1 \le n \le 2^v - 1$; or vice versa f with values in $\mathbf{C_i}$ or $\mathbf{H}_{J,K,L}$ and coefficients a in $\mathcal{A}_v$ but with $1 \le n \le 4$. Thus all such variants of operator coefficients $\mathbf{a}_j$ and values of functions f can be treated by the noncommutative transform. Henceforward, we suppose that these variants take place.

We suppose that $g(t)$ is an original function, that is satisfying Conditions $1(1-4)$. Consider at first the case of constant coefficients $\mathbf{a}_j$ on a quadrant domain Q^n. Let Q^n be oriented so that $\mathsf{a}_k = -\infty$ and $b_k = +\infty$ for each $k \le n - \kappa$; either $\mathsf{a}_k = -\infty$ or $b_k = +\infty$ for each $k > n - \kappa$, where $0 \le \kappa \le n$ is a marked integer number. If conditions of Theorem 25 are satisfied, then

$$(6) \quad \mathcal{F}^n(A[f](t), u; p; \zeta) = \sum_{|j| \le \alpha} \mathbf{a}_j \Big\{ [\mathsf{R}_{e_1}(p)]^{j_1} [\mathsf{R}_{e_2}(p)]^{j_2} ... [\mathsf{R}_{e_n}(p)]^{j_n} \mathcal{F}^n(f(t)\chi_{Q^n}(t), u; p; \zeta)$$

$$+ \sum_{1 \le |(lj)|;\ m_k+q_k+h_k=j_k;\ 0 \le m_k;\ 0 \le q_k;\ h_k = sign(l_k j_k);\ q_k=0 \text{ for } l_k j_k = 0, \text{ for each } k=1,...,n;\ (l) \in \{0,1,2\}^n}$$

$$(-1)^{|(lj)|} [\mathsf{R}_{e_1}(p)]^{m_1} [\mathsf{R}_{e_2}(p)]^{m_2} ... [\mathsf{R}_{e_n}(p)]^{m_n} \mathcal{F}^{n-|h(lj)|} \Big(\partial^{|q|} f(t^{(lj)}) \chi_{\partial Q^n_{(lj)}}(t^{(lj)}) / \partial t_1^{q_1} ... \partial t_n^{q_n}, u; p; \zeta \Big) \Big\}$$

$$= \mathcal{F}^n(g(t)\chi_{Q^n}(t), u; p; \zeta)$$

for $u(p,t;\zeta)$ in the $\mathcal{A}_r$ spherical or $\mathcal{A}_r$ Cartesian coordinates, where the operators $\mathsf{R}_{e_j}(p)$ are given by Formulas 25(1.1) or 25(1.2). Here (l) enumerates faces $\partial Q^n_{(l)}$ in ∂Q^n_{k-1} for $|h(l)| = k \ge 1$, so that $\partial Q^n_{k-1} = \bigcup_{|h(l)|=k} Q^n_{(l)}$, $\partial Q^n_{(l)} \cap \partial Q^n_{(m)} = \emptyset$ for each $(l) \ne (m)$ in accordance with §25 and the notation of this section.

Therefore, Equation (6) shows that the boundary conditions are necessary:

$(\partial^{|q|} f(t^{(l)}) / \partial t_1^{q_1} ... \partial t_n^{q_n})|_{\partial Q^n_{(l)}}$ for $|j| \le \alpha$, $|(lj)| \ge 1$, $\mathbf{a}_j \ne 0$, $q_k = 0$ for $l_k j_k = 0$, $m_k + q_k + h_k = j_k$, $h_k = sign(l_k j_k)$, $k = 1, ..., n$, $t^{(l)} \in \partial Q^n_{(l)}$. But $dim_{\mathbf{R}} \partial Q^n = n-1$ for $\partial Q^n \ne \emptyset$, consequently, $(\partial^{|q|} f(t^{(l)}) / \partial t_1^{q_1} ... \partial t_n^{q_n})|_{\partial Q^n_{(l)}}$ can be calculated if know $(\partial^{|\beta|} f(t^{(l)}) / \partial t_{\gamma(1)}^{\beta_1} ... \partial t_{\gamma(m)}^{\beta_m})|_{\partial Q^n_{(l)}}$ for $|\beta| = |q|$, where $\beta = (\beta_1, ..., \beta_m)$, $m = |h(l)|$, a number $\gamma(k)$ corresponds to $l_{\gamma(k)} > 0$, since $q_k = 0$ for $l_k = 0$ and $q_k > 0$ only for $l_k j_k > 0$ and $k > n - \kappa$. That is, $t_{\gamma(1)},...,t_{\gamma(m)}$ are coordinates in $\mathbf{R}^n$ along unit vectors orthogonal to $\partial Q^n_{(l)}$.

Take a sequence U^k of sub-domains $U^k \subset U^{k+1} \subset U$ for each $k \in \mathbf{N}$ so that each $U^k = \bigcup_{l=1}^{m(k)} Q^n_{k,l}$ is the finite union of quadrants $Q^n_{k,l}$, $m(k) \in \mathbf{N}$. We choose them so that each two different quadrants may intersect only by their borders, each U^k satisfies the same conditions as U and

$$(7) \quad \lim_{k\to\infty} dist(U, U^k) = 0.$$

Therefore, Equation (6) can be written for more general domain U also.

For U instead of Q^n we get a face $\partial U_{(l)}$ instead of $\partial Q^n_{(l)}$ and local coordinates $\tau_{\gamma(1)},...,\tau_{\gamma(m)}$ orthogonal to $\partial U_{(l)}$ instead of $t_{\gamma(1)},...,t_{\gamma(m)}$ (see Conditions $(i - iii)$ above).

Thus the sufficient boundary conditions are:

$$(5.1) \quad \Big(\partial^{|\beta|} f(t^{(lj)}) / \partial \tau_{\gamma(1)}^{\beta_1} ... \partial \tau_{\gamma(m)}^{\beta_m} \Big) \Big|_{\partial U_{(lj)}} = \phi_{\beta,(lj)}(t^{(lj)})$$

for $|\beta| = |q|$, where $m = |h(lj)|$, $|j| \le \alpha$, $|(lj)| \ge 1$, $\mathbf{a}_j \ne 0$, $q_k = 0$ for $l_k j_k = 0$, $m_k + q_k + h_k = j_k$,

$h_k = sign(l_k j_k)$, $0 \le q_k \le j_k - 1$ for $k > n - \kappa$; $\phi_{\beta,(l)}(t^{(l)})$ are known functions on $\partial U_{(l)}$, $t^{(l)} \in \partial U_{(l)}$. In the half-space $t_n \ge 0$ only

(5.2) $\partial^\beta f(t)/\partial t_n^\beta|_{t_n=0}$

are necessary for $\beta = |q| < \alpha$ and q as above.

Depending on coefficients of the operator A and the domain U some boundary conditions may be dropped, when the corresponding terms vanish in Formula (6). For example, if $A = \partial^2/\partial t_1 \partial t_2$, $U = U_{1,1}$, $n = 2$, then $\partial f/\partial \nu|_{\partial U_0}$ is not necessary, only the boundary condition $f|_{\partial U}$ is sufficient.

If $U = \mathbf{R}^n$, then no any boundary condition appears. Mention that

(5.3) $\mathcal{F}^0(f(a); u; p; \zeta) = f(a) e^{-u(p,a;\zeta)}$,

which happens in (6), when $a = t^{(l)}$ and $|h(l)| = n$.

Conditions in (5.1) are given on disjoint for different (l) submanifolds $\partial U_{(l)}$ in ∂U and partial derivatives are along orthogonal to them coordinates in $\mathbf{R}^n$, so they are correctly posed.

In $\mathcal{A}_r$ spherical coordinates due to Corollary 4.1 Equation (6) with different values of the parameter ζ gives a system of linear equations relative to unknown functions $\mathsf{S}_{(m)}\mathcal{F}^n(f(t), u; p; \zeta)$, from which $\mathcal{F}^n(f(t), u; p; \zeta)$ can be expressed through a family

$$\left\{\mathsf{S}_{(m)}\mathcal{F}^n(g(t), u; p; \zeta);\ \mathsf{S}_{(m)}\mathcal{F}^{n-|h(l)|}\left(\partial^{|q|} f(t^{(l)})\chi_{\partial Q^n_{(l)}}(t^{(l)})/\partial t_1^{q_1}...\partial t_n^{q_n}, u; p; \zeta\right): (m) \in \mathbf{Z}^n\right\}$$

and polynomials of p, where $\mathbf{Z}$ denotes the ring of integer numbers, where the corresponding term $\mathcal{F}^{n-|h(l)|}$ is zero when $t_j^{(l)} = \pm\infty$ for some j. In the $\mathcal{A}_r$ Cartesian coordinates there are not so well periodicity properties generally, so the family may be infinite. This means that $\mathcal{F}^n(f(t), u; p; \zeta)$ can be expressed in the form:

$$(8) \qquad \mathcal{F}^n(f(t), u; p; \zeta) = \sum_{(m)} \mathsf{P}_{(m)}(p)\mathsf{S}_{(m)}\mathcal{F}^n(g(t), u; p; \zeta)$$

$$+ \sum_{j,(q),(l),|(l)|\ge 1,(m)} \mathsf{P}_{j,(q),(l),(m)}(p)\mathsf{S}_{(m)}\mathcal{F}^{n-|h(lj)|}\left(\partial^{|q|} f(t^{(lj)})\chi_{\partial U_{(lj)}}(t^{(lj)})/\partial t_1^{q_1}...\partial t_n^{q_n}, u; p; \zeta\right),$$

where $\mathsf{P}_{(m)}(p)$ and $\mathsf{P}_{j,(q),(l),(m)}(p)$ are quotients of polynomials of real variables $p_0, p_1, ..., p_n$. The sum in (8) is finite in the $\mathcal{A}_r$ spherical coordinates and may be infinite in the $\mathcal{A}_r$ Cartesian coordinates. To the obtained Equation (8) we apply the theorem about the inversion of the noncommutative multiparameter transform. Thus this gives an expression of f through g as a particular solution of the problem given by $(1, 2, 5.1)$ and it is prescribed by Formulas 6.1(1) and 8.1(1).

For $\mathcal{F}^n(f; u; p; \zeta)$ Conditions 8(1, 2) are satisfied, since $\mathsf{P}_{(m)}(p)$ and $\mathsf{P}_{j,(q),(l),(m)}(p)$ are quotients of polynomials with real, complex or quaternion coefficients and real variables, also G^n and $\mathcal{F}^{n-|h(l)|}$ terms on the right of (6) satisfy them. Thus we have demonstrated the theorem.

28.1. Theorem. *Suppose that $\mathcal{F}^n(f; u; p; \zeta)$ given by the right side of* (8) *satisfies Conditions 8*(3). *Then Problem* $(1, 2, 5.1)$ *has a solution in the class of original functions, when g and $\phi_{\beta,(l)}$ are originals, or in the class of generalized functions, when g and $\phi_{\beta,(l)}$ are generalized functions.*

Mention, that a general solution of $(1, 2)$ is the sum of its particular solution and a general solution of the homogeneous problem $Af = 0$. If $\phi_{\beta,(l)} = \phi^1_{\beta,(l)} + \phi^2_{\beta,(l)}$, $g = g_1 + g_2$, $f = f_1 + f_2$, $Af_j = g_j$ and f_j on ∂U_j satisfies (5.1) with $\phi^j_{\beta,(l)}$, $j = 1, 2$, then $Af = g$ and f on ∂U satisfies Conditions (5.1) with $\phi_{\beta,(l)}$.

28.2. Example. We take the partial differential operator of the second order

$$A = \sum_{h,m=1}^n \mathbf{a}_{h,m}\partial^2/\partial\tau_h\partial\tau_m + \sum_{h=1}^n \alpha_h \partial/\partial\tau_h + \omega,$$

where the quadratic form $a(\tau) := \sum_{h,m} \mathbf{a}_{h,m}\tau_h\tau_m$ is non-degenerate and is not always negative, because otherwise we can consider $-A$. Suppose that $\mathbf{a}_{h,m} = \mathbf{a}_{m,h} \in \mathbf{R}$, $\alpha_h, \tau_h \in \mathbf{R}$ for each $h, m = 1, ..., n$, $\omega \in \mathcal{A}_3$. Then we reduce this form $a(\tau)$ by an invertible $\mathbf{R}$ linear operator C to the sum of squares. Thus

$$(9) \qquad A = \sum_{h=1}^{n} \mathbf{a}_h \partial^2/\partial t_h^2 + \sum_{h=1}^{n} \beta_h \partial/\partial t_h + \omega,$$

where $(t_1, ..., t_n) = (\tau_1, ..., \tau_n)C$ with real $\mathbf{a}_h$ and β_h for each h. If coefficients of A are constant, using a multiplier of the type $\exp(\sum_h \epsilon_h s_h)$ it is possible to reduce this equation to the case so that if $\mathbf{a}_h \neq 0$, then $\beta_h = 0$ (see §3, Chapter 4 in [26]). Then we can simplify the operator with the help of a linear transformation of coordinates and consider that only β_n may be non-zero if $\mathbf{a}_n = 0$. For A with constant coefficients as it is well-known from algebra one can choose a constant invertible real matrix $(c_{h,m})_{h,m=1,...,k}$ corresponding to C so that $\mathbf{a}_h = 1$ for $h \leq k_+$ and $\mathbf{a}_h = -1$ for $h > k_+$, where $0 < k_+ \leq n$. For $k_+ = n$ and $\beta = 0$ the operator is elliptic, for $k_+ = n - 1$ with $\mathbf{a}_n = 0$ and $\beta_n \neq 0$ the operator is parabolic, for $0 < k_+ < n$ and $\beta = 0$ the operator is hyperbolic. Then Equation (6) simplifies:

$$(10) \qquad \mathcal{F}^n(A[f](t), u; p; \zeta) = \sum_{h=1}^{n} \mathbf{a}_h \Big\{ [\mathsf{R}_{e_h}(p)]^2 \mathcal{F}^n(f(t)\chi_{Q^n}(t), u; p; \zeta)$$
$$+ \sum_{l_h \in \{1,2\}; (l) = l_h e_h} (-1)^{l_h} \Big[\mathcal{F}^{n-1}\big(\partial f(t^{(l)})\chi_{\partial Q^n_{(l)}}(t^{(l)})/\partial t_h, u; p; \zeta\big)$$
$$+ [\mathsf{R}_{e_h}(p)]\mathcal{F}^{n-1}(f(t^{(l)})\chi_{\partial Q^n_{(l)}}(t^{(l)}), u; p; \zeta) \Big] \Big\}$$
$$+ \beta_n \Big\{ \mathcal{F}^{n-1;t^{n,2}}\big(f(t^{n,2})\chi_{\partial Q^n_{2e_n}}(t^{n,2}), u; p; \zeta\big) - \mathcal{F}^{n-1;t^{n,1}}\big(f(t^{n,1})\chi_{\partial Q^n_{e_n}}(t^{n,1}), u; p; \zeta\big)$$
$$+ [\mathsf{R}_{e_n}(p)]\mathcal{F}^n\big(f(t)\chi_{Q^n}(t), u; p; \zeta\big) \Big\} + \omega \mathcal{F}^n\big(f(t)\chi_{Q^n}(t), u; p; \zeta\big) = \mathcal{F}^n(g(t), u; p; \zeta)$$

in the $\mathcal{A}_r$ spherical or $\mathcal{A}_r$ Cartesian coordinates, where $e_h = (0, ..., 0, 1, 0, ..., 0) \in \mathbf{R}^n$ with 1 on the h-th place, $\mathsf{S}_0 = I$ is the unit operator, the operators $\mathsf{R}_{e_h}(p)$ are given by Formulas 25(1.1) or 25(1.2) respectively.

We denote by $\delta_S(x)$ the delta function of a continuous piecewise differentiable manifold S in $\mathbf{R}^n$ satisfying conditions $(i - vi)$ so that

$$(\Delta) \qquad \int_{\mathbf{R}^n} \eta(x)\delta_S(x)dx = \int_S \eta(y)\lambda_m(dy)$$

for a continuous integrable function $\eta(x)$ on $\mathbf{R}^n$, where $dim(S) = m < n$, $\lambda_m(dy)$ denotes a volume element on the m dimensional surface S (see Condition (v) above). Thus we can consider a non-commutative multiparameter transform on ∂U for an original f on U given by the formula:

$$(11) \quad \mathcal{F}^{n-1;t'}_{\partial U}(f(t')\chi_{\partial U}(t'), u; p; \zeta) = \mathcal{F}^{n;t}(f(t)\delta_{\partial U}(t), u; p; \zeta).$$

Therefore, terms like $\mathcal{F}^{n-1}$ in (10) correspond to the boundary ∂Q^n. They can be simplified:

$$(12) \qquad \beta_n \Big\{ \mathcal{F}^{n-1;t^{n,2}}\big(f(t^{n,2})\chi_{\partial Q^n_{2e_n}}(t), u; p; \zeta\big) - \mathcal{F}^{n-1;t^{n,1}}\big(f(t^{n,1})\chi_{\partial Q^n_{e_n}}(t), u; p; \zeta\big) \Big\}$$
$$= \mathcal{F}^{n-1;t'}_{\partial Q^n}(\beta(t')f(t')\chi_{\partial Q^n}(t'), u; p; \zeta),$$

where $\beta(t')$ is a piecewise constant function on ∂Q^n equal to β_n on the corresponding faces of Q^n orthogonal to e_n given by condition: either $t_n = \mathsf{a}_n$ or $t_n = b_n$; $\beta(t') = 0$ is zero otherwise.

If $\mathsf{a}_k = -\infty$ or $b_k = +\infty$, then the corresponding term disappears. If $\mathbf{R}^n$ embed into $\mathcal{A}_r$ with $2^{r-1} \le n \le 2^r - 1$ as $\mathbf{R}i_1 \oplus ... \oplus \mathbf{R}i_n$, then this induces the corresponding embedding Θ of Q^n or U into $\mathcal{A}_r$. This permits to make further simplification:

$$(12.1) \quad \sum_{h=1}^n \mathbf{a}_h \Big\{ \sum_{l_h \in \{1,2\}; (l) = l_h e_h} (-1)^{l_h} \Big[[\mathsf{R}_{e_h}(p)] \mathcal{F}^{n-1}(f(t^{(l)})\chi_{\partial Q^n_{(l)}}(t^{(l)}), u; p; \zeta)$$
$$+ \mathcal{F}^{n-1}\Big(\partial f(t^{(l)})\chi_{\partial Q^n_{(l)}}(t^{(l)})/\partial t_h, u; p; \zeta\Big)\Big]\Big\}$$
$$= \mathcal{F}^{n-1}_{\partial Q^n}\Big(a(t')\big(\partial f(t')\chi_{\partial Q^n_0}(t')/\partial\nu\big), u(p, t'; \zeta); p; \zeta\Big)$$
$$+ \mathcal{F}^{n-1}_{\partial Q^n}\Big(\mathsf{P}(t') f(t')\chi_{\partial Q^n_0}(t'), u; p; \zeta\Big),$$

where $\nu = \nu(t')$ denotes a real coordinate along an external unit normal $M(t')$ to $\Theta(\partial U)$ at $\Theta(t')$, so that $M(t')$ is a purely imaginary Cayley-Dickson number, $a(t')$ is a piecewise constant function equal to $\mathbf{a}_h$ for the corresponding t' in the face $\partial Q^n_{l_h e_h}$ with $l_h > 0$; $\mathsf{P}(t', p) := \mathsf{P}(t') := \mathsf{R}_{e_h}(p)$ for $t' \in \partial Q^n_{l_h e_h}$, $h = 1, ..., n$, since $\sin(\psi + \pi) = -\sin(\psi)$ and $\cos(\psi + \pi) = -\cos(\psi)$ for each $\psi \in \mathbf{R}$. Certainly the operator-valued function $\mathsf{P}(t')$ has a piecewise continuous extension $\mathsf{P}(t)$ on Q^n. That is

$$(13) \quad \mathcal{F}^{n-1}_{\partial U}\big(\xi(t') f(t')\chi_{\partial U}(t'), u(p, t'; \zeta); p; \zeta\big)$$
$$:= \int_{\mathbf{R}^n} \xi(t) f(t) \delta_{\partial U}(t) \exp\{-u(p, t; \zeta)\} dt$$

for an integrable operator-valued function $\xi(t)$ so that $[\xi(t) f(t)]$ is an original on U whenever this integral exists. For example, when ξ is a linear combination of shift operators $\mathsf{S}_{(m)}$ with coefficients $\epsilon_{(m)}(t, p)$ such that each $\epsilon_{(m)}(t, p)$ as a function by $t \in U$ for each $p \in W$ and $f(t)$ are originals or f and g are generalized functions. For two quadrants $Q_{m,l}$ and $Q_{m,k}$ intersecting by a common face Υ external normals to it for these quadrants have opposite directions. Thus the corresponding integrals in $\mathcal{F}^{n-1}_{\partial Q_{m,l}}$ and $\mathcal{F}^{n-1}_{\partial Q_{m,k}}$ restricted on Υ summands cancel in $\mathcal{F}^{n-1}_{\partial(Q_{m,l} \cup Q_{m,k})}$.

Using Conditions $(iv - vii)$ and the sequence U^m and quadrants $Q^n_{m,l}$ outlined above we get for a boundary problem on U instead of Q^n the following equation:

$$(14) \quad \mathcal{F}^n(A[f](t), u; p; \zeta) = \Big\{ \sum_{h=1}^n \mathbf{a}_h [\mathsf{R}_{e_h}(p)]^2 \mathcal{F}^n(f(t)\chi_U(t), u; p; \zeta) \Big\} +$$
$$\Big\{ \mathcal{F}^{n-1}_{\partial U}\big([\beta(t') + \mathsf{P}(t', p)] f(t')\chi_{\partial U_0}(t'), u; p; \zeta\big) + \mathcal{F}^{n-1}_{\partial U}(\mathsf{a}(t')(\partial f(t')\chi_{\partial U_0}(t')/\partial\nu), u; p; \zeta) \Big\}$$
$$\mathcal{F}^n\big(\beta_n[\mathsf{R}_n(p)] f(t)\chi_U(t), u; p; \zeta\big) + \omega\mathcal{F}^n\big(f(t)\chi_U(t), u; p; \zeta\big) = \mathcal{F}^n(g(t), u; p; \zeta),$$

where $\mathsf{P}(t', p) := \mathsf{P}(t') := \sum_{h=1}^n \mathbf{a}_h [\mathsf{R}_h(p)](\partial\nu/\partial t_h)$ for each $t' \in \partial U_0$ (see also Stokes' formula in §XIII.3.4 [30] and Formulas (14.2, 14.3) below). Particularly, for the quadrant domain Q^n we have $a(t) = \mathbf{a}_h$ for $t \in \partial Q^n_{l_h e_h}$ with $l_h > 0$, $\beta(t) = \beta_n$ for $t \in \partial Q^n_{l_n e_n}$ with $l_n > 0$ and zero otherwise.

The boundary conditions are:

(14.1) $f(t)|_{\partial U_0} = \phi(t)|_{\partial U_0}, \quad (\partial f(t)/\partial\nu)|_{\partial U_0} = \phi_1(t)|_{\partial U_0}$.

The functions $\mathsf{a}(t)$ and $\beta(t)$ can be calculated from $\{\mathbf{a}_h :\ h\}$ and β_n almost everywhere on ∂U with the help of change of variables from $(t_1, ..., t_n)$ to $(y_1, ..., y_{n-1}, y_n)$, where $(y_1, ..., y_n)$ are local coordinates in ∂U_0 in a neighborhood of a point $t' \in \partial U_0$, $y_n = \nu$, since ∂U_0 is of class C^1. Consider the differential form $\sum_{h=1}^n (-1)^{n-h} \mathbf{a}_h dt_1 \wedge ... \wedge \widehat{dt_h} \wedge ... \wedge dt_n = a dy_1 \wedge ... \wedge dy_{n-1}$ and its external product with $d\nu = \sum_{h=1}^n (\partial \nu / \partial t_h) dt_h$, then

(14.2) $\mathsf{a}(t)|_{\partial U_0} = \sum_{h=1}^n \mathbf{a}_h (\partial\nu/\partial t_h)|_{\partial U_0}$ and

(14.3) $\beta(t)|_{\partial U_0} = \beta_n \chi_{U_{e_n} \cup U_{2e_n}} (\partial \nu/\partial t_n)|_{\partial U_0}$.

This is sufficient for the calculation of $\mathcal{F}^{n-1}_{\partial U}$.

28.3. Inversion procedure in the $\mathcal{A}_r$ spherical coordinates.

When boundary conditions 28(5.1) are specified, this equation 28(6) can be resolved relative to $\mathcal{F}^n(f(t)\chi_U(t), u(p,t;\zeta); p; \zeta)$, particularly, for Equations 28.2(14, 14.1) also. The operators S_{e_j} and T_j of §12 have the periodicity properties: $\mathsf{S}_{e_j}^{4+k} F(p;\zeta) = \mathsf{S}_{e_j}^k F(p;\zeta)$ and $T_j^{4+k} F(p;\zeta) = T_j^k F(p;\zeta)$, $\mathsf{S}_{e_1}^{2+k} F(p;\zeta) = -\mathsf{S}_{e_1}^k F(p;\zeta)$ and $T_1^{2+k} F(p;\zeta) = -T_1^k F(p;\zeta)$ for each positive integer number k and $1 \le j \le 2^r - 1$. We put

(6.1) $\mathbf{F}_j(p;\zeta) := (\mathsf{S}_{e_j}^4 - \mathsf{S}_{e_{j+1}}^4) F(p;\zeta)$ for any $1 \le j \le 2^r - 2$,

(6.2) $\mathbf{F}_{2^r-1}(p;\zeta) := \mathsf{S}_{e_{2^r-1}}^4 F(p;\zeta)$. Then from Formula 28(6) we get the following equations:

$$(6.3) \quad \sum_{|j| \le \alpha} \mathbf{a}_j \Big\{ [p_0 + p_1 T_1]^{j_1} [p_0 + p_1 T_1 + p_2 T_2]^{j_2}$$

$$...[p_0 + p_1 T_1 + ... + p_n T_n]^{j_n} \Big\} \Big|_{p_b = 0 \ \forall b > w} \mathbf{F}_w(p;\zeta) = \Big\{ - \sum_{|j| \le \alpha} \mathbf{a}_j$$

$$\sum_{1 \le |(lj)|;\ m_k + q_k + h_k = j_k;\ 0 \le m_k;\ 0 \le q_k;\ h_k = sign(l_k j_k);\ q_k = 0 \text{ for } l_k j_k = 0, \text{ for each } k = 1,...,n;\ (l) \in \{0,1,2\}^n}$$

$$(-1)^{|(lj)|} \big\{ [p_0 + p_1 T_1]^{m_1} [p_0 + p_1 T_1 + p_2 T_2]^{m_2} ... [p_0 + p_1 T_1 + ... + p_n T_n]^{m_n} \big\} \big|_{p_b = 0 \ \forall b > w}$$

$$\mathcal{F}_w^{n - |h(lj)|} \big(\partial^{|q|} f(t^{(lj)}) \chi_{\partial Q^n_{(lj)}} (t^{(lj)}) / \partial t_1^{q_1} ... \partial t_n^{q_n}, u; p; \zeta \big) \Big\} + \mathbf{G}_w(p;\zeta)$$

for each $w = 1, ..., n$, where $F(p;\zeta) = \mathcal{F}^n(f(t)\chi_{Q^n}(t), u; p; \zeta)$ and $G(p;\zeta) = \mathcal{F}^n(g(t)\chi_{Q^n}(t), u; p; \zeta)$. These equations are resolved for each $w = 1, ..., n$ as it is indicated below. Taking the sum one gets the result

(6.4) $F(p;\zeta) = \mathbf{F}_1(p;\zeta) + ... + \mathbf{F}_n(p;\zeta)$,

since $\Big\{ \big[\sum_{j=1}^{2^r-2} (\mathsf{S}_{e_j}^4 - \mathsf{S}_{e_{j+1}}^4) \big] + \mathsf{S}_{e_{2^r-1}}^4 \Big\} e^{-u(p,t;\zeta)} = \mathsf{S}_{e_1}^4 e^{-u(p,t;\zeta)} = e^{-u(p,t;\zeta)}$.

The analogous procedure is for Equation (14) with the domain U instead of Q^n.

From Equation (6.3) or (14) we get the linear equation:

$$(15) \quad \sum_{(l)} \psi_{(l)} x_{(l)} = \phi,$$

where ϕ is the known function and depends on the parameter ζ, $\psi_{(l)}$ are known coefficients depending on p, $x_{(l)}$ are indeterminates and may depend on ζ, $l_1 = 0, 1$ for $h = 1$, so that $x_{(l)+2e_1} = -x_{(l)}$; $l_h = 0, 1, 2, 3$ for $h > 1$, where $x_{(l)+4e_h} = x_{(l)}$ for each $h > 1$ in accordance with Corollary 4.1, $(l) = (l_1, ..., l_n)$.

Acting on both sides of (6.3) or (14) with the shift operators $T_{(m)}$ (see Formula 25(SO)), where $m_1 = 0, 1$, $m_h = 0, 1, 2, 3$ for each $h > 1$, we get from (15) a system of $2^{1+2(k-1)}$ linear equations with the known functions $\phi_{(m)} := T_{(m)} \phi$ instead of ϕ, $\phi_{(0)} = \phi$:

(15.1) $\sum_{(l)} \psi_{(l)} T_{(m)} x_{(l)} = \phi_{(m)}$ for each (m).

Each such shift of ζ left coefficients $\psi_{(l)}$ intact and $x_{(l)+(m)} = (-1)^\eta x_{(l')}$ with $l'_1 = l_1 + m_1 \ (mod\ 2)$, $l'_h = l_h + m_h \ (mod\ 4)$ for each $h > 1$, where $\eta = 1$ for $l_1 + m_1 - l'_1 = 2$,

$\eta = 2$ otherwise. This system can be reduced, when a minimal additive group $\mathcal{G} := \{(l) : l_1 \ (mod\ 2), \ l_j \ (mod\ 4) \ \forall 2 \le j \le k;$ generated by all (l) with non-zero coefficients in Equation (15)$\}$ is a proper subgroup of $\mathsf{g}_2 \times \mathsf{g}_4^{k-1}$, where $\mathsf{g}_h := \mathbf{Z}/(h\mathbf{Z})$ denotes the finite additive group for $0 < h \in \mathbf{Z}$. Generally the obtained system is non-degenerate for λ_{n+1} almost all $p = (p_0, ..., p_n) \in \mathbf{R}^{n+1}$ or in W, where λ_{n+1} denotes the Lebesgue measure on the real space $\mathbf{R}^{n+1}$.

We consider the non-degenerate operator A with real, complex $\mathbf{C_i}$ or quaternion $\mathbf{H}_{J,K,L}$ coefficients. Certainly in the real and complex cases at each point p, where its determinate $\Delta = \Delta(p)$ is non-zero, a solution can be found by the Cramer's rule.

Generally, the system can be solved by the following algorithm. We can group variables by l_1, l_2,...,l_k. For a given $l_2, ..., l_h$ and $l_1 = 0, 1$ subtracting all other terms from both sides of (15) after an action of $T_{(m)}$ with $m_1 = 0, 1$ and marked m_h for each $h > 1$ we get the system of the form

(16) $\alpha x_1 + \beta x_2 = \mathbf{b}_1,$
$-\beta x_1 + \alpha x_2 = \mathbf{b}_2,$

which generally has a unique solution for λ_{n+1} almost all p:

(17) $x_1 = (\alpha(\alpha^2 + \beta^2)^{-1})\mathsf{b}_1 - (\beta(\alpha^2 + \beta^2)^{-1})\mathsf{b}_2$); $x_2 = (\alpha(\alpha^2 + \beta^2)^{-1})\mathsf{b}_2 + (\beta(\alpha^2 + \beta^2)^{-1})\mathsf{b}_1,$

where $\mathsf{b}_1, \mathsf{b}_2 \in \mathcal{A}_r$, for a given set $(m_2, ...m_n)$.

When l_h are specified for each $1 \le h \le k$ with $h \ne h_0$, where $1 < h_0 \le k$, then the system is of the type:

(18) $ax_1 + bx_2 + cx_3 + dx_4 = \mathsf{b}_1,$
$dx_1 + ax_2 + bx_3 + cx_4 = \mathsf{b}_2,$
$cx_1 + dx_2 + ax_3 + bx_4 = \mathsf{b}_3,$
$bx_1 + cx_2 + dx_3 + ax_4 = \mathsf{b}_4,$

where $a, b, c, d \in \mathbf{R}$ or $\mathbf{C_i}$ or $\mathbf{H}_{J,K,L}$, while $\mathsf{b}_1, \mathsf{b}_2, \mathsf{b}_3, \mathsf{b}_4 \in \mathcal{A}_r$. In the latter case of $\mathbf{H}_{J,K,L}$ it can be solved by the Gauss' exclusion algorithm. In the first two cases of $\mathbf{R}$ or $\mathbf{C_i}$ the solution is:

(19) $x_j = \Delta_j/\Delta$, where
$\Delta = a\xi_1 - d\xi_2 + c\xi_3 - b\xi_4,$
$\Delta_1 = \mathsf{b}_1\xi_1 - \mathsf{b}_2\xi_2 + \mathsf{b}_3\xi_3 - \mathsf{b}_4\xi_4,$
$\Delta_2 = -\mathsf{b}_1\xi_4 + \mathsf{b}_2\xi_1 - \mathsf{b}_3\xi_2 + \mathsf{b}_4\xi_3,$
$\Delta_3 = \mathsf{b}_1\xi_3 - \mathsf{b}_2\xi_4 + \mathsf{b}_3\xi_1 - \mathsf{b}_4\xi_2,$
$\Delta_4 = -\mathsf{b}_1\xi_2 + \mathsf{b}_2\xi_3 - \mathsf{b}_3\xi_4 + \mathsf{b}_4\xi_1,$
$\xi_1 = a^3 + b^2c + cd^2 - ac^2 - 2abd,$
$\xi_2 = a^2b + bc^2 + d^3 - b^2d - 2acd,$
$\xi_3 = ab^2 + c^3 + ad^2 - a^2c - 2bcd,$
$\xi_4 = a^2d + b^3 + c^2d - bd^2 - 2abc.$

Thus on each step either two or four indeterminates are calculated and substituted into the initial linear algebraic system that gives new linear algebraic system with a number of indeterminates less on two or four respectively. May be pairwise resolution on each step is simpler, because the denominator of the type $(\alpha^2 + \beta^2)$ should be λ_{2^r} almost everywhere by $p \in \mathcal{A}_r$ positive (see (6), (14) above). This algorithm acts analogously to the Gauss' algorithm. Finally the last two or four indeterminates remain and they are found with the help of Formulas either (17) or (19) respectively. When for a marked h in (6) or (14) all $l_h = 0 \ (mod\ 2)$ (remains only x_1 for $h = 1$, or remain x_1 and x_3 for $h > 1$) or for some $h > 1$ all $l_h = 0 \ (mod\ 4)$ (remains only x_1) a system of linear equations as in $(15, 15.1)$ simplifies.

Thus a solution of the type prescribed by (8) generally λ_{n+1} almost everywhere by $p \in W$ exists, where W is a domain $W = \{p \in \mathcal{A}_r : a_1 < Re(p) < a_{-1}, \ p_j = 0 \ \forall j > n\}$ of convergence of the noncommutative multiparameter transform, when it is non-void, $2^{r-1} \le n \le 2^r - 1$, $Re(p) = p_0$, $p = p_0 i_0 + ... + p_n i_n$.

This domain W is caused by properties of g and initial conditions on ∂U and by the domain U also. Generally U is worthwhile to choose with its interior $Int(U)$ non-intersecting with a characteristic surface $\phi(x_1, ..., x_n) = 0$, i.e. at each point x of it the condition is satisfied

(CS) $\sum_{|j|=\alpha} \mathbf{a}_j(t(x))(\partial\phi/\partial x_1)^{j_1} ... (\partial\phi/\partial x_n)^{j_n} = 0$
and at least one of the partial derivatives $(\partial\phi/\partial x_k) \ne 0$ is non-zero.

In particular, the boundary problem may be with the right side $g = \varsigma f$ in $(2, 2.1, 14)$, where ς is a real or complex $\mathbf{C_i}$ or quaternion $\mathbf{H}_{J,K,L}$ multiplier, when boundary conditions are non-trivial. In the space either $\mathcal{D}(\mathbf{R}^n, \mathcal{A}_r)$ or $\mathcal{B}(\mathbf{R}^n, \mathcal{A}_r)$ (see §19) a partial differential problem simplifies, because all boundary terms disappear. If $f \in \mathcal{B}(\mathbf{R}^n, \mathcal{A}_r)$, then $\{p \in \mathcal{A}_r : \ Re(p) \ge 0\} \subset W_f$. For $f \in \mathcal{D}(\mathbf{R}^n, \mathcal{A}_r)$ certainly $W_f = \mathcal{A}_r$ (see also §9).

28.4. Examples. Take partial differential equations of the fourth order. In this subsection the noncommutative multiparameter transforms in $\mathcal{A}_r$ spherical coordinates are considered. For

(20) $A = \partial^3/\partial s_1^3 + \sum_{j=2}^n \gamma_j \partial^4/\partial s_j^4$
with constants $\gamma_j \in \mathbf{H}_{J,K,L} \setminus \{0\}$ on the space either $\mathcal{D}(\mathbf{R}^n, \mathcal{A}_r)$ or $\mathcal{B}(\mathbf{R}^n, \mathcal{A}_r)$, where $n \ge 2$, Equation (6) takes the form:

$$(21) \qquad \mathcal{F}^n(A[f](t), u; p; \zeta) =$$

$$\Big\{p_0(p_0^2 + 3(p_1\mathsf{S}_{e_1})^2) + \sum_{j=2}^n \gamma_j (p_j \mathsf{S}_{e_j})^4\Big\}\mathcal{F}^n(f(t), u; p; \zeta) + p_1(3p_0^2 + (p_1\mathsf{S}_{e_1})^2)\mathsf{S}_{e_1}\mathcal{F}^n(f(t), u; p; \zeta)$$

$$= \mathcal{F}^n(g(t), u; p; \zeta)$$

due to Corollary 4.1. In accordance with $(16, 17)$ we get:

(22) $\mathbf{F}_w(p; \zeta) = (\alpha(\alpha^2 + \beta^2)^{-1})\mathbf{G}_w(p; \zeta) - (\beta(\alpha^2 + \beta^2)^{-1})T_1\mathbf{G}_w(p; \zeta))$ for each $w = 1, ..., n$, where $\alpha_w = \alpha = [p_0(p_0^2 - 3p_1^2) + \sum_{j=2}^n \gamma_j p_j^4]\Big|_{p_b=0 \ \forall b>w}$, $\beta_w = \beta = p_1(3p_0^2 - p_1^2)\Big|_{p_b=0 \ \forall b>w}$. From Theorem 6, Corollary 6.1 and Remarks 24 we infer that:

$$(23) \qquad f(t) = (2\pi)^{-n} \int_{\mathbf{R}^n} F(a + p; \zeta) \exp\{u(p, t; \zeta)\} dp_1 ... dp_n$$

supposing that the conditions of Theorem 6 and Corollary 6.1 are satisfied, where $F(p; \zeta) = \mathcal{F}^n(f(t), u; p; \zeta)$.

If on the space either $\mathcal{D}(\mathbf{R}^k, \mathcal{A}_r)$ or $\mathcal{B}(\mathbf{R}^k, \mathcal{A}_r)$ an operator is as follows:

(24) $A = \partial^4/\partial s_1^2 \partial s_2^2 + \sum_{j=3}^n \gamma_j \partial^4/\partial s_j^4$, where $\gamma_j \in \mathbf{H}_{J,K,L} \setminus \{0\}$, where $n \ge 3$, then (6) reads as:

(25) $\mathcal{F}^n(Af(t), u; p; \zeta) = p_2^2(p_0^2 + (p_1\mathsf{S}_{e_1})^2)\mathsf{S}_{e_2}^2\mathcal{F}^n(f(t), u; p; \zeta)$
$+2p_0p_1p_2^2\mathsf{S}_{e_1}\mathsf{S}_{e_2}^2\mathcal{F}^n(f(t)), u; p; \zeta) + \sum_{j=3}^n \gamma_j (p_j\mathsf{S}_{e_j})^4\mathcal{F}^n(f(t)), u; p; \zeta)$
$= \mathcal{F}^n(g(t), u; p; \zeta)$.

If on the same spaces an operator is:

(26) $A = \partial^3/\partial s_1 \partial s_2^2 + \sum_{j=3}^n \gamma_j \partial^4/\partial s_j^4$, where $n \ge 3$, then (6) takes the form:

(27) $\mathcal{F}^n(Af(t), u; p; \zeta) = p_0p_2^2\mathsf{S}_{e_2}^2\mathcal{F}^n(f(t), u; p; \zeta) + p_1p_2^2\mathsf{S}_{e_1}\mathsf{S}_{e_2}^2\mathcal{F}^n(f(t), u; p; \zeta) + \sum_{j=3}^n \gamma_j (p_j\mathsf{S}_{e_j})^4\mathcal{F}^n(f(t), u; p; \zeta) = \mathcal{F}^n(g(t), u; p; \zeta)$.

To find $\mathcal{F}^n(f(t), u; p; \zeta)$ in (25) or (27) after an action of suitable shift operators $T_{(0,2,0,...,0)}$, $T_{(1,0,...,0)}$ and $T_{(1,2,0,...,0)}$ we get the system of linear algebraic equations:

(28) $ax_1 + bx_3 + cx_4 = \mathsf{b}_1$,
$bx_1 + cx_2 + ax_3 = \mathsf{b}_2$,
$ax_2 - cx_3 + bx_4 = \mathsf{b}_3$,
$-cx_1 + bx_2 + ax_4 = \mathsf{b}_4$
with coefficients a, b and c, and Cayley-Dickson numbers on the right side $\mathsf{b}_1, ..., \mathsf{b}_4 \in \mathcal{A}_r$, where $x_1 = \mathbf{F}_w(p;\zeta)$, $x_2 = T_1\mathbf{F}_w(p;\zeta)$, $x_3 = T_2^2\mathbf{F}_w(p;\zeta)$, $x_4 = T_1T_2^2\mathbf{F}_w(p;\zeta)$, $\mathsf{b}_1 = \mathbf{G}_w(p;\zeta) = (\mathcal{F}^n(g(t),u;p;\zeta))_w$, $\mathsf{b}_2 = T_2^2\mathbf{G}_w(p;\zeta)$, $\mathsf{b}_3 = T_1\mathbf{G}_w(p;\zeta)$, $\mathsf{b}_4 = T_1T_2^2\mathbf{G}_w(p;\zeta)$. Coefficients are: $a_w = a = [\sum_{j=3}^n \gamma_j p_j^4]|_{p_b=0\ \forall b>w} \in \mathbf{H}_{J,K,L}$, $b_w = b = p_2^2(p_0^2 - p_1^2) \in \mathbf{R}$, $c_w = c = 2p_0p_1p_2^2|_{p_b=0\ \forall b>w} \in \mathbf{R}$ for A given by (24); $a_w = a = \left[\sum_{j=3}^n \gamma_j p_j^4\right]\Big|_{p_b=0\ \forall b>w} \in \mathbf{H}_{J,K,L}$, $b_w = b = p_0p_2^2\big|_{p_b=0\ \forall b>w} \in \mathbf{R}$, $c_w = c = p_1p_2^2\big|_{p_b=0\ \forall b>w} \in \mathbf{R}$ for A given by (26), $w = 1, ..., n$. If $a = 0$ the system reduces to two systems with two indeterminates (x_1, x_2) and (x_3, x_4) of the type described by (16) with solutions given by Formulas (17). It is seen that these coefficients are non-zero λ_{n+1} almost everywhere on $\mathbf{R}^{n+1}$. Solving this system for $a \neq 0$ we get:

(29) $\mathbf{F}_w(p;\zeta) = a^{-1}\mathsf{b}_1 - [(a^2-b^2+c^2)^2 + 4b^2c^2]^{-1}a^{-1}[(a^2-b^2+c^2)((c^2-b^2)\mathsf{b}_1 + ab\mathsf{b}_2 - 2bc\mathsf{b}_3 + ac\mathsf{b}_4) - 2bc(2bc\mathsf{b}_1 - ac\mathsf{b}_2 + (c^2-b^2)\mathsf{b}_3 + ab\mathsf{b}_4)]$.

Finally Formula (23) provides the expression for f on the corresponding domain W for suitable known function g for which integrals converge. If $\gamma_j > 0$ for each j, then $a > 0$ for each $p_3^2 + ... + p_w^2 > 0$.

For $(21, 24)$ on a bounded domain with given boundary conditions equations will be of an analogous type with a term on the right $\mathcal{F}^n(g(t),u;p;\zeta)$ minus boundary terms appearing in (6) in these particular cases.

For a partial differential equation

$$(30) \qquad \mathbf{a}(t_{n+1})Af(t_1, ..., t_{n+1}) + \partial f(t_1, ..., t_{n+1})/\partial t_{n+1} = g(t_1, ..., t_{n+1})$$

with octonion valued functions f, g, where A is a partial differential operator by variables $t_1, ..., t_n$ of the type given by $(2, 2.1)$ with coefficients independent of $t_1, ..., t_n$, it may be simpler the following procedure. If a domain V is not the entire Euclidean space $\mathbf{R}^{n+1}$ we impose boundary conditions as above in (5.1). Make the noncommutative transform $\mathcal{F}^{n;t_1,...,t_n}$ of both sides of Equation (30), so it takes the form:

$$(31) \qquad \mathbf{a}(t_{n+1})\mathcal{F}^{n;t_1,...,t_n}(Af(t_1, ..., t_{n+1}), u; p; \zeta) + \partial\mathcal{F}^{n;t_1,...,t_n}(f(t_1, ..., t_{n+1}), u; p; \zeta)/\partial t_{n+1}$$
$$= \mathcal{F}^{n;t_1,...,t_n}(g(t_1, ..., t_{n+1}), u; p; \zeta).$$

In the particular case, when
$\mathbf{a}(t_{n+1})\sum_{|j|\le\alpha} \mathbf{a}_j(t_{n+1}) \sum_{0\le k_1\le j_1} \binom{j_1}{k_1} S_{(k_1,j_2,...,j_k)} e^{-u(p,t;\zeta)} = e^{-u(p,t;\zeta)}$
for each t_{n+1}, p, t and ζ, with the help of $(6, 8)$ one can deduce an expression of $F^n(p;\zeta;t_{n+1}) := \mathcal{F}^{n;t_1,...,t_n}(f(t_1, ..., t_{n+1}), u; p; \zeta)$ through $G^n(p;\zeta;t_{n+1}) := \mathcal{F}^{n;t_1,...,t_n}(g(t_1, ..., t_{n+1}), u; p; \zeta)$ and boundary terms in the following form:

$$(32) \qquad \mathbf{b}(p_0, ..., p_n; t_{n+1})F^n(p;\zeta;t_{n+1}) + \partial F^n(p;\zeta;t_{n+1})/\partial t_{n+1} = Q(p_0, ..., p_n; t_{n+1}),$$

where $\mathbf{b}(p_0, ..., p_n; t_{n+1})$ is a real mapping and $Q(p_0, ..., p_n; t_{n+1})$ is an octonion valued function. The latter differential equation by t_{n+1} has a solution analogously to the real case, since t_{n+1} is the real variable, while $\mathbf{R}$ is the center of the Cayley-Dickson algebra $\mathcal{A}_r$. Thus we infer:

$$(33) \qquad F^n(p;\zeta;t_{n+1}) = \exp\left\{-\int_{\tau_0}^{t_{n+1}} \mathbf{b}(p_0, ..., p_n; \xi)d\xi\right\}$$

$$\left\{C_0+\left[\int_{\tau_0}^{t_{n+1}} Q(p_0,...,p_n;\tau)\exp\left\{\int_{\tau_0}^{\tau}\mathbf{b}(p_0,...,p_n;\xi)d\xi\right\}d\tau\right]\right\},$$

since the octonion algebra is alternative and each equation $\mathbf{b}x=\mathbf{c}$ with non-zero $\mathbf{b}$ has the unique solution $x=\mathbf{b}^{-1}\mathbf{c}$, where C_0 is an octonion constant which can be specified by an initial condition. More general partial differential equations as (30), but with $\partial^l f/\partial t^l_{n+1}$, $l\ge 2$, instead of $\partial f/\partial t_{n+1}$ can be considered. Making the inverse transform $(\mathcal{F}^{n;t_1,...,t_n})^{-1}$ of the right side of (33) one gets the particular solution f.

28.5. Integral kernel. We rewrite Equation 28(6) in the form:

(34) $$\mathbf{A}_{\mathcal{S}}\mathcal{F}^n(f\chi_{Q^n},u;p;\zeta)=\mathcal{F}^n(g\chi_{Q^n},u;p;\zeta)-$$

$$\sum_{|j|\le\alpha}\mathbf{a}_j\sum_{1\le|(lj)|,\ 0\le m_k,\ 0\le q_k,\ h_k=sign(l_kj_k),\ m_k+q_k+h_k=j_k;\ q_k=0 \text{ for } l_kj_k=0;\ \forall k=1,...,n;\ (l)\in\{0,1,2\}^n}$$

$$(-1)^{|(lj)|}\mathcal{S}^m\mathcal{F}^{n-|h(lj)|}\left(\partial^{|q|}f(t^{(lj)})/\partial t_1^{q_1}...\partial t_n^{q_n}\right)\chi_{\partial Q^n_{(lj)}}(t^{(lj)}),u;p;\zeta),\ \text{where}$$

(34.1) $\mathcal{S}_k(p):=\mathcal{S}_k:=\mathsf{R}_{e_k}(p)$

in the $\mathcal{A}_r$ spherical or $\mathcal{A}_r$ Cartesian coordinates respectively (see also Formulas 25(1.1, 1.2)), for each $k=1,...,n$,

(34.2) $\mathcal{S}^m(p):=\mathcal{S}^m:=\mathcal{S}_1^{m_1}...\mathcal{S}_n^{m_n}$,

(35) $\mathbf{A}_{\mathcal{S}}:=\sum_{|j|\le\alpha}\mathbf{a}_j\mathcal{S}^j(p)$.

Then we have the integral formula:

(36) $\mathbf{A}_{\mathcal{S}}\mathcal{F}^n(f\chi_{Q^n},u;p;\zeta)=\int_{Q^n}f(t)[\mathbf{A}_{\mathcal{S}}\exp(-u(p,t;\zeta))]dt$

in accordance with 1(7) and 2(4). Due to §28.3 the operator $\mathbf{A}_{\mathcal{S}}$ has the inverse operator for λ_{n+1} almost all $(p_0,...,p_n)$ in $\mathbf{R}^{n+1}$. Practically, its calculation may be cumbersome, but finding for an integral inversion formula its kernel is sufficient. In view of the inversion Theorem 6 or Corollary 6.1 and §§19 and 20 we have

(37) $(2\pi)^{-n}\int_{\mathbf{R}^n}\exp(-u(a+p,t;\zeta))\exp(u(a+p,\tau;\zeta))dp_1...dp_n=\delta(t;\tau)$, where

(38) $[\delta,f)(\tau)=\int_{\mathbf{R}^n}f(t)\delta(t;\tau)dt_1...dt_n=f(\tau)$

at each point $\tau\in\mathbf{R}^n$, where the original $f(\tau)$ satisfies Hölder's condition. That is, the functional $\delta(t;\tau)$ is $\mathcal{A}_r$ linear. Thus the inversion of Equation (36) is:

(39) $$\int_{\mathbf{R}^n}\left(\int_{\mathbf{R}^n}f(t)\chi_{Q^n}(t)\left\{\left[\mathbf{A}_{\mathcal{S}}\exp(-u(p+a,t;\zeta))\right]\xi(p+a,t,\tau;\zeta)\right\}dt\right)dp_1...dp_n=f(\tau),$$

so that

(40) $$\left[\mathbf{A}_{\mathcal{S}}\exp(-u(p+a,t;\zeta))\right]\xi(p+a,t,\tau;\zeta)=(2\pi)^{-n}\exp(-u(p+a,t;\zeta))\exp(-u(p+a,\tau;\zeta)),$$

where the coefficients of $\mathbf{A}_{\mathcal{S}}$ commute with generators i_j of the Cayley-Dickson algebra $\mathcal{A}_r$ for each j. Consider at first the alternative case, i.e. over the Cayley-Dickson algebra $\mathcal{A}_r$ with $r\le 3$.

Let by our definition the adjoint operator $\mathbf{A}^*_{\mathcal{S}}$ be defined by the formula

(41) $\mathbf{A}^*_{\mathcal{S}}\eta(p,t;\zeta)=\sum_{|j|\le\alpha}\mathbf{a}^*_j\mathcal{S}^j\eta^*(p,t;\zeta)$ for any function $\eta:\mathcal{A}_r\times\mathbf{R}^n\times\mathcal{A}_r\to\mathcal{A}_r$, where $t\in\mathbf{R}^n$, p and $\zeta\in\mathcal{A}_r$, $\mathcal{S}^j\eta^*(p,t;\zeta):=\left[\mathcal{S}^j\eta(p,t;\zeta)\right]^*$. Any Cayley-Dickson number $z\in\mathcal{A}_v$ can be written with the help of the iterated exponent (see §3) in $\mathcal{A}_v$ spherical coordinates as

(42) $z=|z|\exp(-u(0,0;\psi))$,

where $v\ge r$, $\psi\in\mathcal{A}_v$, $u\in\mathcal{A}_v$, $Re(\psi)=0$. Certainly the phase shift operator is isometrical:

(43) $|T_1^{k_1}...T_n^{k_n}z|=|z|$

for any $k_1, ..., k_n \in \mathbf{R}$, since $|\exp(-u(0,0;Im(\psi))| = 1$ for each $\psi \in \mathcal{A}_v$, while $T_1^{k_1}...T_n^{k_n} e^{-u(0,0;Im(\psi))} = \exp\{-u(0,0;Im(\psi) - (k_1 i_1 + ... + k_n i_n)\pi/2)\}$ (see §12).

In the $\mathcal{A}_r$ Cartesian coordinates each Cayley-Dickson number can be presented as:

(42.1) $z = |z|\exp(\phi M)$, where $\phi \in \mathbf{R}$ is a real parameter, M is a purely imaginary Cayley-Dickson number (see also §3 in [17, 16]). Therefore, we deduce that

(44) $|\mathbf{A}_{\mathcal{S}} \exp(-u(p+a,t;\zeta))| = \exp(-(p_0+a)s_1 - \zeta_0)|\mathbf{A}_{\mathcal{S}} \exp(-u(Im(p),t;Im(\zeta)))|$,
since $\mathbf{R}$ is the center of the Cayley-Dickson algebra $\mathcal{A}_v$ and $p_0,\ a,\ \zeta_0,\ s_1 \in \mathbf{R},\ \ s_1 = s_1(t)$, where particularly $\mathbf{A}_{\mathcal{S}} 1 := \mathbf{A}_{\mathcal{S}} e^{-u(0,0;\zeta)}|_{\zeta=0}$ (see also Formulas 12(3.1 − 3.7)).

Then expressing ξ from (40) and using Formulas (41, 42, 42.1, 44) we infer, that

(45) $\xi(p,t,\tau;\zeta) = (2\pi)^{-n}[\mathbf{A}^*_{\mathcal{S}} \exp(-u(Im(p),t;Im(\zeta))]$
$[\exp(-u(Im(p),t;Im(\zeta))\exp(u(p,\tau;\zeta))]|\mathbf{A}_{\mathcal{S}} \exp(-u(Im(p),t;Im(\zeta))|^{-2}$,
since $z^{-1} = z^*/|z|^2$ for each non-zero Cayley-Dickson number $z \in \mathcal{A}_v$, $v \ge 1$, where $Im(p) = p_1 i_1 + ... + p_n i_n$, $p = p_0 i_0 + ... + p_n i_n$, $p_0 = Re(p)$.

Generally, for $r \ge 4$, Formula (45) gives the integral kernel $\xi(p,t,\tau;\zeta)$ for any restriction of ξ on the octonion subalgebra $alg_{\mathbf{R}}(N_1,N_2,N_4)$ embedded into $\mathcal{A}_r$. In view of §28.3 ξ is unique and is defined by (45) on each subalgebra $alg_{\mathbf{R}}(N_1,N_2,N_4)$, consequently, Formula (45) expresses ξ by all variables $p, \xi \in \mathcal{A}_r$ and $t,\ \tau \in \mathbf{R}^n$. Applying Formulas (39, 45) and 28.2(Δ) to Equation (34), when Condition 8(3) is satisfied, we deduce, that

$$(46)\quad (f\chi_{Q^n})(\tau) = \int_{\mathbf{R}^n}\left(\int_{\mathbf{R}^n} g(t)\chi_{Q^n}(t)[\exp(-u(p+a,t;\zeta))\xi(p+a,t,\tau;\zeta)]dt\right)dp_1...dp_n -$$

$$\sum_{|j|\le\alpha}\mathbf{a}_j \sum_{1\le|(lj)|,\ 0\le m_k,\ 0\le q_k,\ h_k=sign(l_kj_k);\ m_k+q_k+h_k=j_k;\ q_k=0 \text{ for } l_kj_k=0,\ \forall\ k=1,...,n;\ (l)\in\{0,1,2\}^n} (-1)^{|(lj)|}$$

$$\int_{\mathbf{R}^n}\left(\int_{\partial Q^n_{(lj)}}\left[\partial^{|q|} f(t^{(lj)}/\partial t_1^{q_1}...\partial t_n^{q_n}\right]\left[\{\mathcal{S}^m(p)\exp(-u(p+a,t^{(lj)};\zeta))\}\right.\right.$$

$$\left.\left.\xi(p+a,t^{(lj)},\tau;\zeta)\right]dt^{(lj)}\right)dp_1...dp_n,$$

where $dim_{\mathbf{R}}\partial Q^n_{(lj)} = n - |h(lj)|$, $t^{(lj)} \in \partial Q^n_{(lj)}$ in accordance with §28.1, $\mathcal{S}^m(p)$ is given by Formulas (34.1, 34.2) above.

For simplicity the zero phase parameter $\zeta = 0$ in (46) can be taken. In the particular case $Q^n = \mathbf{R}^n$ all terms with $\int_{\partial Q^n_{(lj)}}$ vanish.

Terms of the form $\int_{\mathbf{R}^n}\left[\{\mathcal{S}^m(p)\exp(-u(p+a,t;\zeta))\}\xi(p+a,t,\tau;\zeta)\right]dp_1...dp_n$ in Formula (46) can be interpreted as left $\mathcal{A}_r$ linear functionals due to Fubini's theorem and §§19 and 20, where $\mathcal{S}^0 = I$.

For the second order operator from (14) one gets:

(47) $\mathbf{A}_{\mathcal{S}} = (\sum_{h=1}^n \mathbf{a}_h[\mathcal{S}_h(p)]^2) + \beta_n\mathcal{S}_n(p) + \omega$ and

$$(48)\qquad (f\chi_U)(t) = \int_{\mathbf{R}^n}\left(\int_{\mathbf{R}^n} g(t)\chi_U(t)\left[\exp(-u(p+a,t;\zeta))\xi(p,t,\tau;\zeta)\right]dt\right)dp_1...dp_n -$$

$$\int_{\mathbf{R}^n}\left(\int_{\partial U_0} f(t')\left[\{(\beta(t') + \mathsf{P}(t',p))\exp(-u(p+a,t;\zeta))\}\xi(p,t',\tau;\zeta)\right]dt'\right)dp_1...dp_n -$$

$$\int_{\mathbf{R}^n}\left(\int_{\partial U_0} a(t')(\partial f(t')/\partial \nu)\Big[\exp(-u(p+a,t;\zeta))\xi(p,t',\tau;\zeta)\Big]dt'\right)dp_1...dp_n.$$

For a calculation of the appearing integrals the generalized Jordan lemma (see §§23 and 24 in [18]) and residues of functions at poles corresponding to zeros $|\mathbf{A}_{\mathcal{S}}\exp(-u(Im(p),t;Im(\zeta)))|$=0 by variables $p_1,...,p_n$ can be used.

Take $g(t)=\delta(y;t)$, where $y\in\mathbf{R}^n$ is a parameter, then

$$(49)\qquad \int_{\mathbf{R}^n}\left(\int_{\mathbf{R}^n}\delta(y;t)\Big[\exp(-u(p+a,t;\zeta))\xi(p+a,t,\tau;\zeta)\Big]dt\right)dp_1...dp_n$$

$$=\int_{\mathbf{R}^n}\Big[\exp(-u(p+a,y;\zeta))\xi(p+a,y,\tau;\zeta)\Big]dp_1...dp_n =: \mathcal{E}(y;\tau)$$

is the fundamental solution in the class of generalized functions, where

(50) $A_t\mathcal{E}(y;t)=\delta(y;t)$,

(51) $\int_{\mathbf{R}^n}\delta(y;t)f(t)dt=f(y)$

for each continuous function $f(t)$ from the space $L^2(\mathbf{R}^n,\mathcal{A}_r)$; A_t is the partial differential operator as above acting by the variables $t=(t_1,...,t_n)$ (see also §§19, 20 and 33-35).

29. The decomposition theorem of partial differential operators over the Cayley-Dickson algebras.

We consider a partial differential operator of order u:

$$(1)\qquad Af(x)=\sum_{|\alpha|\le u}\mathbf{a}_\alpha(x)\partial^\alpha f(x),$$

where $\partial^\alpha f=\partial^{|\alpha|}f(x)/\partial x_0^{\alpha_0}...\partial x_n^{\alpha_n}$, $x=x_0i_0+...x_ni_n$, $x_j\in\mathbf{R}$ for each j, $1\le n=2^r-1$, $\alpha=(\alpha_0,...,\alpha_n)$, $|\alpha|=\alpha_0+...+\alpha_n$, $0\le\alpha_j\in\mathbf{Z}$. By the definition this means that the principal symbol

$$(2)\qquad A_0:=\sum_{|\alpha|=u}\mathbf{a}_\alpha(x)\partial^\alpha$$

has α so that $|\alpha|=u$ and $\mathbf{a}_\alpha(x)\in\mathcal{A}_r$ is not identically zero on a domain U in $\mathcal{A}_r$. As usually $C^k(U,\mathcal{A}_r)$ denotes the space of k times continuously differentiable functions by all real variables $x_0,...,x_n$ on U with values in $\mathcal{A}_r$, while the x-differentiability corresponds to the super-differentiability by the Cayley-Dickson variable x.

Speaking about locally constant or locally differentiable coefficients we shall undermine that a domain U is the union of subdomains U^j satisfying conditions $28(D1,i-vii)$ and $U^j\cap U^k=\partial U^j\cap\partial U^k$ for each $j\ne k$. All coefficients $\mathbf{a}_\alpha$ are either constant or differentiable of the same class on each $Int(U^j)$ with the continuous extensions on U^j. More generally it is up to a C^u or x-differentiable diffeomorphism of U respectively.

If an operator A is of the odd order $u=2s-1$, then an operator E of the even order $u+1=2s$ by variables (t,x) exists so that

(3) $Eg(t,x)|_{t=0}=Ag(0,x)$ for any $g\in C^{u+1}([c,d]\times U,\mathcal{A}_r)$, where $t\in[c,d]\subset\mathbf{R}$, $c\le 0<d$, for example, $Eg(t,x)=\partial(tAg(t,x))/\partial t$.

Therefore, it remains the case of the operator A of the even order $u=2s$. Take $z=z_0i_0+...+z_{2^v-1}i_{2^v-1}\in\mathcal{A}_v$, $z_j\in\mathbf{R}$. Operators depending on a less set $z_{l_1},...,z_{l_n}$ of variables can be considered as restrictions of operators by all variables on spaces of functions constant by variables z_s with $s\notin\{l_1,...,l_n\}$.

Theorem. *Let $A = A_u$ be a partial differential operator of an even order $u = 2s$ with locally constant or variable C^s or x-differentiable on U coefficients $\mathbf{a}_\alpha(x) \in \mathcal{A}_r$ such that it has the form*

(4) $Af = c_{u,1}(B_{u,1}f) + ... + c_{u,k}(B_{u,k}f)$, *where each*

(5) $B_{u,p} = B_{u,p,0} + Q_{u-1,p}$
is a partial differential operator by variables $x_{m_{u,1}+...+m_{u,p-1}+1},...,x_{m_{u,1}+...+m_{u,p}}$ and of the order u, $m_{u,0} = 0$, $c_{u,k}(x) \in \mathcal{A}_r$ for each k, its principal part

(6) $B_{u,p,0} = \sum_{|\alpha|=s} \mathbf{a}_{p,2\alpha}(x)\partial^{2\alpha}$
is elliptic with real coefficients $\mathbf{a}_{p,2\alpha}(x) \ge 0$, either $0 \le r \le 3$ and $f \in C^u(U, \mathcal{A}_r)$, or $r \ge 4$ and $f \in C^u(U, \mathbf{R})$. Then three partial differential operators Υ^s and Υ^s_1 and Q of orders s and p with $p \le u-1$ with locally constant or variable C^s or x-differentiable correspondingly on U coefficients with values in $\mathcal{A}_v$ exist, $r \le v$, such that

(7) $Af = \Upsilon^s(\Upsilon^s_1 f) + Qf$.

Proof. Certainly we have $ordQ_{u-1,p} \le u-1$, $ord(A-A_0) \le u-1$. We choose the following operators:

$$(8)\quad \Upsilon^s f(x) = \sum_{p=1}^{k} \sum_{|\alpha|\le s,\ \alpha_q=0 \forall q<(m_{u,1}+...+m_{u,p-1}+1) \text{ and } q>(m_{u,1}+...+m_{u,p})} (\partial^\alpha f(x))[w_p^*\psi_{p,\alpha}] \text{ and}$$

$$(9)\quad \Upsilon^s_1 f(x) = \sum_{p=1}^{k} \sum_{|\alpha|\le s,\ \alpha_q=0 \forall q<(m_{u,1}+...+m_{u,p-1}+1) \text{ and } q>(m_{u,1}+...+m_{u,p})} (\partial^\alpha f(x))[w_p\psi_{p,\alpha}^*],$$

where $w_p^2 = c_{u,p}$ for all p and $\psi_{p,\alpha}^2(x) = -\mathbf{a}_{p,2\alpha}(x)$ for each p and x, $w_p \in \mathcal{A}_r$, $\psi_{p,\alpha}(x) \in \mathcal{A}_{r,v}$ and $\psi_{p,\alpha}(x)$ is purely imaginary for $\mathbf{a}_{p,2\alpha}(x) > 0$ for all α and x, $Re(w_p Im(\psi_{p,\alpha})) = 0$ for all p and α, $Im(x) = (x - x^*)/2$, $v > r$. Here $\mathcal{A}_{r,v} = \mathcal{A}_v/\mathcal{A}_r$ is the real quotient algebra. The algebra $\mathcal{A}_{r,v}$ has the generators i_{j2^r}, $j = 0, ..., 2^{v-r}-1$. A natural number v so that $2^{v-r}-1 \ge \sum_{p=1}^k \sum_{q=0}^u \binom{m_p+q-1}{q}$ is sufficient, where $\binom{m}{q} = m!/(q!(m-q)!)$ denotes the binomial coefficient, $\binom{m+q-1}{q}$ is the number of different solutions of the equation $\alpha_1 + ... + \alpha_m = q$ in non-negative integers α_j. We have either $\partial^{\alpha+\beta} f \in \mathcal{A}_r$ for $0 \le r \le 3$ or $\partial^{\alpha+\beta} f \in \mathbf{R}$ for $r \ge 4$. Therefore, we can take $\psi_{p,\alpha}(x) \in i_{2^r q}\mathbf{R}$, where $q = q(p, \alpha) \ge 1$, $q(p^1, \alpha^1) \ne q(p, \alpha)$ when $(p, \alpha) \ne (p^1, \alpha^1)$.

Thus Decomposition (7) is valid due to the following. For $b = \partial^{\alpha+\beta} f(z)$ and $\mathbf{l} = i_{2^r p}$ and $w \in \mathcal{A}_r$ one has the identities:

(10) $(b(w\mathbf{l}))(w^*\mathbf{l}) = ((wb)\mathbf{l})(w^*\mathbf{l}) = -w(wb) = -w^2 b$ and

(11) $(((b\mathbf{l})w^*)\mathbf{l})w = (((bw)\mathbf{l})\mathbf{l})w = -(bw)w = -bw^2$ in the considered here cases, since $\mathcal{A}_r$ is alternative for $r \le 3$ while $\mathbf{R}$ is the center of the Cayley-Dickson algebra (see Formulas $(M1, M2)$ in the introduction).

This decomposition of the operator A_{2s} is generally up to a partial differential operator of order not greater, than $(2s-1)$:

(12) $Qf(x) = \sum_{p=1}^k c_{u,p} Q_{u-1,p} +$
$\sum_{|\alpha|\le s, |\beta|\le s; \gamma\le\alpha, \epsilon\le\beta, |\gamma+\epsilon|>0} \left[\prod_{j=0}^{2^v-1} \binom{\alpha_j}{\gamma_j}\binom{\beta_j}{\epsilon_j}\right] \left(\partial^{\alpha+\beta-\gamma-\epsilon} f(x)\right) \left[(\partial^\gamma \eta_\alpha(x))\partial^\epsilon \eta^1_\beta(x)\right],$

where operators Υ^s and Υ^s_1 are already written in accordance with the general form

(13) $\Upsilon^s f(x) = \sum_{|\alpha|\le s}(\partial^\alpha f(x))\eta_\alpha(x)$;

(14) $\Upsilon^s_1 f(x) = \sum_{|\beta|\le s}(\partial^\beta f(x))\eta^1_\beta(x)$.

When A in (3) is with constant coefficients, then the coefficients w_p and $\psi_{p,\alpha}$ for Υ^m and Υ^m_1 can also be chosen constant and $Q - \sum_{p=1}^k c_{u,p}Q_{u-1,p} = 0$.

30. Corollary. *Let suppositions of Theorem 29 be satisfied. Then a change of variables locally affine or variable C^1 or x-differentiable on U correspondingly exists so that the principal part $A_{2,0}$ of A_2 becomes with constant coefficients, when $\mathbf{a}_{p,2\alpha} > 0$ for each p, α and x.*

31. Corollary. *If two operators $E = A_{2s}$ and $A = A_{2s-1}$ are related by Equation 29(3), and A_{2s} is presented in accordance with Formulas 29(4, 5), then three operators Υ^s, Υ^{s-1} and Q of orders s, $s-1$ and $2s-2$ exist so that*

(1) $A_{2s-1} = \Upsilon^s\Upsilon^{s-1} + Q$.

Proof. It remains to verify that $ord(Q) \le 2s-2$ in the case of A_{2s-1}, where $Q = \{\partial(tA_{2s-1})/\partial t - \Upsilon^s\Upsilon^s_1\}|_{t=0}$. Indeed, the form $\lambda(E)$ corresponding to E is of degree $2s-1$ by x and each addendum of degree $2s$ in it is of degree not less than 1 by t, consequently, the product of forms $\lambda(\Upsilon_s)\lambda(\Upsilon^s_1)$ corresponding to Υ^s and Υ^s_1 is also of degree $2s-1$ by x and each addendum of degree $2s$ in it is of degree not less than 1 by t. But the principal parts of $\lambda(E)$ and $\lambda(\Upsilon_s)\lambda(\Upsilon^s_1)$ coincide identically by variables (t,x), hence $ord(\{E - \Upsilon^s\Upsilon^s_1\}|_{t=0}) \le 2s-2$. Let $a(t,x)$ and $h(t,x)$ be coefficients from Υ^s_1 and Υ^s. Using the identities

$a(t,x)\partial_t\partial^\gamma tg(x) = a(t,x)\partial^\gamma g(x)$ and

$h(t,x)\partial^\beta\partial_t[a(t,x)\partial^\gamma g(x)] = h(t,x)\partial^\beta[(\partial_t a(t,x))\partial^\gamma g(x)]$

for any functions $g(x) \in C^{2s-1}$ and $a(t,x) \in C^s$, $ord[(h(t,x)\partial^\beta),(a(t,x)\partial^\gamma)]|_{t=0} \le 2s-2$, where $\partial_t = \partial/\partial t$, $|\beta| \le s-1$, $|\gamma| \le s$, $[A,B] := AB - BA$ denotes the commutator of two operators, we reduce $(\Upsilon^s\Upsilon^s_1 + Q_1)|_{t=0}$ from Formula 29(7) to the form prescribes by equation (1).

32. We consider operators of the form:

(1) $(\Upsilon^k + \beta I_r)f(z) := \{\sum_{0<|\alpha|\le k}(\partial^\alpha f(z)\eta_\alpha(z)\} + f(z)\beta(z)$,
with $\eta_\alpha(z) \in \mathcal{A}_v$, $\alpha = (\alpha_0, ..., \alpha_{2^r-1})$, $0 \le \alpha_j \in \mathbf{N}$
for each j, $|\alpha| = \alpha_0 + ... + \alpha_{2^r-1}$, $\beta I_r f(z) := f(z)\beta$,

$\partial^\alpha f(z) := \partial^{|\alpha|}f(z)/\partial z_0^{\alpha_0}...\partial z_{2^r-1}^{\alpha_{2^r-1}}$, $2 \le r \le v < \infty$, $\beta(z) \in \mathcal{A}_v$, $z_0, ..., z_{2^r-1} \in \mathbf{R}$, $z = z_0 i_0 + ... + z_{2^r-1}i_{2^r-1}$.

Proposition. *The operator $(\Upsilon^k + \beta)^*(\Upsilon^k + \beta)$ is elliptic on the space $C^{2k}(\mathbf{R}^{2^r}, \mathcal{A}_v)$, where $(\Upsilon^k + \beta)^*$ denotes the adjoint operator (i.e. with adjoint coefficients).*

Proof. We establish the identity

(2) $(ay)z^* + (az)y^* = a2Re(yz^*)$

for any $a, y, z \in \mathcal{A}_v$. It is sufficient to prove Equality (2) for any three basic generators of the Cayley-Dickson algebra $\mathcal{A}_v$, since the real field $\mathbf{R}$ is its center, while the multiplication in $\mathcal{A}_v$ is distributive $(a+y)z = az + yz$ and $((\alpha a)(\beta y))(\gamma z^*) = (\alpha\beta\gamma)((ay)z^*)$ for all $\alpha, \beta, \gamma \in \mathbf{R}$ and $a, y, z \in \mathcal{A}_v$. If $a = i_0$, then (2) is evident, since $yz^* + zy^* = yz^* + (yz^*)^* = 2Re(yz^*)$. If $y = i_0$, then $(ay)z^* + (az)y^* = az^* + az = a2Re(z) = a2Re(yz^*)$. Analogously for $z = i_0$.

For three purely imaginary generators i_p, i_s, i_k consider the minimal Cayley-Dickson algebra $\Phi = alg_{\mathbf{R}}(i_p, i_s, i_k)$ over the real field generated by them. If it is associative, then it is isomorphic with either the complex field $\mathbf{C}$ or the quaternion skew field $\mathbf{H}$, so that $(ay)z^* + (az)y^* = a(yz^* + zy*) = a2Re(yz^*)$.

If the algebra Φ is isomorphic with the octonion algebra, then we use Formulas $(M1, M2)$ from the introduction for either $a, y \in \mathbf{H}$ and $z = \mathbf{l}$ or $a, z \in \mathbf{H}$ and $y = \mathbf{l}$. This gives (2) in all cases, since the algebra $alg_{\mathbf{R}}(i_p, i_s)$ with two basic generators i_p and i_s is always associative. Particularly, if $y = i_s \ne z = i_k$, $s,\ k \ge 1$, then the result in (2) is zero.

Using (2) we get more generally, that

(3) $((ay)z^*)b^* + ((az)y^*)b^* = (a2Re(yz^*))b^* = (ab^*)2Re(yz^*)$,
consequently,

(4) $((ay)z^*)b^* + ((az)y^*)b^* + ((by)z^*)a^* + ((bz)y^*)a^* = 4Re(ab^*)Re(yz^*)$

for any Cayley-Dickson numbers $a, b, y, z \in \mathcal{A}_v$. In view of Formulas $(1, 4)$ the form corresponding to the principal symbol of the operator $(\Upsilon^k + \beta)^*(\Upsilon^k + \beta)$ is with real coefficients, of degree $2k$ and non-negative definite, consequently, the operator $(\Upsilon^k + \beta)^*(\Upsilon^k + \beta)$ is elliptic.

33. Fundamental solutions. Let either Y be a real $Y = \mathcal{A}_v$ or complexified $Y = (\mathcal{A}_v)_{\mathbf{C}}$ or quaternionified $Y = (\mathcal{A}_v)_{\mathbf{H}}$ Cayley-Dickson algebra (see §28). Consider the space $\mathcal{B}(\mathbf{R}^n, Y)$ (see §19) supplied with a topology in it is given by the countable family of semi-norms

(1) $p_{\alpha,k}(f) := \sup_{x\in\mathbf{R}^n} |(1 + |x|)^k \partial^\alpha f(x)|$,
where $k = 0, 1, 2, ...$; $\alpha = (\alpha_1, ..., \alpha_n)$, $0 \le \alpha_j \in \mathbf{Z}$. On this space we take the space $\mathcal{B}'(\mathbf{R}^n, Y)_l$ of all Y valued continuous generalized functions (functionals) of the form

(2) $f = f_0 i_0 + ... + f_{2^v-1} i_{2^v-1}$ and $g = g_0 i_0 + ... + g_{2^v-1} i_{2^v-1}$, where f_j and $g_j \in \mathcal{B}'(\mathbf{R}^n, Y)$, with restrictions on $\mathcal{B}(\mathbf{R}^n, \mathbf{R})$ being real or $\mathbf{C_i}$ or $\mathbf{H}_{J,K,L}$ -valued generalized functions $f_0, ..., f_{2^v-1}, g_0, ..., g_{2^v-1}$ respectively. Let $\phi = \phi_0 i_0 + ... + \phi_{2^v-1} i_{2^v-1}$ with $\phi_0, ..., \phi_{2^v-1} \in \mathcal{B}(\mathbf{R}^n, \mathbf{R})$, then

(3) $[f, \phi) = \sum_{k,j=0}^{2^v-1} [f_j, \phi_k) i_k i_j$. We define their convolution as

(4) $[f * g, \phi) = \sum_{j,k=0}^{2^v-1} ([f_j * g_k, \phi) i_j) i_k$ for each $\phi \in \mathcal{B}(\mathbf{R}^n, Y)$. As usually

(5) $(f * g)(x) = f(x - y) * g(y) = f(y) * g(x - y)$
for all $x, y \in \mathbf{R}^n$ due to (4), since the latter equality (5) is satisfied for each pair f_j and g_k. Thus a solution of the equation

(6) $(\Upsilon^s + \beta)f = g$ in $\mathcal{B}(\mathbf{R}^n, Y)$ or in the space $\mathcal{B}'(\mathbf{R}^n, Y)_l$ is:

(7) $f = \mathcal{E}_{\Upsilon^s+\beta} * g$, where $\mathcal{E}_{\Upsilon^s+\beta}$ denotes a fundamental solution of the equation

(8) $(\Upsilon^s + \beta)\mathcal{E}_{\Upsilon+\beta} = \delta$, $(\delta, \phi) = \phi(0)$. The fundamental solution of the equation

(9) $A_0 \mathcal{V} = \delta$ with $A_0 = (\Upsilon^s + \beta)(\Upsilon_1^{s_1} + \beta_1)$
using Equalities $32(2 - 4)$ can be written as the convolution

$$(10) \qquad \mathcal{V} =: \mathcal{V}_{A_0} = \mathcal{E}_{\Upsilon^s+\beta} * \mathcal{E}_{\Upsilon_1^{s_1}+\beta_1}.$$

More generally we can consider the equation

(11) $Af = g$ with $A = A_0 + (\Upsilon_2 + \beta_2)$,

where $A_0 = (\Upsilon + \beta)(\Upsilon_1 + \beta_1)$, Υ, Υ_1, Υ_2 are operators of orders s, s_1 and s_2 respectively given by 32(1) with z-differentiable coefficients. For $\Upsilon_2 + \beta_2 = 0$ this equation was solved above. Suppose now, that the operator $\Upsilon_2 + \beta_2$ is non-zero.

To solve Equation (11) on a domain U one can write it as the system:

(12) $(\Upsilon_1 + \beta_1)f = g_1$, $(\Upsilon + \beta)g_1 = g - (\Upsilon_2 + \beta_2)f$.
Find at first a fundamental solution $\mathcal{V}_A$ of Equation (11) for $g = \delta$. We have:

(13) $f = \mathcal{E}_{\Upsilon_1+\beta_1} * g_1 = \mathcal{E}_{\Upsilon_2+\beta_2} * (g - (\Upsilon + \beta)g_1)$, consequently,

(13.1) $\mathcal{E}_{\Upsilon_1+\beta_1} * g_1 + \mathcal{E}_{\Upsilon_2+\beta_2} * ((\Upsilon + \beta)g_1) = \mathcal{E}_{\Upsilon_2+\beta_2} * g$.
In accordance with $(3 - 5)$ and 32(1) the identity is satisfied: $[\mathcal{E}_{\Upsilon_2+\beta_2} * ((\Upsilon + \beta)g_1), \phi_0) = [(\Upsilon + \beta)(\mathcal{E}_{\Upsilon_2+\beta_2} * g_1), \phi_0)$. Thus (13) is equivalent to

(14) $\mathcal{E}_{\Upsilon_1+\beta_1} * g_1 + (\Upsilon + \beta)(\mathcal{E}_{\Upsilon_2+\beta_2} * g_1) = \mathcal{E}_{\Upsilon_2+\beta_2}$

for $g=\delta$, since $\mathcal{E}_{\Upsilon_2+\beta_2}*\delta=\mathcal{E}_{\Upsilon_2+\beta_2}$.

We consider the Fourier transform F by real variables with the generator $\mathbf{i}$ commuting with i_j for each $j=0,...,2^v-1$ such that

$(F1)$ $(Fg)(y)=\int_{\mathbf{R}^n} e^{-\mathbf{i}(y,x)}g(x)dx_1...dx_n$
for any $g\in L^1(\mathbf{R}^n,\mathcal{A}_v)$, i.e. $\int_{\mathbf{R}^n}|g(x)|dx_1...dx_n<\infty$, where $g:\mathbf{R}^n\to Y$ is an integrable function, $(y,x)=x_1y_1+...+x_ny_n$, $x=(x_1,...,x_n)\in\mathbf{R}^n$, $x_j\in\mathbf{R}$ for every j. The inverse Fourier transform is:

$(F2)$ $(F^{-1}g)(y)=(2\pi)^{-n}\int_{\mathbf{R}^n} e^{\mathbf{i}(y,x)}g(x)dx_1...dx_n$.

For a generalized function f from the space $\mathcal{B}'(\mathbf{R}^n,Y)_l$ its Fourier transform is defined by the formula

$(F3)$ $(Ff,\phi)=(f,F\phi)$, $(F^{-1}f,\phi)=(f,F^{-1}\phi)$.

In view of $(2-5)$ the Fourier transform of (14) gives:

(15) $[F(\mathcal{E}_{\Upsilon_1+\beta_1})][F(g_1)]+\sum_{j=0}^{2^v-1}[F((\Upsilon+\beta)_j\mathcal{E}_{\Upsilon_2+\beta_2})][F(g_1)]i_j=F(\mathcal{E}_{\Upsilon_2+\beta_2})$
for $g=\delta$. With generators $i_0,...,i_{2^v-1},i_0\mathbf{i},...,i_{2^v-1}\mathbf{i}$ the latter equation gives the linear system of 2^{v+1} equations over the real field, or 2^{v+2} equations when $Y=(\mathcal{A}_v)_{\mathbf{H}}$. From it $F(g_1)$ and using the inverse transform F^{-1} a generalized function g_1 can be found, since $F(g)=F(g_0)i_0+...+F(g_{2^v-1})i_{2^v-1}$ and $F^{-1}(g)=F^{-1}(g_0)i_0+...+F^{-1}(g_{2^v-1})i_{2^v-1}$ (see also the Fourier transform of real and complex generalized functions in [5, 29]). Then

(16) $\mathcal{V}_A=\mathcal{E}_{\Upsilon_1+\beta_1}*g_1$ and $f=\mathcal{V}_A*g$ gives the solution of (11), where g_1 was calculated from (15).
Let $\pi^v_r:(\mathcal{A}_v)_{\mathbf{H}}\to(\mathcal{A}_r)_{\mathbf{H}}$ be the $\mathbf{R}$-linear projection operator defined as the sum of projection operators $\pi_0+...+\pi_{2^r-1}$, where $\pi_j:(\mathcal{A}_v)_{\mathbf{H}}\to\mathbf{H}i_j$,

(17) $\pi_j(h)=h_ji_j$, $h=\sum_{j=0}^{2^v-1}h_ji_j$, $h_j\in\mathbf{H}_{J,K,L}$, that gives the corresponding restrictions when $h_j\in\mathbf{C_i}$ or $h_j\in\mathbf{R}$ for $j=0,...,2^r-1$. Indeed, Formulas $2(5,6)$ have the natural extension on $(\mathcal{A}_v)_{\mathbf{H}}$, since the generators J, K and L commute with i_j for each j.

Finally, the restriction from the domain in $\mathcal{A}_v$ onto the initial domain of real variables in the real shadow and the extraction of $\pi^v_r\circ f\in\mathcal{A}_r$ with the help of Formulas $2(5,6)$ gives the reduction of a solution from $\mathcal{A}_v$ to $\mathcal{A}_r$.

Theorems 29, Proposition 32 and Corollaries 30, 31 together with formulas of this section provide the algorithm for subsequent resolution of partial differential equations for $s,s-1,...,2$, because principal parts of operators A_2 on the final step are with constant coefficients. A residue term Q of the first order can be integrated along a path using a non-commutative line integration over the Cayley-Dickson algebra [17, 16].

34. Multiparameter transforms of generalized functions.

If $\phi\in\mathcal{B}(\mathbf{R}^n,Y)$ and $g\in\mathcal{B}'(\mathbf{R}^n,Y)_l$ (see §§19 and 33) we put

(1) $\sum_{j=0}^{2^v-1}[\mathcal{F}^n(g_j;u;p;\zeta),\phi)i_j:=\sum_{j=0}^{2^v-1}[g_j,\mathcal{F}^n(\phi;u;p;\zeta))i_j$ or shortly

(2) $\sum_{j=0}^{2^v-1}[g_je^{-u(p;t;\zeta)},\phi)i_j=\sum_{j=0}^{2^v-1}[g_j,\phi e^{-u(p;t;\zeta)})i_j$.

If the support $supp(g)$ of g is contained in a domain U, then it is sufficient to take a base function ϕ with the restriction $\phi|_U\in\mathcal{B}(U,Y)$ and any $\phi|_{\mathbf{R}^n\setminus U}\in C^\infty$.

34.1. Remark. It is possible to use Theorem 29, Corollaries 30 and 31, Proposition 32 and §33 for solutions of definite differential equations with variable coefficients. For this purpose one can present an operator A as the composition $A=\Upsilon\Upsilon_1+Q$, where $ord(A)=ord(\Upsilon)+ord(\Upsilon_1)$, $ord(Q)\le ord(A)-1$, Υ and Υ_1 are operators with variable coefficients, $\Upsilon^*\Upsilon$ and $\Upsilon_1^*\Upsilon_1$ are elliptic operators with constant coefficients of their principal symbols at least. Then use Formulas $33(1-16)$ to find fundamental solutions $\mathcal{E}_\Upsilon$, $\mathcal{E}_{\Upsilon_1}$ and $\mathcal{E}_A$ or iterate this procedure (see also §35). A generalization of Feynman's formula over the Cayley-Dickson algebras for the second order partial differential operators with the first order addendum Q with variable coefficients from [20] also can be used.

35. Examples.
Let

(1) $Af(t) = \sum_{j=1}^n (\partial^2 f(t)/\partial t_j^2) c_j$
be the operator with constant coefficients $c_j \in \mathcal{A}_r$, $|c_j| = 1$, by the variables $t_1, ..., t_n$, $n \ge 2$. We suppose that c_j are such that the minimal subalgebra $alg_{\mathbf{R}}(c_j, c_k)$ containing c_j and c_k is alternative for each j and k and $|(...(c_1^{1/2} c_2^{1/2})...)c_n^{1/2}| = 1$. Since

(2) $\partial f(t)/\partial t_j = \sum_{k=1}^n (\partial f(t(s))/\partial s_k)(\partial s_k/\partial t_j) = \sum_{k=1}^j \partial f(t(s))/\partial s_k$, the operator A takes the form

(3) $Af = \sum_{j=1}^n (\sum_{1 \le k,b \le j} (\partial^2 f(t(s))/\partial s_k \partial s_b)) c_j$,
where $s_j = t_j + ... + t_n$ for each j. Therefore, by Theorem 12 and Formulas 25(SO) and 28(6) we get:

(4) $\mathcal{F}^n(Af; u; p; \zeta) = \sum_{j=1}^n \{[\mathsf{R}_{e_j}(p)]^2 F^n_u(p; \zeta)\} c_j$ for $u(p, t; \zeta)$ either in $\mathcal{A}_r$ spherical or $\mathcal{A}_r$ Cartesian coordinates with the corresponding operators $\mathsf{R}_{e_j}(p)$ (see also Formulas 25(1.1, 1.2)). On the other hand,

(5) $\mathcal{F}^n(\delta; u; p; \zeta) = e^{-u(p,0;\zeta)} = e^{-u(0,0;\zeta)}$ in accordance with Formula 20(2). The delta function $\delta(t)$ is invariant relative to any invertible linear operator $C : \mathbf{R}^n \to \mathbf{R}^n$ with the determinant $|\det(C)| = 1$, since

$$\int_{\mathbf{R}^n} \delta(Cx)\phi(x)dx = \int_{\mathbf{R}^n} \delta(y)\phi(C^{-1}y)|\det(C)|dy = \phi(C^{-1}0) = \phi(0).$$

Thus

(5) $\mathcal{F}^n(C(Af); u; p; \zeta) = \mathcal{F}^n(Af; u; p; \zeta)$
for any Fundamental solution f, where $Cg(t) := g(Ct)$, $Af = \delta$. If $C : \mathbf{R}^n \to \mathbf{R}^n$ is an invertible linear operator and $\xi = Ct$, $q = Cp$, $\zeta' = C\zeta$, then $t = C^{-1}\xi$, $p = C^{-1}q$ and $\zeta = C^{-1}\zeta'$. In the multiparameter noncommutative transform $\mathcal{F}^n$ there are the corresponding variables (t_j, p_j, ζ_j). This is accomplished in particular for the operator $C(t_1, ..., t_n) = (s_1, ..., s_n)$. The operator C^{-1} transforms the right side of Formula (4), when it is written in the $\mathcal{A}_r$ spherical coordinates, into $\sum_{j=1}^n \{(p_0 + q_j \mathsf{S}_{e_j})^2 F^n_u(q; \zeta)\} c_j$. The Cayley-Dickson number $q = q_0 + q_1 i_1 + ... + q_n i_n$ can be written as $q = q_0 + q_M M$, where $|M| = 1$, M is a purely imaginary Cayley-Dickson number, $q_M \in \mathbf{R}$, $q_M M = q_1 i_1 + ... + q_n i_n$, since $q_0 = Re(q)$. After a suitable automorphism $\theta : \mathcal{A}_r \to \mathcal{A}_r$ we can take $\theta(q) = q_0 + q_M i_1$, so that $\theta(x) = x$ for any real number. The functions $[\sum_{j=1}^n q_j^2 \mathsf{S}^2_{e_j} c_j]$ and $[\sum_{j=1}^n p_j^2 \mathsf{S}^2_{e_j} c_j]$ are even by each variable q_j and p_j respectively. Therefore, we deduce in accordance with (5) and 2(3, 4) and Corollary 6.1 with parameters $p_0 = 0$ and $\zeta = 0$ and $c_j \in \{-1, 1\}$ for each j that

(6) $(\mathcal{F}^n)^{-1}\left(1/\left[\sum_{j=1}^n \left\{\sum_{1 \le k,b \le j} p_k \mathsf{S}_{e_k} p_b \mathsf{S}_{e_b}\right\} c_j\right]; u; y; \zeta\right) = -\left[g, e^{N([y],[q])}\right)$
in the $\mathcal{A}_r$ spherical coordinates, where $g = 1\Big/\left[\sum_{j=1}^n q_j^2 c_j\right]$, or

(6.1) $(\mathcal{F}^n)^{-1}(1/[\sum_{j=1}^n \{p_j^2 \mathsf{S}^2_{e_j}\} c_j]; u; y; \zeta) = -[g, e^{N([y],[p])})$
in the $\mathcal{A}_r$ Cartesian coordinates, where $g = 1/[\sum_{j=1}^n p_j^2 c_j]$, $N = y/|y|$ for $y \ne 0$, $N = i_1$ for $y = 0$, $y = y_1 i_1 + ... + y_n i_n \in \mathcal{A}_r$, $[y] = (y_1, ..., y_n) \in \mathbf{R}^n$, $([y], [q]) = \sum_{j=1}^n y_j q_j$, since $\mathsf{S}^2_{e_k} \cos(\phi + \zeta_k) = \cos(\phi + \zeta_k + \pi) = -\cos(\phi + \zeta_k)$ and $\mathsf{S}^2_{e_k} \sin(\phi + \zeta_k) = \sin(\phi + \zeta_k + \pi) = -\sin(\phi + \zeta_k)$

for each k.

Particularly, we take $c_j = 1$ for each $j = 1, ..., k_+$ and $c_j = -1$ for any $j = k_+ + 1, ..., n$, where $1 \le k_+ \le n$. Thus the inverse Laplace transform for $q_0 = 0$ and $\zeta = 0$ in accordance with Formulas 2(1 − 4) reduces to

(7) $(\mathcal{F}^n)^{-1}\Big(1\Big/\Big[\sum_{j=1}^n \Big\{\sum_{1\le k,b\le j} p_k \mathsf{S}_{e_k} p_b \mathsf{S}_{e_b}\Big\} c_j\Big]; u; y; \zeta\Big) =$
$(2\pi)^{-n} \int_{\mathbf{R}^n} \exp(\mathbf{i}(q_1 y_1 + ... + q_n y_n)) \Big(1\Big/\Big[\sum_{j=1}^{k_+} q_j^2 - \sum_{j=k_++1}^n q_j^2\Big]\Big) dq_1...dq_n$
in the $\mathcal{A}_r$ spherical coordinates and

(7.1) $(\mathcal{F}^n)^{-1}\Big(1\Big/\Big[\sum_{j=1}^n p_j^2 \mathsf{S}_{e_j}^2 c_j\Big]; u; y; \zeta\Big) =$
$(2\pi)^{-n} \int_{\mathbf{R}^n} \exp(\mathbf{i}(p_1 y_1 + ... + p_n y_n)) \Big(1\Big/\Big[\sum_{j=1}^{k_+} p_j^2 - \sum_{j=k_++1}^n p_j^2\Big]\Big) dp_1...dp_n$
in the $\mathcal{A}_r$ Cartesian coordinates,
since for any even function its cosine Fourier transform coincides with the Fourier transform.

The inverse Fourier transform $(F^{-1}g)(x) = (2\pi)^{-n}(Fg)(-x) =: \Psi_n$ of the functions $g = 1/(\sum_{j=1}^n z_j^2)$ for $n \ge 3$ and $\mathcal{P}(1/(\sum_{j=1}^2 z_j^2))$ for $n = 2$ in the class of the generalized functions is known (see [5] and §§9.7 and 11.8 [29]) and gives

(8) $\Psi_n(z_1, ..., z_n) = C_n(\sum_{j=1}^n z_j^2)^{1-n/2}$ for $3 \le n$, where $C_n = -1/[(n-2)\sigma_n]$, $\sigma_n = 4\pi^{n/2}/\Gamma((n/2)-1)$ denotes the surface of the unit sphere in $\mathbf{R}^n$, $\Gamma(x)$ denotes Euler's gamma-function, while

(9) $\Psi_2(z_1, z_2) = C_2 \ln(\sum_{j=1}^2 z_j^2)$ for $n = 2$, where $C_2 = 1/(4\pi)$.
Thus the technique of §2 over the Cayley-Dickson algebra has permitted to get the solution of the Laplace operator.

For the function

(10) $P(x) = \sum_{j=1}^{k_+} x_j^2 - \sum_{j=k_++1}^n x_j^2$ with $1 \le k_+ < n$ the generalized functions $(P(x) + \mathbf{i}0)^\lambda$ and $(P(x) - \mathbf{i}0)^\lambda$ are defined for any $\lambda \in \mathbf{C} = \mathbf{R} \oplus \mathbf{i}\mathbf{R}$ (see Chapter 3 in [5]). The function P^λ has the cone surface $P(z_1, ..., z_n) = 0$ of zeros, so that for the correct definition of generalized functions corresponding to P^λ the generalized functions

(11) $(P(x) + c\mathbf{i}0)^\lambda = \lim_{0<c\epsilon,\epsilon\to 0}(P(x)^2 + \epsilon^2)^{\lambda/2} \exp(\mathbf{i}\lambda arg(P(x) + \mathbf{i}c\epsilon))$
with either $c = -1$ or $c = 1$ were introduced. Therefore, the identity

(12) $F(\Psi_{k_+,n-k_+})(x) = -\Big(\sum_{j=1}^{k_+} x_j^2 - \sum_{j=k_++1}^n x_j^2\Big)\Big[F(\Psi_{k_+,n-k_+})(x)\Big]^2$ or

(13) $F(\Psi) = -1/(P(x) + c\mathbf{i}0)$ follows, where $c = -1$ or $c = 1$.

The inverse Fourier transform in the class of the generalized functions is:

(14) $F^{-1}((P(x)+c\mathbf{i}0)^\lambda)(z_1, ..., z_n) = \exp(-\pi c(n-k_+)\mathbf{i}/2)2^{2\lambda+n}\pi^{n/2}\Gamma(\lambda+n/2)(Q(z_1, ..., z_n) - c\mathbf{i}0)^{-\lambda-n/2)}/[\Gamma(-\lambda)|D|^{1/2}]$
for each $\lambda \in \mathbf{C}$ and $n \ge 3$ (see §IV.2.6 [5]), where $D = \det(g_{j,k})$ denotes a discriminant of the quadratic form $P(x) = \sum_{j,k=1}^n g_{j,k} x_j x_k$, while $Q(y) = \sum_{j,k=1}^n g^{j,k} x_j x_k$ is the dual quadratic form so that $\sum_{k=1}^n g^{j,k} g_{k,l} = \delta_l^j$ for all j, l; $\delta_l^j = 1$ for $j = l$ and $\delta_l^j = 0$ for $j \ne l$. In the particular case of $n = 2$ the inverse Fourier transform is given by the formula:

(15) $F^{-1}((P(x) + c\mathbf{i}0)^{-1})(z_1, z_2) = -4^{-1}|D|^{-1/2} \exp(-\pi c(n-k_+)\mathbf{i}/2) \ln(Q(z_1, ..., z_n) - c\mathbf{i}0)$.
Making the inverse Fourier transform F^{-1} of the function $-1/(P(x) + \mathbf{i}0)$ in this particular case of $\lambda = -1$ we get two complex conjugated fundamental solutions

(16) $\Psi_{k_+,n-k_+}(z_1, ..., z_n) = -\exp(\pi c(n-k_+)\mathbf{i}/2)\Gamma((n/2)-1)(Q(z_1, ..., z_n) + c\mathbf{i}0)^{1-(n/2)}/(4\pi^{n/2})$ for $3 \le n$ and $1 \le k_+ < n$, while

(17) $\Psi_{1,1}(z_1, z_2) = 4^{-1} \exp(\pi c(n-k_+)\mathbf{i}/2) \ln(Q(z_1, z_2) + c\mathbf{i}0)$ for $n = 2$,

where either $c = 1$ or $c = -1$.

Generally for the operator A given by Formula (1) we get $P(x) = P_0(x) + P_i(x)$, where $P_0(x) = \sum_{j=1}^n x_j^2 Re(c_j)$ and $P_i(x) = \sum_{j=1}^n x_j^2 Im(c_j)$ are the real and imaginary parts of P, $Im(z) = z - Re(z)$ for any Cayley-Dickson number z. Take $\mathbf{l} = i_{2^r}$ and consider the form $P(x) + \epsilon c\mathbf{l}$ with $\epsilon \ne 0$ and either $c = 1$ or $c = -1$, then $P(x) + \epsilon c\mathbf{l} \ne 0$ for each $x \in \mathbf{R}^n$. We put

(18) $(P(x) + c\mathbf{l}0)^\lambda = \lim_{0<c\epsilon,\epsilon\to 0}(P(x)^2 + \epsilon^2)^{\lambda/2} \exp(\mathbf{i}\lambda Arg(P(x) + \mathbf{l}c\epsilon))$. Consider $\lambda \in \mathbf{R}$, the generalized function $(P(x)^2 + \epsilon^2)^{\lambda/2} \exp(\mathbf{i}\lambda Arg(P(x) + \mathbf{l}c\epsilon))$ is non-degenerate and for it the Fourier transform is defined. The limit $\lim_{0<c\epsilon,\epsilon\to 0}$ gives by our definition the Fourier transform of $(P(x) + c\mathbf{l}0)^\lambda$. Since

(19) $c_j(\beta_j + \sum_{1\le k\le n, k\ne j} c_j^{-1} c_k\beta_k) = \sum_{j=1}^n c_j\beta_j$
for all $\beta_j \in \mathbf{R}$ and any $1 \le j \le n$ in accordance with the conditions imposed on c_j at the beginning of this section and $\mathbf{i}N_j = N_j\mathbf{i}$ for each j, the Fourier transform with the generator $\mathbf{i}$ can be accomplished subsequently by each variable using Identity (19). The transform $x_j \mapsto (c_j)^{1/2}x_j$ is diagonal and $|(...((c_1^{1/2}c_2^{1/2})...)c_n^{1/2}| = 1$, so we can put $|D| = 1$.

Each Cayley-Dickson number can be presented in the polar form $z = |z|e^{\phi M}$, $\phi \in \mathbf{R}$, $|\phi| \le \pi$, M is a purely imaginary Cayley-Dickson number $|M| = 1$, $Arg(z) = (\phi + 2\pi k)M$ has the countable number of values, $k \in \mathbf{Z}$ (see §3 in [17, 16]). Therefore, we choose the branch $z^{1/2} = |z|^{1/2} \exp((Arg z)/2)$, $|z|^{1/2} > 0$ for $z \ne 0$, with $|Arg(z)| \le \pi$, $Arg(M) = M\pi/2$ for each purely imaginary M with $|M| = 1$.

We treat the iterated integral as in §6, i.e. with the same order of brackets. Taking initially $c_j \in \mathbf{R}$ and considering the complex analytic extension of formulas given above in each complex plane $\mathbf{R} \oplus N_j\mathbf{R}$ by c_j for each j by induction from 1 to n, when c_j is not real in the operator A, $Im(c_j) \in \mathbf{R}N_j$, we get the fundamental solutions for A with the form $(P(x) + c\mathbf{l}0)^\lambda$ instead of $(P(x) + c\mathbf{i}0)^\lambda$ with multipliers $(...(c_1^{c/2}c_2^{c/2})...)c_n^{c/2}$ instead of $\exp(\pi c(n - k_+)\mathbf{i}/2)$ as above and putting $|D| = 1$. Thus

(20) $\Psi(z_1, ..., z_n) = -\Gamma((n/2) - 1)(P^*(z_1, ..., z_n) - c\mathbf{l}0)^{1-(n/2)}[(...(c_1^{c/2}c_2^{c/2})...)c_n^{c/2}]^*/(4\pi^{n/2})$ for $3 \le n$, while

(21) $\Psi(z_1, z_2) = 4^{-1}[c_1^{c/2}c_2^{c/2}]^* Ln(P^*(z_1, z_2) - c\mathbf{l}0)$ for $n = 2$,
since $c_j^* = c_j^{-1}$ for $|c_j| = 1$, $y_jq_j = y_j(c_j^{c/2})^*q_jc_j^{1/2}$, while
$(...(dc_1^{c/2}q_1dc_2^{c/2}q_2)...)dc_n^{c/2}q_n] = dq_1...dq_n[(...(c_1^{c/2}c_2^{c/2})...)c_n^{c/2}]$ and
$|(...(c_1^{c/2}c_2^{c/2})...)c_n^{c/2}| = 1$.

36. Partial differential equations with polynomial real coefficients.
Let

(1) $A = \sum_{|\alpha|\le m} a_\alpha(q)\partial_q^\alpha$, $a_\alpha(q) = \sum_\beta a_{\alpha,\beta}q^\beta$, $q^\beta := q_1^{\beta_1}...q_n^{\beta_n}$, $a_{\alpha,\beta}$ and f have values as in §28, and Af be an original. Using the transform in the $\mathcal{A}_r$ Cartesian coordinates we take $q_j = t_j$ for each j, while using the transform in $\mathcal{A}_r$ spherical coordinates we choose $q_j = s_j(t)$ for each j. Then

(2) $\mathcal{F}^n(Af; u; p; \zeta) = \sum_\beta (-1)^{|\beta|} \mathsf{S}_\beta(p)\partial_p^\beta$
$[\sum_\beta a_{\alpha,\beta}([p_0 + p_1\mathsf{S}_{e_1}]^{\alpha_1} p_2^{\alpha_2}\mathsf{S}_{e_2}^{\alpha_2}...p_n^{\alpha_n}\mathsf{S}_{e_n}^{\alpha_n})]F^n(p;\zeta) = G^n(p;\zeta)$
in the $\mathcal{A}_r$ spherical coordinates and

(2.1) $\mathcal{F}^n(Af; u; p; \zeta) = \sum_\beta (-1)^{|\beta|} \mathsf{S}_\beta(p)\partial_p^\beta$
$(\sum_\beta a_{\alpha,\beta}[p_0 + p_1\mathsf{S}_{e_1}]^{\alpha_1}[p_0 + p_2\mathsf{S}_{e_2}]^{\alpha_2}...[p_0 + p_n\mathsf{S}_{e_n}]^{\alpha_n})F^n(p;\zeta) = G^n(p;\zeta)$
in the $\mathcal{A}_r$ Cartesian coordinates (see Theorems 12 and 13 above). It may happen that the

second differential equation is simpler than the initial one:

(3) $Af = g$.

For example, when coefficients depend only on one variable t_n, then the second differential equation is ordinary and linear.

37. Noncommutative transforms of products and convolutions of functions in the $\mathcal{A}_r$ spherical coordinates.

For any Cayley-Dickson number $z = z_0 i_0 + ... + z_{2^r-1} i_{2^r-1}$ we consider projections

(1) $\theta_j(z) = z_j$, $z_j \in \mathbf{R}$ or $\mathbf{C_i}$ or $\mathbf{H}_{J,K,L}$, $j = 0, ..., 2^r - 1$, $\theta_j(z) = \pi_j(z) i_j^*$,
given by Formulas 2(5, 6) and 33(17). We define the following operators

(2) $\mathcal{R}_{\alpha,j}(F^n(p;\zeta)) := F^n(p_0, (-1)^{\alpha_1} p_1, ..., (-1)^{\alpha_{j+1-\delta_{j,n}}} p_{j+1-\delta_{j,n}}, p_{j+2-\delta_{j,n}},$
$..., p_n; \zeta_0, (-1)^{\alpha_1}\zeta_1 + \pi\alpha_1/2, ..., (-1)^{\alpha_{j+1-\delta_{j,n}}} \zeta_{j+1-\delta_{j,n}} + \pi\alpha_{j+1-\delta_{j,n}}/2, \zeta_{j+2-\delta_{j,n}}, ..., \zeta_n)$
on images F^n, $2^{r-1} \le n \le 2^r - 1$, $j = 0, ..., n$. For α_j and $\beta_j \in \{0, 1\}$ their sum $\alpha_j + \beta_j$ is considered by (mod 2), i.e. in the ring $\mathbf{Z}_2 = \mathbf{Z}/(2\mathbf{Z})$, for two vectors α and $\beta \in \{0,1\}^{2^r-1}$ their sum is considered componentwise in $\mathbf{Z}_2$. Let

$$(3) \qquad \mathcal{F}^n(f; u; p; \zeta) = \sum_{j=0}^{n} \sum_{k=0}^{2^r-1} \theta_j(\mathcal{F}^n(\theta_k(f); u; p; \zeta)) i_k i_j,$$

also $F_j^n(p;\zeta) := \sum_{k=0}^{2^r-1} \theta_j(\mathcal{F}^n(\theta_k(f); u; p; \zeta)) i_k$ for an original f, where $u(p, t; \zeta)$ is given by Formulas 2(1, 2, 2.1). If f is real or $\mathbf{C_i}$ or $\mathbf{H}_{J,K,L}$ -valued, then $F_j^n = \theta_j(F^n)$.

Theorem. *If f and g are two originals, then*

(4) $\mathcal{F}^n(fg; u; p; \zeta) = \sum_{j=0}^n \sum_{\alpha,\beta \in \{0,1\}^n} (-1)^{\alpha_{j+1}(1-\delta_{j+1,n})} (\mathcal{R}_{\alpha,j}(F_j^n(p - q_0; \zeta - \eta))*(\mathcal{R}_{\beta,j}(G_j^n(p + q_0 - p_0; \eta)) i_j$,

(4.1) $\mathcal{F}^n(f * g; u; p; \zeta) = \sum_{j=0}^n \sum_{\alpha,\beta \in \{0,1\}^n} (-1)^{\alpha_{j+1}(1-\delta_{j+1,n})} (\mathcal{R}_{\alpha,j}(F_j^n(p; \zeta - \eta))(\mathcal{R}_{\beta,j}(G_j^n(p; \eta)) i_j$,
whenever $\mathcal{F}^n(fg)$, $\mathcal{F}^n(f)$, $\mathcal{F}^n(g)$ exist, where $1 \le n \le 2^r - 1$, $2 \le r$; $\alpha_k + \beta_k = 1$ (mod 2) for $k \le j$ or $k = j + 1 = n$, $\alpha_k + \beta_k = 0$ (mod 2) for $k = j + 1 < n$ and $\alpha_k = \beta_k = 0$ for $k > j + 1$ in the j-th addendum on the right of Formulas (4, 4.1); the convolution is by $(p_1, ..., p_n)$ in (4), at the same time $q_0 \in \mathbf{R}$ and $\eta \in \mathcal{A}_r$ are fixed.

Proof. The product of two originals can be written in the form:

(5) $f(t)g(t) = \sum_{j=0}^{2^r-1} \sum_{k,l:\ i_k i_l = i_j} \theta_k(f(t)) \theta_l(g(t)) i_j$.

The functions $\theta_k(f)$ and $\theta_l(g)$ are real or $\mathbf{C_i}$ or $\mathbf{H}_{J,K,L}$ valued respectively. The noncommutative transform of fg is:

$$(6) \qquad \mathcal{F}^n(fg)(p;\zeta) = \int_{\mathbf{R}^n} f(t)g(t) \exp(-u(p, t; \zeta)) dt =$$
$$\left\{ \int_{\mathbf{R}^n} (f(t)g(t)) e^{-p_0 s_1} \cos(p_1 s_1 + \zeta_1) i_0 dt \right\} +$$
$$\left\{ \sum_{j=2}^{n-1} \int_{\mathbf{R}^n} (f(t)g(t)) e^{-p_0 s_1} \sin(p_1 s_1 + \zeta_1) ... \sin(p_{j-1} s_{j-1} + \zeta_{j-1}) \cos(p_j s_j + \zeta_j) i_{j-1} dt \right\} +$$
$$\int_{\mathbf{R}^n} (f(t)g(t)) e^{-p_0 s_1} \sin(p_1 s_1 + \zeta_1) ... \sin(p_n s_n + \zeta_n) i_n dt.$$

On the other hand,

$$(7) \qquad \int_{\mathbf{R}^n} f(t)g(t) e^{-p_0 s_1 + \mathbf{i} \sum_{j=1}^k (p_j s_j + \zeta_j)\gamma_j} dt =$$

$$\int_{\mathbf{R}^n}\left(\int_{\mathbf{R}^n} f(t)e^{-(p_0-q_0)s_1+\mathbf{i}\sum_{j=1}^k((p_j-q_j)s_j+\zeta_j-\eta_j)\gamma_j}dt\right)\left(\int_{\mathbf{R}^n} g(t)e^{-q_0s_1+\mathbf{i}\sum_{j=1}^k(q_js_j+\eta_j)\gamma_j}dt\right)dq,$$

where $k = 1, 2, ..., n$, $\gamma_j \in \{-1, 1\}$. Therefore, using Euler's formula $e^{\mathbf{i}\phi} = \cos(\phi) + \mathbf{i}\sin(\phi)$ and the trigonometric formulas $\cos(\phi + \psi) = \cos(\phi)\cos(\psi) - \sin(\phi)\sin(\psi)$, $\sin(\phi + \psi) = \sin(\phi)\cos(\psi) + \cos(\phi)\sin(\psi)$ for all ϕ, $\psi \in \mathbf{R}$, and Formulas $(6, 7)$, we deduce expressions for $\theta_j(\mathcal{F}^n(fg))$. We get the integration by $q_1, ..., q_n$, which gives convolutions by the $p_1, ..., p_n$ variables. Here $q_0 \in \mathbf{R}$ and $\eta \in \mathcal{A}_r$ are any marked numbers. Thus from Formulas $(5-7)$ and $2(1, 2, 2.1, 4)$ we deduce Formula (4).

Moreover, one certainly has

$$(8)\qquad \int_{\mathbf{R}^n}(f * g)(t)e^{-p_0s_1+\mathbf{i}\sum_{j=1}^k(p_js_j+\zeta_j)\gamma_j}dt =$$

$$\left(\int_{\mathbf{R}^n} f(t)e^{-p_0s_1+\mathbf{i}\sum_{j=1}^k(p_js_j+\zeta_j-\eta_j)\gamma_j}dt\right)\left(\int_{\mathbf{R}^n} g(t)e^{-p_0s_1+\mathbf{i}\sum_{j=1}^k(p_js_j+\eta_j)\gamma_j}dt\right)$$

for each $1 \le k \le n$, $\gamma_j \in \{-1, 1\}$, since $s_j(t) = s_j(t-\tau)+s_j(\tau)$ for all $j = 1, ..., n$ and t, $\tau \in \mathbf{R}^n$. Thus from Relations $(6, 8)$ and $2(1, 2, 2.1, 4)$ and Euler's formula one deduces expressions for $\theta_j(\mathcal{F}^n(f * g))$ and Formula (4.1).

38. Moving boundary problem.

Let us consider a boundary problem

(1) $Af = g$ in the half-space $t_n \ge \phi(t_n)$, where $\phi(0) = 0$ and $\phi(t_n) < t_n$ for each $0 \le t_n \in \mathbf{R}$. Suppose that the function $t_n - \phi(t_n) =: \psi(t_n)$ is differentiable and bijective. For example, if $0 < v < 1$ and $\phi(t_n) = vt_n$, then the boundary is moving with the speed v. Make the change of variables $y_n = \psi(t_n)$, $y_1 = t_1$,...,$y_{n-1} = t_{n-1}$, then $t_n = \psi^{-1}(y_n)$ and $dt_n = ds_n = (dt_n/dy_n)dy_n$ and due to Theorem 25 we infer that

$$(2)\qquad \mathcal{F}^n\left(\sum_{|\alpha|\le m}\mathbf{b}_\alpha\partial_s^\alpha\chi_{y_n\ge0}f(t); p; \zeta\right) = \sum_{|\alpha|\le m, 0\le q_n\le\alpha_n-1}\mathbf{b}_\alpha(\delta_{0,\alpha_n}-1)$$

$$(p_0 + \mathsf{S}_{e_1}p_1)^{\alpha_1}p_2^{\alpha_2}...p_{n-1}^{\alpha_{n-1}}p_n^{\alpha_n-q_n-1}\mathsf{S}_{\alpha-\alpha_1e_1-(q_n+1)e_n}\mathcal{F}^{n-1,y^n}(\partial_{t_n}^{q_n}w(y), u(p, (y^n); \zeta); p; \zeta)$$

$$+\sum_{|\alpha|\le m}\mathbf{b}_\alpha(p_0 + \mathsf{S}_{e_1}p_1)^{\alpha_1}p_2^{\alpha_2}...p_n^{\alpha_n}\mathsf{S}_{\alpha-\alpha_1e_1}\mathcal{F}^n(\chi_{y_n\ge0}(y)w(y); p; \zeta) = G^n(p; \zeta)$$

in the $\mathcal{A}_r$ spherical coordinates and

$$(2.1)\qquad \mathcal{F}^n\left(\sum_{|\alpha|\le m}\mathbf{a}_\alpha\partial_t^\alpha\chi_{y_n\ge0}f(t); p; \zeta\right) = \sum_{|\alpha|\le m, 0\le q_n\le\alpha_n-1}\mathbf{a}_\alpha(\delta_{0,\alpha_n}-1)$$

$$(p_0 + \mathsf{S}_{e_1}p_1)^{\alpha_1}(p_0 + p_2\mathsf{S}_{e_2})^{\alpha_2}...(p_0 + p_{n-1}\mathsf{S}_{e_{n-1}})^{\alpha_{n-1}}(p_0 + p_n\mathsf{S}_{e_n})^{\alpha_n-q_n-1}$$

$$\mathcal{F}^{n-1,y^n}(\partial_{t_n}^{q_n}w(y), u(p, (y^n); \zeta); p; \zeta)$$

$$+\sum_{|\alpha|\le m}\mathbf{a}_\alpha(p_0 + \mathsf{S}_{e_1}p_1)^{\alpha_1}(p_0 + p_2\mathsf{S}_{e_2})^{\alpha_2}...(p_0 + p_n\mathsf{S}_{e_n})^{\alpha_n}\mathcal{F}^n(\chi_{y_n\ge0}(y)w(y); p; \zeta) = G^n(p; \zeta)$$

in the $\mathcal{A}_r$ Cartesian coordinates, where $w(y) := f(t(y))(dt_n/dy_n)$.

Expressing $\mathcal{F}^n(\chi_{y_n\ge0}(y)w(y); p; \zeta)$ through $G^n(p; \zeta)$ and the boundary terms $\mathcal{F}^{n-1,y^n}(\partial_{t_n}^{q_n}w(y), u(p, (y^n); \zeta); p; \zeta)$ as in §28.3 and making the inverse transform 8(4) or 8.1(1), or using the integral kernel ξ as in §28.5, one gets a solution $w(y)$ or $f(t) = w(y(t))(dy_n(t_n)/dt_n)$.

39. Partial differential equations with discontinuous coefficients. Consider a domain U and its subdomains $U \supset U_1 \supset ... \supset U_k$ satisfying Conditions $28(D1, D4, i - vii)$ so that coefficients of an operator A (see 28(2)) are constant on $Int(U_k)$ and on $V_1 = U \setminus Int(U_1)$, $V_2 = U_1 \setminus Int(U_2)$,...,$V_k = U_{k-1} \setminus Int(U_k)$ and are allowed to be discontinuous at the common borders $\partial V_j \cap \partial U_j$ for each $j = 1, ..., k$. Each function $f\chi_{U_j}$ is an original on U or a generalized function with the support $supp(f\chi_{U_j}) \subset U_j$ if f is an original or a generalized function on U. Choose operators A^j with constant coefficients on U^j and $A^j|_{Int(V_j)} = 0$, where $U^0 = U$, so that $A|_{U_k} = A^k$,..., $A|_{U_j} = A^j + ... + A^k$,..., $A|_U = A^0 + ... + A^k$. Therefore, in the class of originals or generalized functions on U the problem (see $28(1, 2)$) can be written as

(1) $Af = g$, or

(2) $A^0 f\chi_{V_1} = g\chi_{V_1}$,...,$A^{k-1} f\chi_{V_k} = g\chi_{V_k}$, $A^k f\chi_{U_k} = g\chi_{U_k}$,

since $\chi_{V_1} + ... + \chi_{V_k} + \chi_{U_k} = \chi_U$. Thus the equivalent problem is:

(3) $A^0 f^0 = g^0$, $A^1 f^1 = g^1$,...,$A^k f^k = g^k$

with $f^k = f\chi_{U_k}$, $g^k = g\chi_{U_k}$, also $f^j = f\chi_{V_{j+1}}$, $g^j = g\chi_{V_{j+1}}$ for each $j = 0, ..., k-1$. On ∂U take the boundary condition in accordance with 28(5.1). With any boundary conditions in the class of originals or generalized functions on additional borders $\partial U_j \setminus \partial U$ given in accordance with 28(5.1) a solution f^j on U^j exists, when the corresponding condition 8(3) is satisfied (see Theorems 8 and 28.1).

Each problem $A^j f^j = g^j$ can be considered on U_j, since $supp(g^j) \subset U_j$. Extend f^j by zero on $U \setminus V_j$ for each $0 \le j \le k-1$. When the right side of 28(6) is non-trivial, then f^j is non-trivial. If f^{j-1} is calculated, then the boundary conditions on $\partial U^j \setminus \partial U$ can be chosen in accordance with values of f^{j-1} and its corresponding derivatives $(\partial^\beta f^{j-1}/\partial \nu^\beta)|_{(\partial U^j \setminus \partial U)}$ for some $\beta < ord(A^j)$ in accordance with the operator A^j and the boundary conditions 28(5.1) on the boundary $\partial U^j \setminus \partial U$. Having found f^j for each $j = 0, ..., k$ one gets the solution $f = f^0 + ... + f^k$ on U of Problem (1) with the boundary conditions 28(5.1) on ∂U.

40. Remark. The multiparameter noncommutative transform over the Cayley-Dickson algebras presented above is the natural generalization of the usual complex one-parameter Laplace transform. It opens new opportunities for solving partial differential equations of different types.

It may happen that Theorem 13 is simpler to use, than Theorem 21 for partial differential equations with real variables. Theorem 13 has an advantage that it can be simpler used for partial differential equations of complex and hyper-complex variables, because each pair $(p_l + p_j i_l^* i_j)$ for $l \ne j$ is the complex variable. In these variants boundary conditions may be for $F^k(p; \zeta)$ on a hyperplane $Re(p) = a$ in $\mathcal{A}_r$.

As it was seen above the appearing integrals are by multidimensional domains. For their calculations the Fubini's theorem, residues, Jordan Lemma and tables of known integrals also can be used. Generally in computational mathematics integrals are easier to calculate, than to solve partial differential equations numerically. As a rule iterations of algorithms for integrals converge faster, than iterations of numerical methods for partial differential equations.

Functions with octonion values may be used to resolve systems of partial differential equations. Using conjugations of Cayley-Dickson numbers one gets the transition between operators with coefficients either on the left or on the right of partial derivatives: $[(\partial^\alpha f(x))c_\alpha]^* = c_\alpha^*(\partial^\alpha f(x))^*$, particularly, $(\partial^\alpha f(x))^* = \partial^\alpha f^*(x)$ for $x \in \mathbf{R}^n$, $\partial^\alpha = \partial_x^\alpha$.

Using of Formulas $2(5, 6)$ gives variables $t_j = z_j$ for $z \in \mathcal{A}_r$. So one can consider a class of super-differentiable originals $f(z)$, $z \in V \subset \mathcal{A}_r$. In the class of piecewise on open subsets super-differentiable originals $f(z)$, $z \in V \subset \mathcal{A}_r$, with $t_j = z_j$ for each $j = 1, ..., n$, $n = 2^r - 1$, in the fixed z-representations we get the noncommutative transform for $f(z)\chi_V(z)$ relative to the Cayley-Dickson variable $z \in \mathcal{A}_r$. Therefore, the results given above transfer on this variant also.

Theorem 17 also opens new opportunities to investigate and solve certain types of nonlinear partial differential equations using previous results on spectral theory of functions of operators [21, 22]. For example, analytic functions $q(z)$ in Theorem 17 permit to consider nonlinear operators $q(\sigma)$, where $\sigma f(z) := \sum_{j=0}^{2^r-1}(\partial f(z)/\partial z_j)i_j$. It is planned to study in the next paper.

Partial differential equations with periodic g and f with vector period corresponding to Q^n may be considered also. Certainly others classes of smoothness, for example, Sobolev's or generalized functions can also be considered. It is planned in a next paper to consider this and also problems with boundary conditions as well as with non-constant coefficients in more details.

The technique described above permits to consider partial differential equations of different types and write their solutions in integral forms. If appearing integrals can be calculated in elementary or special of generalized functions, then this gives the explicit formulas in terms of known functions. In conjunction with the line integration over the Cayley-Dickson algebras it permits to solve some types of non linear partial differential equations. The multiparameter Laplace transform over the Cayley-Dickson algebras takes into account the boundary conditions. It naturally means the treatment of systems of partial differential equations due to the multidimensionality of the Cayley-Dickson algebras.

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МНОГОПАРАМЕТРИЧЕСКИЕ ПРЕОБРАЗОВАНИЯ ЛАПЛАСА НАД АЛГЕБРАМИ КЭЛИ-ДИКСОНА И ДИФФЕРЕНЦИАЛЬНЫЕ УРАВНЕНИЯ С ЧАСТНЫМИ ПРОИЗВОДНЫМИ

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Изучаются многомерные некоммутативные преобразования Лапласа над алгебрами Кэли-Диксона. Доказываются теоремы о прямом и обратном преобразованиях Лапласа над алгебрами Кэли-Диксона. Исследуются применения к дифференциальным уравнениям с частными производными, включая эллиптические, параболические и гиперболические. Более того, рассматриваются дифференциальные уравнения с частными производными более высоких порядков с вещественными и комплексными коэффициентами, которые могут быть переменными, с граничными условиями или без них.

Ключевые слова: многомерное некоммутативное преобразование Лапласа, алгебры Кэли-Диксона, дифференциальные уравнения с частными производными, граничные условия