

THE PROLONGATIONS OF A FINSLER METRIC TO THE TANGENT BUNDLE $T^k(M)$ ($k > 1$) OF THE HIGHER ORDER ACCELERATIONS

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An old problem in differential geometry is that of prolongation of a Riemannian structure $g(x)$ on a real n -dimensional C^∞ -manifold $M, x \in M$, to the bundle of k -jets $(J_0^k M, \pi^k, M)$ or, equivalently the tangent bundle $(T^k M, \pi^k, M)$ of the higher order accelerations. The problem belongs to so-called geometry of higher order. It was solved in [18] for $k = 1$ and partially in [19] for $k = 2$. The same problem of prolongation can be considered for a Finslerian structure $F(x, y^{(1)})$. In the paper [15] are given these solutions in the general cases, using the Sasaki-Matsumoto N -lift (for $k = 2$, see [3] and [6]).

But, the terms of Sasaki-Matsumoto prolongation of a Riemannian metric (or Finslerian metric) to $T^k M$ have not the same physical dimensions because these prolongations is not homogeneous on the fibres of the tangent bundle of order k . This is a disadvantage in the study of the geometry of $T^k M$ using the Riemannian metrics determined by these prolongations.

In this paper, only for a Finsler space $F^n = (M, F(x, y^{(1)}))$, we correct this disadvantage introducing a new kind of prolongation $\overset{\circ}{\mathbf{G}}$ of the Finsler metric $g_{ab}(x, y^{(1)}) = \partial^2 F / \partial y^{(1)a} \partial y^{(1)b}$ given by (2.1), which is 0-homogeneous. Some properties of the Riemannian space $(\widetilde{T^k M}, \overset{\circ}{\mathbf{G}})$ are studied. The almost $(k-1)n$ -contact structure $\overset{\circ}{\mathbf{F}}$ from (2.13) is introduced. It has the property of homogeneity and $(\overset{\circ}{\mathbf{G}}, \overset{\circ}{\mathbf{F}})$ is a metrical almost $(k-1)n$ -contact structure on $T^k M$. It depend only on the fundamental function $F(x, y^{(1)})$ of the Finsler space F^n . The space $(\widetilde{T^k M}, \overset{\circ}{\mathbf{G}}, \overset{\circ}{\mathbf{F}})$ is **the geometrical model** of the Finsler space $F^n = (M, F(x, y^{(1)}))$.

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The Sasaki-Matsumoto N -lift of a Finsler metric

Let M be a real n -dimensional C^∞ -manifold and $(T^k M, \pi^k, M)$ its tangent bundle of order k (or k -jet bundle, or tangent bundle of the higher order accelerations).

Let us consider the Finsler space $F^n = (M, F)$ with the fundamental function $F(x, y^{(1)})$, $F : T^1 M \rightarrow R$, and the fundamental tensor $g_{ab}(x, y^{(1)})$ on $T^1 M$ given by

$$g_{ab}(x, y^{(1)}) = \frac{1}{2} \frac{\partial^2 F}{\partial y^{(1)a} \partial y^{(1)b}}, \quad (1.1)$$

where $g_{ab}(x, y^{(1)})$ is positively defined on $T^1 M$.

The indices a, b, \dots run over set $\{1, 2, \dots, n\}$ and Einstein convention of summing is adopted all over this work.

Let $\gamma_{bc}^a(x, y^{(1)})$ be the formal Christoffel symbols of the $g_{ab}(x, y^{(1)})$, *i.e.* :

$$\gamma_{bc}^a(x, y^{(1)}) = \frac{1}{2} g^{ad} \left(\frac{\partial g_{bd}}{\partial x^c} + \frac{\partial g_{dc}}{\partial x^b} - \frac{\partial g_{bc}}{\partial x^d} \right). \quad (1.2)$$

Then, the canonical semispray of F^n is given by

$$\frac{d^2x^a}{dt^2} + 2G_{(1)}^a \left(x, \frac{dx}{dt} \right) = 0, \tag{1.3}$$

where

$$G_{(1)}^a = \frac{1}{2} \gamma_{bc}^a (x, y^{(1)}) y^{(1)b} y^{(1)c}. \tag{1.3'}$$

The canonical nonlinear connection (determined only by the function F of the Finsler space F^n) is the Cartan nonlinear connection with the coefficients

$$G_b^a (x, y^{(1)}) = \frac{\partial G_{(1)}^a}{\partial y^{(1)b}}. \tag{1.4}$$

Then, on the domain of chart $(\pi^k)^{-1}(U) \subset T^k M, U \subset M$, we can consider the functions

$$\begin{aligned} F^* (x, y^{(1)}, \dots, y^{(k)}) &= (F \circ \pi_1^k) (x, y^{(1)}, \dots, y^{(k)}), \\ g_{ab}^* (x, y^{(1)}, \dots, y^{(k)}) &= (g_{ab} \circ \pi_1^k) (x, y^{(1)}, \dots, y^{(k)}), \\ &\forall (x, y^{(1)}, \dots, y^{(k)}) (U), \end{aligned}$$

where $\pi_1^k : T^k M \rightarrow TM, \pi_1^k (x, y^{(1)}, \dots, y^{(k)}) = (x, y^{(1)})$ is the natural projection. For simplicity, F^* and g_{ab}^* will be denote by the same letters F and g_{ab} .

We have

1⁰. The canonical nonlinear connection N on $\widetilde{T^k M} = T^k M \setminus \{0\}$ has the dual coefficients

$$\begin{aligned} M_1^a_b &= G^a_b, \\ M_2^a_b &= \frac{1}{2} \left(C M_1^a_b + M_1^a_c M_1^c_b \right), \\ &\dots\dots\dots \\ M_k^a_b &= \frac{1}{k} \left(C M_{k-1}^a_b + M_1^a_c M_{k-1}^c_b \right), \end{aligned} \tag{1.5}$$

where C is the operator

$$C = y^{(1)a} \frac{\partial}{\partial x^a} + 2y^{(2)a} \frac{\partial}{\partial y^{(1)a}} + \dots + ky^{(k)a} \frac{\partial}{\partial y^{(k-1)a}} \tag{1.6}$$

2⁰. The Liouville d -vector field $z^{(k)}$ corresponding to the canonical nonlinear connection N is given by

$$kz^{(k)a} = ky^{(k)a} + (k - 1) y^{(k-1)b} M_1^a_b + \dots + y^{(1)b} M_{k-1}^a_b. \tag{1.7}$$

3⁰. The following Lagrangian

$$L (x, y^{(1)}, \dots, y^{(k)}) = g_{ab} (x, y^{(1)}) z^{(k)a} z^{(k)b}, \tag{1.8}$$

is a regular Lagrangian on $\widetilde{T^k M}$, determined only by $F (x, y^{(1)})$ because g_{ab} and $z^{(k)}$ have this property.

4⁰. Its fundamental tensor field coincide with the fundamental tensor field on Finsler space F^n , namely on $\widetilde{T^k M}$ we have

$$\frac{1}{2} \frac{\partial^2 L}{\partial z^{(k)a} \partial z^{(k)b}} = g_{ab}(x, y^{(1)}). \tag{1.9}$$

5⁰. N determines the direct decomposition

$$T_u T^k M = N_0(u) \oplus N_1(u) \oplus \dots \oplus N_{k-1}(u) \oplus V_k(u), \quad \forall u \in T^k M. \tag{1.10}$$

6⁰. The adapted cobasis $\{dx^a, \delta y^{(1)a}, \dots, \delta y^{(k)a}\}$ and the adapted basis $\left\{ \frac{\delta}{\delta x^a}, \frac{\delta}{\delta y^{(1)a}}, \dots, \frac{\delta}{\delta y^{(k-1)a}}, \frac{\delta}{\delta y^{(k)a}} \right\}$ to N are depending only on fundamental function $F(x, y^{(1)})$ of Finsler space F^n , where

$$\begin{aligned} \delta y^{(1)a} &= dy^{(1)a} + M_{1c}^a dx^c, \\ \delta y^{(2)a} &= dy^{(2)a} + M_{1c}^a dy^{(1)c} + M_{2c}^a dx^c, \\ &\dots\dots\dots \\ \delta y^{(k)a} &= dy^{(k)a} + M_{1c}^a dy^{(k-1)c} + \dots + M_{k-1c}^a dy^{(1)c} + M_k^a dx^c, \end{aligned} \tag{1.11}$$

and

$$\begin{aligned} \frac{\delta}{\delta x^a} &= \frac{\partial}{\partial x^a} - N_{1c}^a \frac{\partial}{\partial y^{(1)c}} - N_{2c}^a \frac{\partial}{\partial y^{(2)c}} - \dots - N_{kc}^a \frac{\partial}{\partial y^{(k)c}}, \\ \frac{\delta}{\delta y^{(1)a}} &= \frac{\partial}{\partial y^{(1)a}} - N_{1c}^a \frac{\partial}{\partial y^{(2)c}} - \dots - N_{k-1c}^a \frac{\partial}{\partial y^{(k)c}}, \\ &\dots\dots\dots \\ \frac{\delta}{\delta y^{(k-1)a}} &= \frac{\partial}{\partial y^{(k-1)a}} - N_{1c}^a \frac{\partial}{\partial y^{(k)c}}. \end{aligned} \tag{1.11'}$$

We know that

$$\begin{aligned} N_{1b}^a &= M_{1b}^a, N_{2b}^a = M_{2b}^a - M_{1b}^c M_{1c}^a, \dots, \\ N_{kb}^a &= M_{kb}^a - M_{1b}^c N_{k-1c}^a - \dots - M_{k-2b}^c N_{2c}^a - M_{k-1b}^c N_{1c}^a, \end{aligned} \tag{1.12}$$

and conversely

$$\begin{aligned} M_{1b}^a &= N_{1b}^a, M_{2b}^a = N_{2b}^a + N_{1c}^a M_{1c}^b, \dots, \\ M_{kb}^a &= N_{kb}^a + N_{k-1c}^a M_{1c}^b + \dots + N_{2c}^a M_{k-2c}^b + N_{1c}^a M_{k-1c}^b. \end{aligned} \tag{1.12'}$$

Then, the Sasaki-Matsumoto N -lift of $g_{ab}(x, y^{(1)})$ to $T^k M$ is defined by

$$\mathbf{G}(u) = g_{ab}(x, y^{(1)}) dx^a \otimes dx^b + \sum_{\beta=1}^k g_{ab}(x, y^{(1)}) \delta y^{(\beta)a} \otimes \delta y^{(\beta)b}, \quad \forall u \in \widetilde{T^k M}. \tag{1.13}$$

The following properties hold:

- 7⁰. \mathbf{G} is globally defined on $T^k M$.
- 8⁰. \mathbf{G} is a Riemannian structure on $T^k M$ determined only by the Finsler space F^n .
- 9⁰. \mathbf{G} is not homogeneous on the fibres of $T^k M$.

Namely, for the homothety $h_t : (x, y^{(1)}, \dots, y^{(k)}) \rightarrow (x, ty^{(1)}, \dots, t^k y^{(k)})$, $\forall t \in R_*^+$, we get

$$(G \circ h_t)(u) = g_{ab}(x, y^{(1)}) dx^a \otimes dx^b + \sum_{\beta=1}^k t^{2\beta} g_{ab}(x, y^{(1)}) \delta y^{(\beta)a} \otimes \delta y^{(\beta)b} \neq G(u).$$

Let us consider the $\mathcal{F}(T^k M)$ – linear mapping $\mathbf{F} : \chi(T^k M) \rightarrow \chi(T^k M)$ given in the adapted basis (1.11') by

$$\begin{aligned} \mathbf{F} \left(\frac{\delta}{\delta x^a} \right) &= -\frac{\partial}{\partial y^{(k)a}}, \\ \mathbf{F} \left(\frac{\delta}{\delta y^{(1)a}} \right) &= \dots = \mathbf{F} \left(\frac{\delta}{\delta y^{(k-1)a}} \right) = 0, \\ \mathbf{F} \frac{\partial}{\partial y^{(k)a}} &= \frac{\delta}{\delta x^a}. \end{aligned} \tag{1.14}$$

It follows that:

- 10⁰. \mathbf{F} is globally defined on $T^k M$ and it is a d – tensor field of type $(1, 1)$.
- 11⁰. \mathbf{F} is an $(k - 1)n$ –contact structure : $\mathbf{F}^3 + \mathbf{F} = 0$.
- 12⁰. \mathbf{F} depend only on the fundamental function $F(x, y^{(1)})$ of Finsler space F^n .
- 13⁰. The pair (\mathbf{G}, \mathbf{F}) is a Riemannian almost $(k - 1)n$ –contact structure on $T^k M$:

$$\mathbf{G}(\mathbf{F}X, Y) = -\mathbf{G}(X, \mathbf{F}Y), \forall X, Y \in \chi(T^k M).$$

Consequently, we get

Theorem 1.1 *The space $(T^k M, \mathbf{G}, \mathbf{F})$ is a Riemannian almost $(k - 1)n$ –contact space depending only on the fundamental function $F(x, y^{(1)})$ of the Finsler space $F^n = (M, F)$.*

The previous space, called "the geometrical model on $T^k M$ of the Finsler space" (M, F) is important in the study of the geometry of the initial Finsler space $F^n = (M, F)$.

The homogeneous prolongation to $T^k M$ of a Finsler metric

We define a new prolongation $\overset{\circ}{G}$ on $T^k M$ of the fundamental tensor field $g_{ab}(x, y^{(1)})$ of a Finsler space $F^n = (M, F)$, which satisfies the following conditions:

- 1⁰. $\overset{\circ}{G}$ is 0– homogeneous with respect to $y^{(1)a}, y^{(2)a}, \dots,$ and $y^{(k)a}$.
- 2⁰. It depends only on the fundamental function $F(x, y^{(1)})$.
- 3⁰. In the mechanical meaning the terms of $\overset{\circ}{G}$ have the same physical dimensions.

Definition 2.1. *We call the homogeneous prolongation to $T^k M$ of the fundamental tensor field $g_{ab}(x, y^{(1)})$ of a Finsler space $F^n = (M, F)$, the following tensor field on $T^k M$:*

$$\overset{\circ}{G}(u) = g_{ab}(x, y^{(1)}) dx^a \otimes dy^b + \sum_{\beta=1}^k \frac{\mathbf{a}^{2\beta}}{\|y^{(1)}\|^{2\beta}} g_{ab}(x, y^{(1)}) \delta y^{(\beta)a} \otimes \delta y^{(\beta)b}, \forall u \in \widetilde{T^k M}, \tag{2.1}$$

where $\mathbf{a} > \mathbf{0}$ is a constant imposed by application in order to preserve the physical dimension of the components of $\overset{\circ}{G}$, and $\|y^{(1)}\|^2$ is the square of the norm of the first Liouville vector field

$$\|y^{(1)}\|^2 = g_{ab}(x, y^{(1)}) y^{(1)a} y^{(1)b}. \tag{2.2}$$

We get, without difficulties:

Theorem 2.1. 1. The pair $(\widetilde{T^k M}, \overset{\circ}{G})$ is a Riemann space.

2. $\overset{\circ}{G}$ is a 0-homogeneous tensor field with respect to $y^{(\beta)a}$, $(\beta = 1, \dots, k)$.

3. $\overset{\circ}{G}$ depends only on the fundamental function $F(x, y^{(1)})$ of the Finsler space F^n .

4. The distributions N_0, N_1, \dots, N_{k-1} and V_k are orthogonal, in pairs, with respect to $\overset{\circ}{G}$.

We can write $\overset{\circ}{G}$ in the form

$$\overset{\circ}{G} = \overset{\circ}{G}^H + \overset{\circ}{G}^{V_1} + \dots + \overset{\circ}{G}^{V_k}, \tag{2.3}$$

where

$$\overset{\circ}{G}^H = g_{ab}(x, y^{(1)}) dx^a \otimes dx^b, \overset{\circ}{G}^{V_\beta} = g_{ab}(x, y^{(1)}) dy^{(\beta)a} \otimes dy^{(\beta)b} \tag{2.4}$$

and

$$g_{ab}^{(\beta)}(x, y^{(1)}) = \frac{\mathbf{a}^{2\beta}}{\|y^{(1)}\|^{2\beta}} g_{ab}(x, y^{(1)}), \quad (\beta = 1, \dots, k). \tag{2.5}$$

As usually, let us denote

$$\partial_a = \frac{\partial}{\partial x^a}, \dot{\partial}_{1a} = \frac{\partial}{\partial y^{(1)a}}, \dots, \dot{\partial}_{ka} = \frac{\partial}{\partial y^{(k)a}},$$

and from now on we denote the adapted basis (1.11') by

$$\{\delta_a, \delta_{1a}, \dots, \delta_{(k-1)a}, \delta_{ka}\}.$$

In order to study the geometry of Riemann space $(\widetilde{T^k M}, \overset{\circ}{G})$, we can apply the theory of the (h, v_1, \dots, v_k) -Riemannian metric given by author in [5] (for $k = 2$, see [2], [4]).

A linear connection D on $T^k M$ is called a metrical N -linear connection with respect to $\overset{\circ}{G}$ if $D_X \overset{\circ}{G} = 0, \forall X \in \chi(T^k M)$ and it preserves by paralelism the horizontal and vertical distributions $N_0, N_1, \dots, N_{k-1}, V_k$.

We can easily prove the existence of the metrical N -linear connections in the adapted basis. To this aim we represent a linear connection D in the adapted basis in the following form:

$$\begin{aligned} D_{\delta_c} \delta_b &= \overset{0}{(00)} L^a{}_{bc} \delta_a + \sum_{\beta=1}^k \overset{\beta}{(00)} L^a{}_{bc} \delta_{\beta a}, \\ D_{\delta_c} \delta_{\gamma b} &= \overset{0}{(\gamma 0)} L^a{}_{bc} \delta_a + \sum_{\beta=1}^k \overset{\beta}{(\gamma 0)} L^a{}_{bc} \delta_{\beta a}, \quad (\gamma = 1, \dots, k; \delta_{ka} = \dot{\partial}_{ka}), \\ D_{\delta_{1c}} \delta_b &= \overset{0}{(01)} C^a{}_{bc} \delta_a + \sum_{\beta=1}^k \overset{\beta}{(01)} C^a{}_{bc} \delta_{\beta a}, \\ D_{\delta_{1c}} \delta_{\gamma b} &= \overset{0}{(\gamma 1)} C^a{}_{bc} \delta_a + \sum_{\beta=1}^k \overset{\beta}{(\gamma 1)} C^a{}_{bc} \delta_{\beta a}, \quad (\gamma = 1, \dots, k; \delta_{ka} = \dot{\partial}_{ka}), \\ &\dots\dots\dots \\ D_{\delta_{kc}} \delta_b &= \overset{0}{(0k)} C^a{}_{bc} \delta_a + \sum_{\beta=1}^k \overset{\beta}{(0k)} C^a{}_{bc} \delta_{\beta a}, \\ D_{\delta_{kc}} \delta_{\gamma b} &= \overset{0}{(\gamma k)} C^a{}_{bc} \delta_a + \sum_{\beta=1}^k \overset{\beta}{(\gamma k)} C^a{}_{bc} \delta_{\beta a}, \quad (\gamma = 1, \dots, k; \delta_{ka} = \dot{\partial}_{ka}). \end{aligned} \tag{2.6}$$

The system of functions

$$\left(\overset{\alpha}{L}_{(00)}{}^a{}_{bc}, \overset{\alpha}{L}_{(\beta 0)}{}^a{}_{bc}, \overset{\alpha}{C}_{(01)}{}^a{}_{bc}, \overset{\alpha}{C}_{(\beta 1)}{}^a{}_{bc}, \dots, \overset{\alpha}{C}_{(0k)}{}^a{}_{bc}, \overset{\alpha}{C}_{(\beta k)}{}^a{}_{bc} \right), \quad (\alpha = 0, 1, \dots, k; \quad \beta = 1, \dots, k),$$

are the coefficients of D and

$$\left(\overset{0}{L}_{(00)}{}^a{}_{bc}, \overset{\beta}{L}_{(\beta 0)}{}^a{}_{bc}, \overset{0}{C}_{(01)}{}^a{}_{bc}, \overset{\beta}{C}_{(\beta 1)}{}^a{}_{bc}, \dots, \overset{0}{C}_{(0k)}{}^a{}_{bc}, \overset{\beta}{C}_{(\beta k)}{}^a{}_{bc} \right), \quad (\beta = 1, \dots, k),$$

are the coefficients of an N -linear connection $D\Gamma(N)$ on $T^k M$.

Also, we will denote the coefficients of $D\Gamma(N)$ with

$$\left(\overset{H}{L}_{(00)}{}^a{}_{bc}, \overset{V_\beta}{L}_{(\beta 0)}{}^a{}_{bc}, \overset{H}{C}_{(01)}{}^a{}_{bc}, \overset{V_\beta}{C}_{(\beta 1)}{}^a{}_{bc}, \dots, \overset{H}{C}_{(0k)}{}^a{}_{bc}, \overset{V_\beta}{C}_{(\beta k)}{}^a{}_{bc} \right), \quad (\beta = 1, \dots, k).$$

It is not difficult to prove

Theorem 2.2. *There exist metrical N -linear connection $D\Gamma(N)$ on $T^k M$ with respect to the homogeneous prolongation $\overset{\circ}{\mathbf{G}}$, which depend only of the fundamental function $F(x, y^{(1)})$ of the Finsler space F^n . One of them has the "horizontal" coefficients:*

$$\begin{aligned} \overset{H}{L}_{(00)}{}^a{}_{bc} &= \frac{1}{2} g^{ad} (\delta_b g_{dc} + \delta_c g_{bd} - \delta_d g_{bc}), \\ \overset{V_\beta}{L}_{(\beta 0)}{}^a{}_{bc} &= \frac{1}{2} g^{ad} \left(\underset{(\beta)}{\delta_b} g_{dc} + \underset{(\beta)}{\delta_c} g_{bd} - \underset{(\beta)}{\delta_d} g_{bc} \right), \quad (\beta = 1, \dots, k), \end{aligned} \tag{2.7}$$

the "v₁-vertical" coefficients:

$$\begin{aligned} \overset{V_\beta}{C}_{(\beta 1)}{}^a{}_{bc} &= \frac{1}{2} g^{ad} (\delta_{1b} g_{dc} + \delta_{1c} g_{bd} - \delta_{1d} g_{bc}), \\ \overset{V_\beta}{C}_{(\beta 1)}{}^a{}_{bc} &= \frac{1}{2} g^{ad} \left(\underset{(\beta)}{\delta_{1b}} g_{dc} + \underset{(\beta)}{\delta_{1c}} g_{bd} - \underset{(\beta)}{\delta_{1d}} g_{bc} \right), \quad (\beta = 1, \dots, k), \end{aligned} \tag{20}$$

and the "v_γ-vertical" coefficients vanish:

$$\overset{H}{C}_{(0\gamma)}{}^a{}_{bc} = \overset{V_1}{C}_{(1\gamma)}{}^a{}_{bc} = \dots = \overset{V_k}{C}_{(k\gamma)}{}^a{}_{bc} = 0, \quad (\gamma = 2, \dots, k). \tag{2.9}$$

Let us remark the particular form of the metrical N -linear connection $D\Gamma(N)$ in (2.7), (2.8) and (2.9). Because it depends only on the fundamental function $F(x, y^{(1)})$ of the Finsler space F^n , $D\Gamma(N)$ from the Theorem 2.2 will be called the **canonical** metrical N -linear connection of the space $\left(\widetilde{T^k M}, \overset{\circ}{\mathbf{G}} \right)$.

Let us denote

$$\sigma_c = -\frac{1}{2} \frac{1}{F^2} \delta_c F^2, \quad \tau_c = -\frac{1}{2} \frac{1}{F^2} \delta_{1c} F^2. \tag{2.10}$$

We obtain

Theorem 2.3. *The coefficients of the metrical N -linear connection $D\Gamma(N)$ with respect to $\overset{\circ}{\mathbf{G}}$ given by (2.1), satisfy the following equations:*

$$\begin{aligned} \overset{V_\beta}{C}_{(\beta 0)}^a{}_{bc} &= \overset{H}{L}_{(00)}^a{}_{bc} + \beta (\delta_b^a \sigma_c + \delta_c^a \sigma_b - g_{bc} g^{ad} \sigma_a), \\ \overset{V_\beta}{C}_{(\beta 1)}^a{}_{bc} &= \overset{V_\beta}{C}_{(01)}^a{}_{bc} + \beta (\delta_b^a \sigma_c + \delta_c^a \sigma_b - g_{bc} g^{ad} \sigma_a), \quad (\beta = 1, \dots, k). \end{aligned} \tag{2.11}$$

Indeed, substituting the tensors $\overset{g}{g}_{(1)ab}, \overset{g}{g}_{(2)ab}, \dots, \overset{g}{g}_{(k)ab}$ given by (2.5) in (2.7) and (2.8) and using (2.10), one obtains (2.11).

It is not difficult to prove

Theorem 2.4. *The coefficients of the canonical metrical N -linear connection $D\Gamma(N)$ = $\left(\overset{H}{L}_{(00)}^a{}_{bc}, \overset{V_\beta}{L}_{(\beta 0)}^a{}_{bc}, \overset{H}{C}_{(01)}^a{}_{bc}, \overset{V_\beta}{C}_{(\beta 1)}^a{}_{bc}, \dots, \overset{H}{C}_{(0k)}^a{}_{bc}, \overset{V_\beta}{C}_{(\beta k)}^a{}_{bc} \right)$, $(\beta = 1, \dots, k)$ satisfy the equations*

$$\begin{aligned} \overset{V_1}{L}_{(10)}^a{}_{bc} &= \overset{H}{L}_{(00)}^a{}_{bc} + \delta_b^a \sigma_c + \delta_c^a \sigma_b - g_{bc} g^{ad} \sigma_a, \\ \overset{V_2}{L}_{(20)}^a{}_{bc} &= \overset{V_2}{L}_{(10)}^a{}_{bc} + \delta_b^a \sigma_c + \delta_c^a \sigma_b - g_{bc} g^{ad} \sigma_a, \\ &\dots\dots\dots \\ \overset{V_k}{L}_{(k0)}^a{}_{bc} &= \overset{V_{k-1}}{L}_{(k-10)}^a{}_{bc} + \delta_b^a \sigma_c + \delta_c^a \sigma_b - g_{bc} g^{ad} \sigma_a, \\ \overset{V_1}{C}_{(11)}^a{}_{bc} &= \overset{H}{C}_{(01)}^a{}_{bc} + \delta_b^a \sigma_c + \delta_c^a \sigma_b - g_{bc} g^{ad} \sigma_a, \\ \overset{V_2}{C}_{(21)}^a{}_{bc} &= \overset{V_1}{C}_{(11)}^a{}_{bc} + \delta_b^a \sigma_c + \delta_c^a \sigma_b - g_{bc} g^{ad} \sigma_a, \\ &\dots\dots\dots \\ \overset{V_k}{C}_{(k1)}^a{}_{bc} &= \overset{V_{k-1}}{C}_{(k-11)}^a{}_{bc} + \delta_b^a \sigma_c + \delta_c^a \sigma_b - g_{bc} g^{ad} \sigma_a, \\ \overset{H}{C}_{(0\gamma)}^a{}_{bc} &= \overset{V_1}{C}_{(1\gamma)}^a{}_{bc} = \dots = \overset{V_k}{C}_{(k\gamma)}^a{}_{bc} = 0, \quad (\gamma = 2, \dots, k). \end{aligned} \tag{2.12}$$

The particular form (2.12) of the canonical metrical N -linear connection shows that the curvature of the v_k -connection $\left(\overset{V_k}{L}_{(k0)}^a{}_{bc}, \overset{V_k}{C}_{(k1)}^a{}_{bc}, \dots, \overset{V_k}{C}_{(kk)}^a{}_{bc} \right)$ lead to the Weyl's conformal curvature tensor with respect to the curvature of the v_{k-1} -connection $\left(\overset{V_{k-1}}{L}_{(k-10)}^a{}_{bc}, \overset{V_{k-1}}{C}_{(k-11)}^a{}_{bc}, \dots, \overset{V_{k-1}}{C}_{(k-1k)}^a{}_{bc} \right), \dots,$ and the curvature of the v_1 -connection $\left(\overset{V_1}{L}_{(10)}^a{}_{bc}, \overset{V_1}{C}_{(11)}^a{}_{bc}, \dots, \overset{V_2}{C}_{(1k)}^a{}_{bc} \right)$ lead to the Weyl's conformal curvature of the h -connection $\left(\overset{H}{L}_{(00)}^a{}_{bc}, \overset{H}{C}_{(01)}^a{}_{bc}, \dots, \overset{H}{C}_{(0k)}^a{}_{bc} \right)$.

This property shows the necessity to construct a gauge theory in the Asanov sense, [1], for the Riemannian metric given on $\widetilde{T^k M}$ by the prolongation $\overset{\circ}{\mathbf{G}}$, from (4.1).

Now, we remark that the almost $(k-1)n$ -contact structure \mathbf{F} defined in (1.14) has not the property of homogeneity. The $\mathcal{F}(\widetilde{T^k M})$ -linear mapping $\mathbf{F} : \chi(\widetilde{T^k M}) \rightarrow \chi(\widetilde{T^k M})$, applies the 1-homogeneous vector field δ_a into the $(1-k)$ -homogeneous vector field $\delta_{ka} = \overset{\circ}{\partial}_{ka}$, $(a = 1, \dots, n)$.

Therefore, we consider the $\mathcal{F}(\widetilde{T^k M})$ -linear mapping $\overset{\circ}{\mathbf{F}} : \chi(\widetilde{T^k M}) \rightarrow \chi(\widetilde{T^k M})$, given in the adapted basis by

$$\begin{aligned} \overset{\circ}{\mathbf{F}}(\delta_a) &= -\frac{\|y^{(1)}\|^k}{\mathbf{a}^k} \dot{\partial}_{ka}, \\ \overset{\circ}{\mathbf{F}}(\delta_{1a}) &= \dots = \overset{\circ}{\mathbf{F}}(\delta_{k-1a}) = 0, \\ \overset{\circ}{\mathbf{F}}(\dot{\partial}_{ka}) &= \frac{\mathbf{a}^k}{\|y^{(1)}\|^k} \delta_a. \end{aligned} \tag{2.13}$$

By direct calculus, we can prove:

Theorem 2.5. $\overset{\circ}{\mathbf{F}}$ has the following properties:

1. $\overset{\circ}{\mathbf{F}}$ is a tensor field of type (1.1) on $(\widetilde{T^k M})$.
2. $\overset{\circ}{\mathbf{F}}$ is an almost $(k-1)n$ -contact structure on $\widetilde{T^k M} : \mathbf{F}^3 + \mathbf{F} = 0$.
3. $\overset{\circ}{\mathbf{F}}$ depends only the fundamental function $F(x, y^{(1)})$ of the Finsler space F^n .
4. $\overset{\circ}{\mathbf{F}}$ is homogeneous on the fibres on $\widetilde{T^k M}$.
5. The pair $(\overset{\circ}{\mathbf{G}}, \overset{\circ}{\mathbf{F}})$ is a metrical $(k-1)n$ -contact structure on $\widetilde{T^k M} :$

$$\overset{\circ}{\mathbf{G}}(\overset{\circ}{\mathbf{F}}X, Y) = -\overset{\circ}{\mathbf{G}}(X, \overset{\circ}{\mathbf{F}}Y), \quad \forall X, Y \in \chi(\widetilde{T^k M}).$$

The space $(\widetilde{T^k M}, \overset{\circ}{\mathbf{G}}, \overset{\circ}{\mathbf{F}})$ is the **geometrical model** of the Finsler space $F^n = (M, F)$, with respect to the homogeneous lift $\overset{\circ}{\mathbf{G}}$ given by (2.1). It can be used for studying the Finslerian higher order gauge theory and, in general, the geometry of the Finsler space $F^n = (M, F)$.

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